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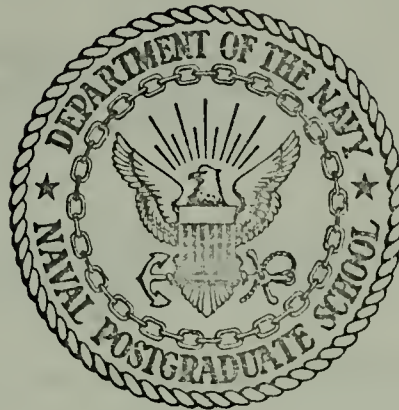
HOLOGRAPHIC INTERFEROMETRY OF THE FLOW  
FIELD BETWEEN A FIN AND FLAT PLATE

Robert Ward Heyer



# NAVAL POSTGRADUATE SCHOOL

## Monterey, California



# THESIS

Holographic Interferometry of The Flow  
Field Between a Fin And Flat Plate

by

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Thesis Advisor:

D. G. Collins

March 1972

*Approved for public release; distribution unlimited.*

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Holographic Interferometry of The Flow  
Field Between a Fin And Flat Plate

by

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Submitted in partial fulfillment of the  
requirements for the degree of

MASTER OF SCIENCE IN AERONAUTICAL ENGINEERING

from the

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## ABSTRACT

This study was an attempt to map the density field in a fin-flat plate junction three-dimensionally using holographic interferometry. This investigation has extended the density studies by Matulka [12, 13] and Jagota [3, 4] to include, for the first time, interferometric fringe information obtained through a transparent model in supersonic flow. The fringe information was then inverted by a FORTRAN computer program to produce a plot of the density field around the model. The feasibility of the method was demonstrated.

The factors which are thought to have limited the success of the experiment include vibration of the model, fluctuations of the tunnel flow and the fact that the model was somewhat too large in relation to the size of the wind tunnel test section. Schlieren photography was used to look through and around the model and to verify that the same flow was established as was reported by Thomas [23, 24] and Winkelmann [26, 27].

The data reduction of holographic interferograms was, for the first time, accomplished using photographic enlargements. This technique is considered to be much easier and more accurate than the one used in the previous investigations. However, the data reduction step, because of the time and labor involved, is considered to be the rate controlling process of the whole analysis.





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## I. INTRODUCTION

The determination of the flow field in the wing-body junction of an aircraft in supersonic flight presents many problems. The typical approach has been to measure the static pressure on the surface using many small pressure taps [23, 24] and to determine the velocity field in the junction using translatable pressure probes [18, 19, 25]. Thomas [23, 24] and Winkelmann [26, 27] used azobenzene and oil-smear tests to view the flow field around a flat plate-fin junction. From the streak patterns on the plate and fin, they were able to illustrate the three-dimensional flow field, although in a partially speculative fashion. This study has attempted to map the density field in a fin-flat plate junction three-dimensionally using holographic interferometry. Although this objective was not fully achieved, the feasibility of the method has been demonstrated.

By using a Q-switched laser with exposure times of about twenty nanoseconds, it was possible to obtain three-dimensional holographic interferograms of the density field in the fin-flat plate junction. From holograms taken at a number of viewing angles the fringe shifts in different planes could be obtained. By integrating this information using a FORTRAN computer program, the density field can be determined. This technique has been previously demonstrated for the flow field of a free jet by Matulka [12, 13] and for the supersonic flow field around a cone at angle of attack by Jagota [3, 4].

The tests were performed at the Naval Postgraduate School, using the four-inch supersonic wind tunnel.





## II. EXPERIMENTAL APPARATUS

### A. THE WIND TUNNEL

The investigation was conducted in the Naval Postgraduate School blowdown-to-atmosphere supersonic wind tunnel. The test section is four inches by four inches in cross section and six inches long with two different sets of side walls. The two-inch thick plexiglas side walls, which have a refraction index of 1.49, present a complete field of view of the flow from the nozzle throat to aft of the test section mounting bracket (Figure 1 (a)). The second set of sidewalls used are aluminum with high quality optical glass portholes located in the test section area (Figure 1 (b)). The interchangeable nozzle for a test section Mach number of 2.8 was used for all tests. The nominal run time is five minutes at Mach 2.8 with the maximum stagnation pressure of about 105 pounds per square inch.

### B. THE HOLOGRAPHIC ARRANGEMENT

The holographic arrangement is illustrated in Figure 2 and shown in photographs included as Figures 3, 4, and 5. The equipment stand was rested on a portion of the building floor that was vibrationally isolated. A Konrad K-1 pulsed ruby laser with a Pockels cell Q-switching unit was used to produce monochromatic light at a wave length of 6943 Angstroms and exposure time of twenty nano-seconds. The laser cavity length was seventy-three cm. giving a coherence length of about ten cm. To maintain the laser head and output etalon at a constant temperature of 27.5 degrees centigrade, a Lauda constant temperature circulator Model N was used. This was controlled by an electronic relay type R-10 coupled with a Culligan de-ionizer.



Holograms were obtained by routing the reference beam under the wind tunnel and the scene beam through the test section, and intersecting the two beams on the hologram plate at an intersecting angle of approximately 50 degrees. The beam sizes were controlled by translating the concave lenses located between the beam splitter and hologram plate in each beam (Figure 2). The Q-switched laser and optics were aligned using a continuous wave helium-neon laser. For reference purposes, grids were mounted on the outside tunnel walls and aligned using a surveyor's transit. The holographic stand and test section were completely enclosed in a wooden box to enable holograms to be taken in the daylight (Figure 6).

### C. THE WIND TUNNEL MODELS

The fin-flat plate models used are shown in Figures 7 (a), (b), (c), and (d). The metal portions of both models were stainless steel. The center section, part of one strake, and all of the other strake of the first model in Figure 7 (a) were made of epoxy while the center section of the second model, Figures 7 (b) and 7 (c), was fabricated from plexiglas. The flat plate grids in both models were etched into the plastic surface and coated with a clear plastic to achieve a smooth surface.

The models were rotated about their sting mounts as shown in Figure 7 (d). Alignment for the desired rotation angle was accomplished by aligning prescribed lines on the sting mount collar with a scribed mark on the sting stand using a surveyor's transit.



### III. ANALYTICAL EVALUATION OF THE DENSITY FIELD

#### A. THE BASIC INTERFEROMETRIC EQUATION

Interferograms are created when two coherent light beams are superimposed on each other and projected on a viewing screen. The light and dark regions observed correspond to the relative phase difference between the two beams which are caused by a difference in the two optical path lengths. Consider a coherent beam which is split and then recombined on a viewing screen. A difference in optical path lengths of the two beams may be achieved in two ways in order to create an interferogram. The first is to make the physical distance traveled by the two beams different. In a vacuum this path length difference is expressed as  $L = C_0 \Delta t$  where  $C_0$  is the speed of light in a vacuum. The second way is to maintain equal physical path lengths but to have the beams traverse through different media prior to recombining. In this case each light beam will travel at a speed  $\frac{C_0}{n}$  where  $n$  is the index of refraction for the medium traversed. The optical path length difference then becomes:

$$\Delta L = L (n_2 - n_1) = C_0 \Delta t \quad (1)$$

The interference pattern or fringes observed may be expressed as a function of the optical path length difference or

$$g = \frac{\Delta L}{\lambda} \quad (2)$$

where:  $g$  = fringe shift

$\lambda$  = wave length of the light source

$\Delta L$  = change in optical path



Combining equations (1) and (2), the fringe shift is then

$$g = \frac{L}{\lambda} (n_2 - n_1) \quad (3)$$

The index of refraction is known to be a function of density. Since the speed of light is only slightly less in gases than in a vacuum, the index of refraction could be closely approximated by the series expansion [8]

$$n = 1 + \beta \frac{\rho}{\rho_s} \quad (4)$$

where

$\beta$  = dimensionless constant related to the Gladstone-Dale constant by  $K = \beta/\rho_s$

$\rho_s$  = reference density of 0° C, 760 mm. Hg.

The variation of  $\beta$  with wavelength is small and has a value of 0.000292 for  $\lambda = 5893$  angstroms.

Considering a fixed difference in the index of refraction between the two beams in Equation (3), then

$$g = \beta \frac{L}{\lambda} \left( \frac{\rho_2 - \rho_\infty}{\rho_s} \right) \quad (5)$$

If the density varies in a beam path, the net change in the optical path length will be the integrated effect along the beam path or

$$g = \frac{\beta}{\lambda \rho_s} \int_0^L (\rho - \rho_\infty) ds = Q \int_0^L f(x, y, z_c) ds \quad (6)$$

where

$$Q = \frac{\beta \rho_\infty}{\lambda \rho_s} \quad (6a)$$





$$f(x, y, z_c) = \frac{\rho(x, y, z_c)}{\rho_\infty} - 1 \quad (6b)$$

$z_c$  = a plane of constant  $z$

$ds$  = incremental distance along the ray

In order to determine the density along the beam path where the fringe shift is known from an interferogram, Equation (6) must be inverted.

## B. THE INTEGRAL INVERSION

The integral inversion technique was first reported by C. D. Maldonado et al in 1965 [9, 10, 11]. R.D. Matulka [12, 13] and R. C. Jagota [3, 4] used this method to determine the density variation in an asymmetric free jet and about a cone at angle of attack, respectively. The technique involves representing the function,  $f(x, y, z_c)$  in Equation (6) by a complete set of orthogonal functions where the unknown coefficients are evaluated using the orthogonality relationship between the set of functions. The functions are orthogonal over the entire plane and also have the property of being invariant in form to any rotation of the coordinate system. Figure 8 illustrates the coordinate system for the inversion where  $x$  and  $y$  are the fixed laboratory coordinates and  $x'$  and  $y'$  are the coordinates in which the fringe number function is defined. As the view through the test section is varied the primed coordinates are rotated with respect to the fixed coordinates  $x$  and  $y$ .

The fringe shift expressed in Equation (6) may be written as the transform

$$g(\xi, y, z_c) = f(x, y, z_c) \quad (7)$$



or, inverting the equation, the density function,  $f$ , is equal to:

$$f(x, y, z_c) = T_g^{-1}(\xi, y, z_c) \quad (8)$$

The density function can be expanded in the following manner using a set of polynomial functions,  $U_{m+2k}^{\pm m}(\alpha x, \alpha y)$ , and unknown complex coefficients,  $C_{m+2k}^{\pm m}(\alpha)$

$$f(x, y, z_c) = \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \epsilon_m \left[ C_{m+2k}^m(\alpha) U_{m+2k}^m(\alpha x, \alpha y) + C_{m+2k}^{-m}(\alpha) U_{m+2k}^{-m}(\alpha x, \alpha y) \right] e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \quad (9)$$

where

$$\epsilon_m = \begin{cases} 1/2 & m=0 \\ 1 & m=1, 2, 3, \dots \end{cases}$$

$\alpha$  = arbitrary scale factor

The polynomial functions,  $U_{m+2k}^{\pm m}(\alpha x, \alpha y)$ , are invariant in form to a rotation of the coordinate system [10, 12, 13]. They also have a Gauss transform which makes them adaptable to the physical situation and to manipulating into the form of Equation (7). The functions are defined as:

$$U_{m+2k}^{\pm m}(\alpha x, \alpha y) = (-1)^k \alpha \left[ \frac{k! (\alpha^2 x^2 + \alpha^2 y^2)^m}{\pi (m+k)!} \right]^{\frac{1}{2}} e^{\pm i m \phi} L_k^m(\alpha^2 x^2 + \alpha^2 y^2) \quad (10)$$

where  $\phi = \tan^{-1}\left(\frac{y}{x}\right) - \frac{\pi}{2}$  (10a)

$L_k^m$  = Laguerre polynomial

$$= \sum_{s=0}^k \left[ \frac{(m+k)!}{(k-s)!(m-s)!s!} \right] \left[ (-1)^s (\alpha^2 x^2 + \alpha^2 y^2) \right]^s \quad (10b)$$

And the Gauss Transform of  $U_{m+2k}^{\pm m}$  is:

$$I_{m+2k}^{\pm m}(\alpha y'; \xi) = \int_{-\infty}^{\infty} U_{m+2k}^{\pm m}(\alpha x, \alpha y) e^{-\alpha^2 x'^2} dx' = \frac{e^{\pm i m \xi} H_{m+2k}(\alpha y')}{[k! (m+k)!]^{\frac{1}{2}} 2^{m+2k}} \quad (11)$$

where  $H_{m+2k}(\alpha y') =$  Hermite polynomials

By applying the transform above to Equation (9), the fringe function in



Equation (7) can be written as:

$$g(\xi, y, z_c) = \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \frac{\epsilon_m [C_{m+2k}^m(\alpha) e^{im\xi} + C_{m+2k}^{-m}(\alpha) e^{-im\xi}] H_{m+2k}(\alpha) e^{-\alpha^2 y'^2}}{[k!(m+k)! 2^{2(m+k)}]^{1/2}} \quad (12)$$

using the following orthogonality relationship on Equation (12)

$$\int_{-\pi}^{\pi} e^{\pm im\xi} e^{\mp in\xi} d\xi \int_{-\infty}^{\infty} H_{m+2k}(\alpha y') H_{n+2l}(\alpha y') e^{-\alpha^2 y'^2} dy' = \frac{2\pi}{\alpha} \left[ (m+2k)!(n+2l)! 2^{m+2k} 2^{n+2l} \delta_{mn} \delta_{(m+2k)(n+2l)} \right] \quad (13)$$

where  $\delta$  is the kroneker delta, the expansion coefficients  $C_{m+2k}^{\pm m}$ , can be determined by:

$$C_{m+2k}^{\pm m}(\alpha) = \frac{\alpha}{2\pi^{3/2}} \left[ \frac{(k!(m+k)!)^{1/2}}{(m+2k)!} \right] \int_{-\pi}^{\pi} g(y', \xi, z_c) H_{m+2k}(\alpha y') e^{\mp im\xi} dy' d\xi \quad (14)$$

Substitution of the coefficients in Equation (14) back into Equation

(9) results in the density variation being expressed as:

$$f(x, y, z_c) = \left(\frac{\alpha}{\pi}\right)^2 \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \epsilon_m [k!(m+k)!]^{1/2} e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \cdot \text{REAL} \left[ \int_{-\pi}^{\pi} \int_{-\infty}^{\infty} g(y', \xi, z_c) e^{-im\xi} H_{m+2k}(\alpha y') dy' d\xi \right] U_{m+2k}^m(\alpha x, \alpha y) \quad (15)$$

or by inserting Equation (10):

$$f(x, y, z_c) = \left(\frac{\alpha}{\pi}\right)^2 \sum_{m=0}^{\infty} \sum_{k=0}^{\infty} \epsilon_m \frac{(-1)^k k!}{(m+2k)!} (\alpha^2 x^2 + \alpha^2 y^2)^{m/2} L_k^m(\alpha^2 x^2 + \alpha^2 y^2) \cdot [B_{m+2k}^m(\alpha) \cos(m\phi) + D_{m+2k}^m(\alpha) \sin(m\phi)] e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \quad (16)$$

where:

$$B_{m+2k}^m(\alpha) = \int_{-\pi}^{\pi} \int_{-\infty}^{\infty} g(y', \xi, z_c) \cos(m\xi) H_{m+2k}(\alpha y') dy' d\xi \quad (17)$$

$$D_{m+2k}^m(\alpha) = \int_{-\pi}^{\pi} \int_{-\infty}^{\infty} g(y', \xi, z_c) \sin(m\xi) H_{m+2k}(\alpha y') dy' d\xi \quad (18)$$



Equations (16), (17), and (18) are the basic equations used to calculate the density distribution from the experimentally determined fringe variations.

### C. THE NUMERICAL PROCEDURE

The form of the density distribution in Equations (16), (17), and (18) must be modified in order to input the experimentally determined fringe distribution. First from Figure 8 and Equation (6b) it can be seen that it is only necessary to integrate Equations (17) and (18) over an area where the density is changing from a known density,  $\rho_{\infty}$ . Outside of this region where there is no change in density, the function  $f(x,y,z_0) = 0$ , i.e. outside the test section. Also since the fringe distribution is taken in small increments over the test area the coefficients, B and D, can be approximated as

$$B_{m+2k}^m(\alpha) = \sum_{i=1}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta\xi_j, x_i + \Delta x_i) \int_{\xi_j}^{\xi_{j+1}} \cos(m\xi) d\xi \int_{x_i}^{x_{i+1}} H_{m+2k}(\alpha x) dx \quad (19)$$

$$D_{m+2k}^m(\alpha) = \sum_{i=1}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta\xi_j, x_i + \Delta x_i) \int_{\xi_j}^{\xi_{j+1}} \sin(m\xi) d\xi \int_{x_i}^{x_{i+1}} H_{m+2k}(\alpha x) dx \quad (20)$$

The integral of  $\xi$  is easily determined and by using the derivative formula for Hermite polynomials the integral of  $x$  may be manipulated to yield

$$B_{m+2k}^m(\alpha) = \left[ \frac{1}{2\alpha m(m+2k+1)} \right] \sum_{i=0}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta\xi_j, x_i + \Delta x_i) \cdot [\sin(m\xi_{j+1}) - \sin(m\xi_j)] [H_{m+2k+1}(\alpha x_{i+1}) - H_{m+2k+1}(\alpha x_i)] \quad (21)$$

$$D_{m+2k}^m(\alpha) = - \left[ \frac{1}{2\alpha m(m+2k+1)} \right] \sum_{i=0}^{I-1} \sum_{j=0}^{J-1} g(\xi_j + \Delta\xi_j, x_i + \Delta x_i) \cdot [\cos(m\xi_{j+1}) - \cos(m\xi_j)] [H_{m+2k+1}(\alpha x_{i+1}) - H_{m+2k+1}(\alpha x_i)] \quad (22)$$

In the computation of the density function from Equation (16), obtaining the infinite summations experimentally is not plausible or





possible. It has been demonstrated that by using a finite number of terms and by adjusting the values of  $\Delta \xi$ ,  $\Delta x$  and  $\alpha$ , it is possible to obtain the density distribution with very good accuracy [3, 4, 12, 13]. Equation (16) then becomes:

$$\begin{aligned} f(x, y, z_c) = \left(\frac{\alpha}{\pi}\right)^2 \sum_{k=0}^K \sum_{m=0}^M \epsilon_m (-1)^k \left[ \frac{k!}{(m+2k)!} \right] (\alpha^2 x^2 + \alpha^2 y^2) \cdot \\ L_k^m(\alpha^2 x^2 + \alpha^2 y^2) \left[ B_{m+2k}^m(\alpha) \cos(m\phi) + D_{m+2k}^m(\alpha) \sin(m\phi) \right] \cdot \\ e^{-(\alpha^2 x^2 + \alpha^2 y^2)} \end{aligned} \quad (23)$$

#### IV. EXPERIMENTAL PROCEDURE

##### A. LABORATORY TECHNIQUES

The analysis of a free jet by Matulka [12, 13] illustrated how holographic interferometry can be used to obtain a complete three dimensional plot of the density within a moving transparent flow field. Jagota [3, 4] in his study of a cone at angle of attack in supersonic flow went one step further by introducing an opaque object into an assumed steady state flow field and describing the density field three-dimensionally.

This investigation has attempted to determine the three-dimensional density field around a transparent object in a supersonic flow field by passing a light beam through the object. Specifically the interest was to describe the flow field existing in the junction of a fin-root intersection.

##### 1. Model Considerations

In order to obtain uniform flow around the fin-root area, a model of the form shown in Figure 9 was selected. The flat plate has a knife



edge and is intended to remain at zero degrees angle of attack so as to establish the flow conditions illustrated. The fin edges were made circular in order to approach flow conditions similar to those established in Winkelmann's [26, 27] investigation. Plastic and metal strakes were added on the model sides so as to maintain two-dimensional flow as well as to add strength to the flat plate. In the first model constructed (Figure 7(a)), maximum visibility of the plastic fin-flat plate center section was achieved by bolting the leading edge and aft plate together through the plastic center section. Due to model flexure, this design was found to be unsuitable and the second model in Figures 7(b) and 7(c) was constructed. The model was made from a single piece of stainless steel. The model rigidity was satisfactory but unfortunately the strengthening borders around the plastic center section reduced the holographic visibility somewhat.

Since the wind tunnel blocks were fixed, it was possible to make multiple test runs with the same flow conditions over the model provided supersonic flow had been established over the model.

## 2. Holographic Techniques

In order to obtain holographic interferograms it was necessary to ensure that the optical path lengths of the scene and reference beams remained approximately equal. Since the ruby laser is believed to have a coherency length of approximately ten centimeters, the equality of lengths is far less critical than in the classical Mach-Zehner interferometric approach. Consequently a string was used in the experiment to trace the reference beam and then adjust the scene beam. This method kept the two beam lengths within one centimeter of each other. Since the scene beam traversed approximately 4.5 inches of plastic tunnel walls and



grids which the reference beam did not, it was necessary to compensate by making the scene beam physically 2.25 inches shorter than the reference beam. The reference beam varied in length from 61 inches to 68 inches during the experimentation.

To determine the fringe/density field, finite fringe interferograms were made by three different techniques. In the direct fin-root flow approach the diffuser plate was part of the model. Either the mirror,  $M_5$ , in Figure 2 or the hologram plate holder was translated between the no-flow exposure and the flow-established exposure. In the total model flow method the diffuser plate was located between scene beam lens,  $L_3$ , and the test section (See Figure 2) and it was translated horizontally or vertically. Translations were varied from .001 inches to .006 inches with the translation distance of .003 inches yielding the best fringe separation.

Most of the holograms taken using basically the holographic arrangement shown in Figure 2 gave well-defined fringe patterns. In order to improve upon the fringe definition, a variety of techniques was attempted. The transverse mode selector was varied from 1.0 mm. to 2.5 mm. in increments of 0.5 mm. to determine the best lighting of the model. The hologram plate holder was rotated horizontally to various positions. These positions varied between being perpendicular to the scene beam to being perpendicular to the bisection of the angle between the scene and reference beams. Polarization plates were added in both the scene and reference beams between the test section and hologram plate in the scene beam and between the last mirror,  $M_4$ , and the hologram plate in the reference beam. A one-quarter wave plate was also placed between the first lens,  $L_1$ , and the beam splitter as recommended by Okayama and Emori [14].



The holograms were taken using 4" x 5" Agfa-Gavaert 8E-75 hologram plates. The developing process involved:

1. Five minutes in Kodak D-19 developer
  2. Thirty seconds in an acetic acid stop bath
  3. Five minutes in standard fixer
  4. Five minutes in a flowing water bath
  5. One minute emersion in Kodak Photo Flo wetting agent
  6. Drying using blowing cool air
- a. Direct Fin-Root Flow Method

Since the flow field in the fin-root junction is assumed to be identical on either side of the fin, then it is only necessary to determine the density on one side of the fin. To accomplish this it is necessary to obtain holograms for  $180^\circ$  of view as shown in Figure 10. The holograms for the views from  $0^\circ$  to  $90^\circ$  can be obtained by using a frosted fin and flat plate as shown in Figure 11. The advantage of having the frosted plate as part of the model is that the fringe/density information obtained by the interferogram is believed to be only for the area between the fin-root intersection to the tunnel wall vice the whole test section, but this was not verified. In order to obtain the fringe information for angles greater than  $90^\circ$  but less than  $180^\circ$  the fin would be exchanged for one containing a stainless steel reflective surface. The scene beam would then enter the test area from the viewing port below the tunnel and be reflected to the hologram plate as shown in Figure 12.

- b. Total Model Flow Method

In this method the diffuser plate was located in the scene beam outside the test section as shown in Figure 14 and the flat plate





center section and fin were made of optically clear plexiglas. By translating the diffuser plate between exposures of the hologram, an interferogram of the whole density field in the test section about the model was obtained. From Figure 13 it can be seen that due to symmetry only  $90^\circ$  of view was required to obtain the density field. This makes it much easier experimentally to take the holograms than the previous method described.

### 3. Schlieren Analysis

A standard Schlieren knife-edge system was used to verify the establishment of the supersonic flow network around the model as shown in Figure 9. Photographs were also taken of the flow with the model at  $0^\circ$  and  $90^\circ$  rotation in order to compare the fin shock conditions with those obtained by Winkelmann [26, 27].

#### B. PHOTOGRAPHIC TECHNIQUES

By illuminating the holograms with a helium-neon laser beam which has a wave length of 6328 Angstroms, the original scene was reconstructed. Since the original scene beam and the reconstructed beam were of different wave lengths, there is actually a small distortion in the reconstructed scene but of neglectable effect because the hologram plate emulsion in the development process also shrinks.

The typical method for reconstructing the scene is to illuminate the hologram as illustrated in Figure 15. The diffuse glass used in the construction of the hologram appears to act as an infinite light source of non-parallel rays which illuminates the scene. If a small aperture is positioned at the focal plane of the imaging lense, an almost parallel set of rays may then be selected as shown in Figure 16. A third method



illustrated in Figure 17 uses a small diameter conjugate beam to illuminate the hologram. A large depth of field is achieved because the narrow beam acts as an aperture. This effect was of considerable advantage since it enabled both the front and rear grids, the model, and the fringe patterns to be simultaneously projected on the screen. The best photographs were obtained by focusing on the plane of the fringes.

### C. DATA REDUCTION

Photographic interferograms were obtained by illuminating the hologram scene with a thin laser beam and using a camera with a viewing screen located in the film plane as shown in Figure 15. The line of sight in the plane desired was achieved by translating the hologram until common points on the front and rear grids were aligned. The camera, with the aperture set wide open at  $f/7.7$ , was then adjusted to give the best focus on the fringe plane. The best photographic results were achieved by using an exposure time of  $1/10$  second with Polaroid Type 55 P/N film.

It was felt that the density field could be well defined three-dimensionally along the fin if the density fields were determined in four planes perpendicular to the Z-axis and equally spaced along the fin as shown in Figure 18. In obtaining the fringe data across a constant Z-plane it would be necessary to take six photographs per rotation angle, aligned at appropriate intervals down the  $y'$  axis, and then graphically mate the fringe data to form one complete set. The six photos across the field were felt necessary because the optical path length from the model to the hologram plate varied for those points not on the aligned plane.

The fringe shift reduction was accomplished using two different techniques. The first was to project the negative, using a photo enlarger,



onto a sheet of paper and trace out the fringe pattern, model surfaces and grid lines. The light fringes were traced out since it was much easier to judge their center line. From the fringe lines forward of the fin in the region of uniform flow one fringe line which appeared the straightest and compatible with most others was selected as the bench mark. A straight line was then drawn over that fringe and extended past the aligned Z plane (i.e.  $y'$  axis). Lines parallel to the bench mark line were then drawn likewise over the remaining fringe lines. The fringe displacements were then read relative to the lines drawn at the points of intersection of the fringes with the  $y'$  axis. The radius of the inversion circle was selected so that the fin-root intersection was the origin and the fin tip was the 100 percent point.

In the second technique an enlarged positive photograph was made from the negative. Again one fringe line in the uniform flow region just forward of the fin which appeared to be parallel with the majority of the fringes was selected as the bench mark. The remaining fringes were likewise traced over with lines parallel to the first. The fringe displacements were then measured relative to the lines drawn. For further details see Appendix A.

The locations at which reference lines crossed the  $y'$  axis in both techniques above were further adjusted to account for the tunnel wall refraction displacement as shown in Figure 19 and computed by a computer program in Appendix B. Once these corrections were made, the radial variation of the fringe number could then be plotted for the various  $y'$  alignment planes and a smooth curve drawn through the data points. The fringe number at 201 equidistant points across the field can then be obtained for input into the computer program, HOLOFER, in MODE 3. For



further details on how to use the computer program, HOLOFER, see Appendix C. Once the data from all the rotation angles of the model have been put into the computer program, the program will then calculate the density field across the inversion circle for that Z-plane. After the density has been calculated for all four Z-planes in Figure 18, a three-dimensional plot of the density field can be made by connecting points of equal density across the fin.

## V. EXPERIMENTAL RESULTS AND DISCUSSION

The initial attempts to establish uniform flow over the flat plate shown in Figure 7(a) were unsuccessful due to model vibration and flexure. Movement of the model sting within its holder and flexure of the model plastic center section allowed the model leading edge to establish a little over  $1^\circ$  angle of attack upward when flow was established. Due to various modifications made in attempts to eliminate the vibration, the sting finally fractured.

In an attempt to eliminate these problems, the second model shown in Figures 7(b) and 7(c) was made of a single piece of stainless steel and the sting was mated to its holder to within .001 inches. The model center section was made of poured epoxy and the strakes were both made of stainless steel with plexiglas inserts. The model rotation about the sting was reduced to approximately  $0.3^\circ$  angle of attack upward. Due to the vibration in passing through the transonic range and a weak glue seal between the metal strakes and plexiglas inserts, the inserts were found to break loose. They were subsequently removed and not replaced.

In order to determine the flow field using the direct fin-root approach, the epoxy fin and model center section were frosted on one side using fine





emery paper (see Figure 20). Holograms were taken of the model at rotation angles of  $0^\circ$ ,  $12^\circ$ ,  $45^\circ$  and  $90^\circ$  using the holographic arrangement in Figure 2 excluding the diffuser plate between lens,  $L_3$ , and the test section. The mirror,  $M_5$ , was translated in various directions from towards to parallel to the test section between the exposures without flow and with flow established. Fringe patterns were obtained around the model, but only at  $0^\circ$  and  $12^\circ$  rotation could any fringe patterns be observed across the fin. It was found that the fin fringe pattern appeared to remain almost unchanged no matter how or how much the mirror,  $M_5$ , was translated while the fringe around the model changed appropriately. For instance, in Figure 21 the mirror was not translated and in Figure 22 the mirror was translated .006 inches horizontally parallel to the tunnel.

Since fringes could not be observed across the flat plate at  $45^\circ$  and  $90^\circ$ , it was felt that the epoxy center section might be too imperfect optically. Consequently double-exposed holograms were taken with no flow through the test section and various translation distances from 0 to .005 inches. The fringe patterns were excellent across the whole model and their spacing decreased according to the increase in the translation of  $M_5$ . Next the diffuser plate in the scene beam in Figure 2 was inserted with the model at  $0^\circ$  rotation angle and translated between no flow and flow exposures of the hologram. The fringes about the model were of excellent quality but the double diffusion of the scene beam through the model caused all fringe patterns on the model (fin) to disappear.

It was felt at this point that the fringe pattern obtained across the fin was caused by the movement of the model to an angle of attack and possibly by model vibration, although none was observed visually. The



lack of any fringes across the flat plate center section is not well understood but is believed to be caused by vibration of the model.

Since it was not possible to obtain acceptable interferograms with the diffuser plate as part of the model due to model motion, it was felt that the effect of minor model movements could be eliminated by using an optically clear model and an external translating diffuser plate. Therefore the fin and model center sections were replaced with optically clear plexiglas. With these changes it was found that the flat-plate leading-edge angle of attack had been reduced to approximately  $0.1^\circ$ .

Initially double-exposure holograms were taken using the arrangement in Figure 2 and translating the diffuser horizontally. At  $0^\circ$  model rotation the holographic interferograms were excellent. But as observed in Figure 23 it would be extremely difficult to determine the fringe change across the fin since no free stream reference fringes were available, due to the flat plate leading edge Prandtl-Meyer expansion and fin shock intersecting the tunnel top just above the fin. With a larger tunnel or smaller model this would be an excellent technique.

The diffuser plate was then translated vertically between exposures and excellent horizontal fringe patterns were obtained as shown in Figure 24. The model was then rotated to  $22\frac{1}{2}^\circ$  and the same holographic technique was used. Excellent fringe patterns were obtained above and below the flat plate center section. Fringe patterns across the flat plate center section and fin in this area were very light and usable interferograms could not be photographed with the polaroid camera. In an attempt to improve on the fringe quality, polarizer plates were inserted in the reference beam between the mirror,  $M_4$ , and the hologram plate and in the scene beam between the lens,  $L_3$ , and the diffuser plate in order to ensure



polarization of both beams. No significant improvement could be noticed and they were subsequently removed. A one-quarter wave polarizer plate was then placed between the lens,  $L_1$  and the beam splitter in order to utilize circularly polarized light for the reference and scene beams. Okayama and Emori [14] found that their image resolution improved considerably; however, with this particular arrangement little to no improvement in the fringe resolution was observed and the approach was abandoned.

The transverse mode selector was then varied from 2.0 mm. to 1.5 mm. and later to 1.0 mm. in an effort to improve the coherency length of the laser light and consequently the image resolution. Due to the decrease in output light intensity, up to six exposures were taken during a run with flow established. Image and fringe definition were not found to increase possibly due to model vibration which was not visible to the eye.

The model was rotated to  $45^\circ$  and  $67\frac{1}{2}^\circ$  and double exposure holograms were taken with a 2.5 mm. transverse mode selector. There were no observable fringe patterns in any portion of the test section. Consequently it was believed that supersonic flow was not established due to tunnel blockage caused by the shock wave from the model and by slight model vibrations in passing through the transonic range. In the transition to supersonic flow, the model leading edge would sometimes flex as much as  $0.2^\circ$  depending upon the transition time.

The test section walls were changed from the total plexiglas side walls shown in Figure 1(a) to the aluminum walls with the optical quality glass port holes (Figure 1(b)). The better quality glass would hopefully improve the viewing and the port holes made the model much more accessible. The metal strakes were also removed from the model since they appeared to have a minimal effect on the flow, were an interference optically,





and appeared to add very little structurally. Flow tests were run and the model was found to attain an angle of attack of about  $0.14^\circ$  upward when flow was established at Mach 2.84. The angle of attack was tolerable, however the model also maintained a visible high frequency vibration of about  $\pm 0.05^\circ$ . Double exposure holograms were taken with the model at  $0^\circ$  rotation and fringes could only be observed forward of the Prandtl-Meyer expansion on the top of the model and forward of the flat plate leading edge shock on the bottom. The tunnel seals were found to leak and were replaced and resealed. The same angle of attack and high frequency model vibration still persisted. Double exposure holograms were again taken and a very faint fringe pattern was observed across the model. In reproducing the scene when aligned along a constant Z-plane passing through the fin, it was impossible to photograph the fringe pattern across the fin because of a bad speckle effect. Double exposure holograms were made at model rotation angles of  $45^\circ$  and  $90^\circ$  without and with flow. The holograms taken without flow gave excellent fringe patterns across the model however no fringes could be observed with flow.

At this point a knife edge Schlieren system was set up to determine just what the flow conditions were in the tunnel and around the model. The first runs were made at a model rotation angle of  $90^\circ$  and presented the shock knee in front of the fin (Figures 25(a) and 25 (b)) which was pointed out by Thomas [ 23, 24 ] and Winkelman [ 26, 27 ]. Due to the model vibration, the shock extending out from the shock knee was observed to flutter about  $\pm \frac{1}{2}$  cm.

The model was then rotated back to  $0^\circ$  and Schlieren photographs were taken as shown in Figure 26. During all but one of the runs at this rotation angle, the model would pitch up approximately  $0.14^\circ$  angle of attack and continue to vibrate about  $\pm 0.05^\circ$  at a high frequency when





flow was established. The plate leading edge shock and fin shock appeared quite fuzzy and light. Due to oil from the tunnel supply reservoir mixing in the flow, a light oil smear pattern can be observed across the fin in Figure 26. During one run, which could not be duplicated, the model pitched up as flow was established but did not vibrate. The plate shock and Prandtl-Meyer waves, fin shock system, and tunnel Mach lines became very distinct and well defined as seen in Figure 27. The schematic of the Schlieren photographs in Figure 28 points out the cause for the various flow lines observed. The non-uniformity of the free stream caused by tunnel leakage can be easily seen. Consequently the experimental work was discontinued due to tunnel conditions and time considerations.

It was felt at this point that if the holographic interferogram taken of the clear plexiglas fin at  $0^\circ$  rotation angle could be reduced to useful data for the computer program then the holographic method would be to a certain extent verified even though the actual density field could not yet be determined.

In obtaining interferograms from the hologram, the reconstruction technique shown in Figure 15 was used. The plane of constant  $Z$  across the model to be reduced was chosen to be the forward most vertical grid line crossing the fin. It should be pointed out that in order to simplify the hologram alignment process, all four reduction planes across the model should have been scribed on the exterior grids. The three photographic points across the  $Z$  plane on the  $y'$  axis were, for convenience, chosen to be where horizontal exterior grid lines crossed the  $y'$  axis on the fin. Five interferogram photographs were taken at points A, B, and C as shown in Figure 29. Photographs 2 and 5 were taken using a Kodak Wratten Gelatin Filter N.D. 2.0 placed in the reconstruction beam in order to reduce the



beam intensity and increase the fringe definition. These two photographs were also used to provide a check on the consistency of the reduction process.

All five negatives were blown up on the photo enlarger in the dark room and the plate, fin, grid lines and fringes were traced out on a sheet of white paper as shown in Figures 30-34. It was very difficult to trace the fringes in the region of the fin tip and fin root due to the photographic resolution and the fin tip shadow. Also fringes were not visible in the region of the fin leading edge shock. Consequently, connecting the fringe in front of the fin to the correct fringe on the fin was a best guess effort. Typical of the problem was that fringe lines in two of the drawings were initially improperly connected across the fin leading edge shock. After being checked against the photograph negatives, the fringes were reconnected and the data taken correlated well with the data from other drawings. The fringes in the fin root area were extremely light in some photographs which made it quite difficult to determine their centerline crossing the  $y'$  axis. Another point of difficulty was determining the location of the top and bottom of the fin. An error in drawing here will have an effect later when the fringe locations are normalized with respect to the fin height.

Enlarged positives were then made from the negatives as shown in Figures 35-39 to see if the data accuracy could be improved by eliminating the difficulties of tracing the fringes in the dark room. This method also made it possible usually to reassess the assumed path of the fringe lines at a later time if the data did not appear to correlate properly. The same problems of locating the fringe lines in the region of the fin tip and of connecting the fringe lines across the fin shock are readily apparent in the figures.



To obtain the fringe change across the fin one of the straightest fringes in the free stream region which paralleled the majority of other free stream fringes was selected. A straight line was drawn through its centerline and extended to cross the  $y'$  axis. The other free stream reference lines were then drawn parallel to the first and adjusted so as to best follow the centerline of the selected fringe. This method actually averages the free stream conditions and is only valid if the free stream is essentially uniform. Since it was not possible to obtain in the photographs the free stream fringe pattern forward of the leading edge Prandtl-Meyer expansion for the lower fin region, the fringe reference lines in photo enlargements were drawn over the fringes just prior to the fin shock. In order to adjust them to free stream conditions their location was moved downward a distance equal to 1.1 times the average fringe interval for that photograph. The figure 1.1 was observed in all the enlarged photographs to be the approximate fringe change across the Prandtl-Meyer expansion. In the drawings (Figures 30-34), fringe reference lines for the lower fringes (generally  $y' \leq .8$ ) were initially aligned along the straightest portion of the fringe prior to intercepting the fin shock. By comparing the reference line location on the drawing with the fringe pattern in the photographic enlargements the reference line location and fringe change were adjusted by an appropriate portion or all of the 1.1 fringe change caused by the Prandtl-Meyer expansion. All of the reference lines were then corrected for the tunnel wall and grid plastic parrallax shown in Figure 19 and computed in Appendix B. The reference fringe locations were then normalized with respect to the wing height and the fringe numbers calculated by dividing the fringe change by the average fringe interval. For further details and calculations see Appendix A.





The data obtained for alignment points A and C by the two different reduction processes are plotted in Figures 40-43. The data obtained from the photo enlargements aligned at point C and shown in Figure 43 gave the best data agreement between two photographs. The worst data agreement was obtained from the reduction of the drawings aligned at the same point. The inconsistency in the data was probably caused by connecting the wrong fringe lines across the fin leading edge shock. The data fluctuations and discrepancies between the curves in the figures could have been caused by a number of things. It could have been caused, for instance, by not drawing the fringe reference line parallel to or exactly on the free stream fringe center line or slightly missing the fringe center as it crosses the  $y'$  axis or by misconnecting fringe lines across the fin shock as was illustrated in Figure 42. Fringe location errors could be caused by misdrawing the fin tip and root lines as was mentioned earlier. Another contributor would be measurement errors.

To analyze these sources for error, first consider that the typical fin size in the drawings and photo enlargements averaged about 2.5 inches high and 2.75 inches wide and that the fringe spacing averaged about 0.10 inches. All measurements were taken using a ruler graduated in 0.01 inches and readings were made to the nearest .005 inches. Since the average difference between the data points and curves ran around  $\frac{1}{4}$  fringe, some figures were calculated to determine what measurement errors could produce this fringe error. It was found that an error of .025 inches in alignment of the fringe reference line with the free stream fringe center line and/or fringe center line crossing the  $y'$  axis could produce  $\frac{1}{4}$  fringe error. This fringe error will also occur if the fringe reference line differs from the free stream center line by more than  $1.4^\circ$  when drawn 1 inch from the  $y'$  axis or by more than  $0.36^\circ$  when drawn 4 inches from the





y' axis. A difference of .01 inches in the average fringe interval could also produce a  $\frac{1}{4}$ -fringe error; however, this is not too likely since it is an average of fifteen to twenty-five intervals. By comparing the actual height-to-width ratio with those found in all the drawings and photo enlargements it was found that average error was around 2% or a distance of .02 on the y'-axis.

In order to compare the interferogram negative quality and the two different reduction techniques, a plot of the data for each negative was made as shown in Figures 44-48. Photographs 1 and 3 in Figures 44 and 46 gave the smoothest curves indicating the highest interferogram resolution. Photograph 5 (Figure 48) gave the worst dispersion indicating poor interferogram resolution; yet looking at the enlargement in Figure 39, the fringe line contrast is very good. In comparing Figures 40-48 it appears that the reduction technique using the photographic enlargements gave the most consistent data. This technique also provides a much easier and faster recheck on the proper tracing of fringe lines because it is extremely difficult and tedious to duplicate the fringe pattern to the same scale over the drawings using the photo-enlarger in the dark room.

All the data obtained by either method for one alignment point were then plotted in Figures 49 and 50 in order to observe the dispersion. A fringe number dispersion of about 0.6 fringes was observed in the data taken from the negatives aligned at point A and a dispersion of about 0.8 fringes for the data about point C. The dispersion is attributable to the inaccuracies in the drawings, to the photographic quality of the interferograms, and to the inaccuracies in correcting the lower fringe reference lines to free stream conditions. The last point is based upon the increased dispersion between the fin tip and fin root data.



Figures 51 and 52 show an integrated curve of the fringe change across the wing as determined by each reduction method. In constructing the curve, the data from the three aligned points was plotted so as to just overlap each other. These two curves were then compared in Figure 53. The maximum variation in the fringe number is about  $\frac{1}{2}$  a fringe but the variations in the location of the fringe maximums and minimums average about 0.1 inches on the fin. The location difference is probably due to improper drawing of the fringe reference lines compounded with not being able to measure the wing height accurately. Due to the inaccuracies introduced in tracing the negative and then reducing the data it is felt that the photo enlargement method is the more accurate reduction method.

With fringe data from only one field of view, the density field, obviously, could not be obtained. It was felt useful to consider the flow field axisymmetric in order to exercise the computer program and to provide a check on the program's ability to handle these particular curve shapes. Fringe data at 101 equidistant points across the fin from  $0 \leq Y' \leq 1.0$  was obtained from Figure 53 for both curves and fed in HOLOFER in Mode 3 for the axisymmetric case. For further details on HOLOFER see Appendix C. The scale factor,  $\alpha$ , in Equation 9 was then varied from 0.2 to 2.5 and a value to 1.0 was determined to yield the most accurate density solutions. This value was verified by feeding the function data,  $(\frac{\rho}{\rho_\infty} - 1)$ , calculated by Mode 3 back into the program in Mode 1 and comparing the fringe data calculated with the original fringe data obtained from Figure 53.

The density distribution for both the drawing and photographic reduction cases is plotted in Figure 54. The density as  $Y'$  approaches zero actually goes as the drawn lines even though the points indicate a dip.



Matulka [12, 13] pointed out that the computer program accuracy does not converge at the origin. The large variations of the density curves at values of  $Y' > 0.8$  are caused by the program trying to adapt to a step or shock wave type function at  $Y' = 1.0$ . If the remainder of the fringe data in Figure 53 for values of  $Y' > 1.0$  had been included as input data the density curves would have smoothed out. This shock wave step function effect was demonstrated and analyzed by Matulka [12, 13].

The low values of density for the photographic data curve around  $Y' = 0.38$  were unexpected but not surprising. First, the photographic fringe data curve is more extreme than the drawing fringe data curve and second, the density curves are not true values anyway since the field was considered axisymmetric and this is not the case in reality.

In general the density curves are felt to be reasonable under the assumed conditions. Consequently it is believed that the computer program could very easily and accurately handle a complete analysis of the flow field around the fin-flat plate.

In completing the analysis there are some data reduction problems which would have been encountered in reducing the interferograms at other model rotation angles which merit discussion. With no model rotation, adjusting the fringe reference lines close to the fin root to account for the leading edge Prandtl-Meyer expansion was relatively easy. However, when the model is rotated, the fringe shift across the Prandtl-Meyer expansion can no longer be considered a constant and it will be very difficult to adjust the fringe lines at lower values of  $Y'$  to free stream conditions. The best solution would be to increase the scene beam diameter and/or reduce the model size in order to photograph the free stream fringe lines forward of the plate leading edge. The free





stream lines could be connected to the appropriate fringe lines forward of the fin and this would eliminate the numerical correction and increase the fringe data accuracy.

Also at the rotated angles, fringe information will not be available for portions of the reduction plane due to shadows cast by the model sides as shown in Figure 55. In the lower portion of the hologram the fringe curve can be connected with a smooth line because the density field in that portion should be essentially constant. Care must be taken in completing the curve, though, since this information will be used in the integration of other rotation angles. The fringe curve in the upper portion must also be completed carefully but should be easier since the shadow will be smaller. For this particular model it was calculated that for angles greater than  $46.6^{\circ}$  the upper shadow would not penetrate the fin.

There are two more problems to be coped with which do not have apparent solutions at this time. The first has to do with the superimposing of fringe information from different locations onto one scene beam line as illustrated in Figure 56. In case I the superimposing of the fringe change in Region A onto that in Region C and locating the fringe reference line at point C' is tolerable because the density field in Region A should be fairly constant and uniform. However in case II where the fringe change in all three regions are superimposed and located at point C'', the answer is not readily apparent. If the plate and fin were of equal thicknesses then points A'' and C'' would be the same and the error would be somewhat reduced. If, in addition, the thicknesses were made as thin as structurally possible, the error would be reduced to a minimum. Also it might be possible to integrate the model geometry into the computer program but this has not been attempted.





With the model at rotation angles between  $0^\circ$  and  $90^\circ$  scene beam defraction at the fin root and tip areas would also present problems as seen in Figure 57. A minimization of the problem could again be achieved by minimizing the fin and plate thicknesses.

## VI. CONCLUSIONS AND RECOMMENDATIONS

The investigation, although not totally successful, has demonstrated the feasibility of using holographic interferometry to determine the flow field around a transparent model by looking through the model. The problems of model vibration and movement, of superimposing different fringe information on one beam, and of scene beam divergence through the fin root and tip regions must still be solved. The model vibrations and rotation to an angle of attack are attributable to the tunnel pressure fluctuations caused by tunnel leaks and to possible movement of the sting holder. It is felt that all three problems could be decreased to neglectable effects by reducing the overall model size and by incorporating the model geometry into the computer program.

The basic holographic arrangement was found to work quite well for this type of experiment. The holograms were generally high quality except when unfavorable tunnel flow conditions existed. The use of circularly polarized light recommended by Okayama and Emori [14] did not increase the hologram resolution appreciably but the method merits further considerations because of their excellent results.

The data reduction process was found to be the rate-controlling step in the investigation due to the time and labor involved. Reducing the data from enlarged photographs of the interferogram saved some time and appeared to increase the data accuracy. The data scatter of  $\pm 1/8$  fringe



was considered acceptable considering the fringe resolution in the holograms. It was felt that this could be reduced by using either a larger tunnel or a smaller model. With either of the changes it would be possible either to use the free stream fringe pattern forward of the model as the reference conditions across the whole model or to use a vertical fringe pattern, since the leading edge Prandtl-Meyer expansion and fin shock would not block out the free stream fringes above the fin. It appears that the use of vertical fringes would also considerably reduce the data reduction time.

With the use of Schlieren the flow network described by both Thomas [23, 24] and Winkelmann [26, 27] was verified to exist. From the top view ( $90^\circ$  model rotation) the fin shock and its fluctuations were observed and photographed through the flat plate plastic center section.

The computer program, HOLOFER, was found to be quite capable of handling the type of flow field fringe data which would be generated in a complete analysis of this type. As pointed out before it is believed that the program should be modified to incorporate the model geometry.



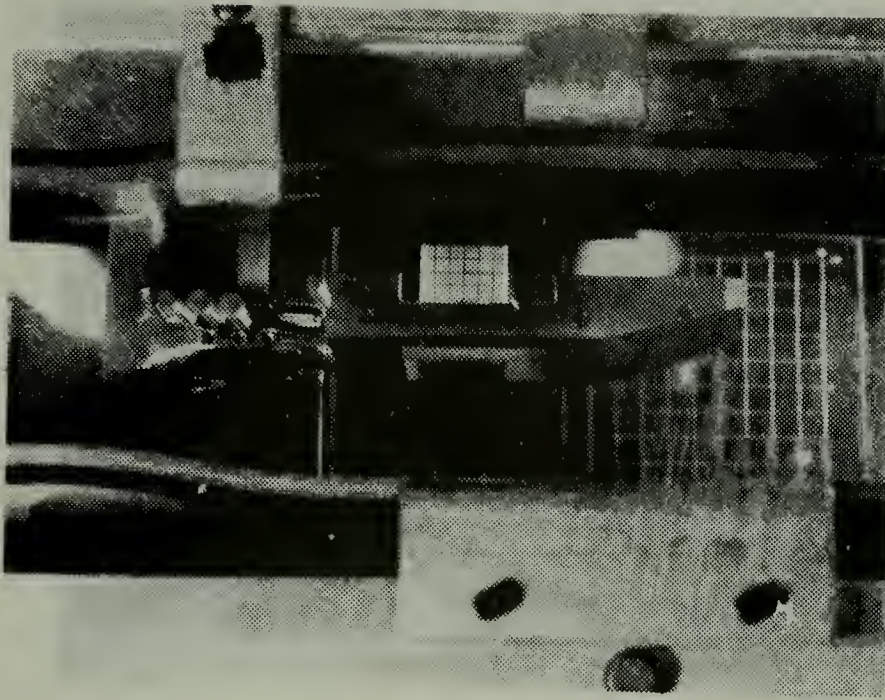


Figure 1(a). Wind Tunnel Test Section with Clear Plastic Side Walls





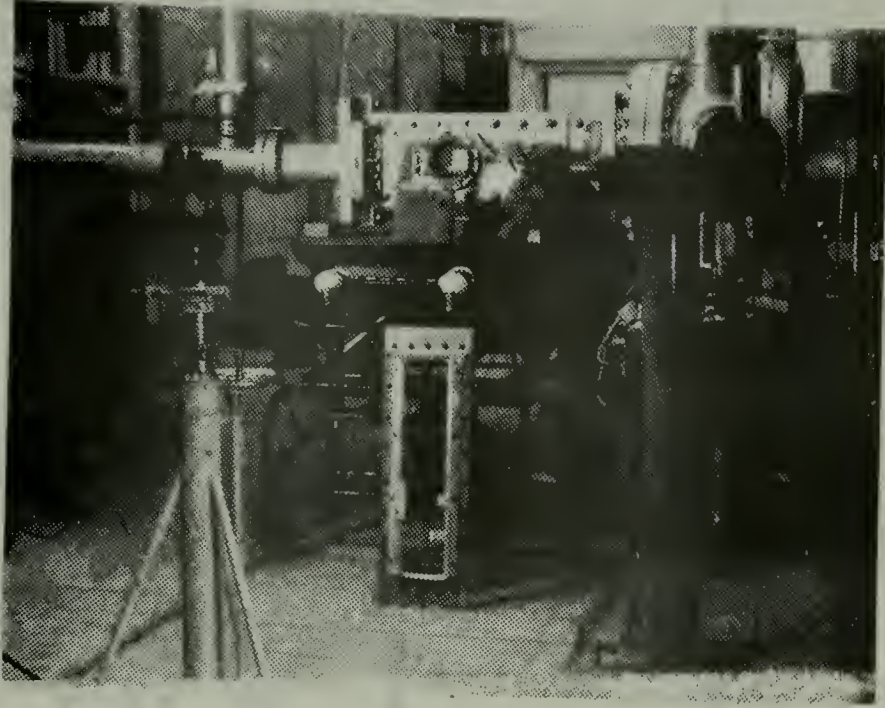


Figure 1(b). Wind Tunnel Test Section with Aluminum Side Walls





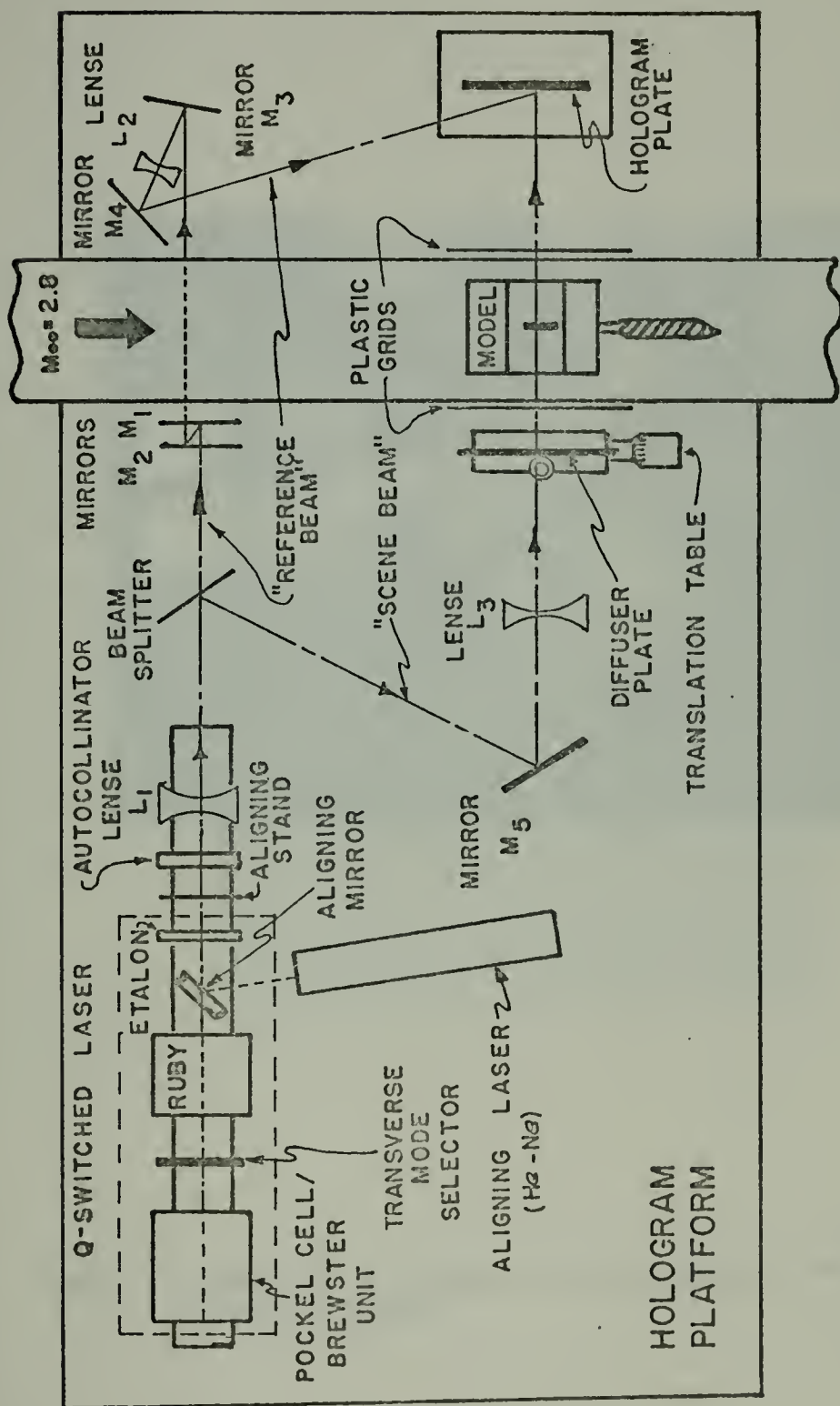


Figure 2. Schematic Representation of the Holographic Arrangement



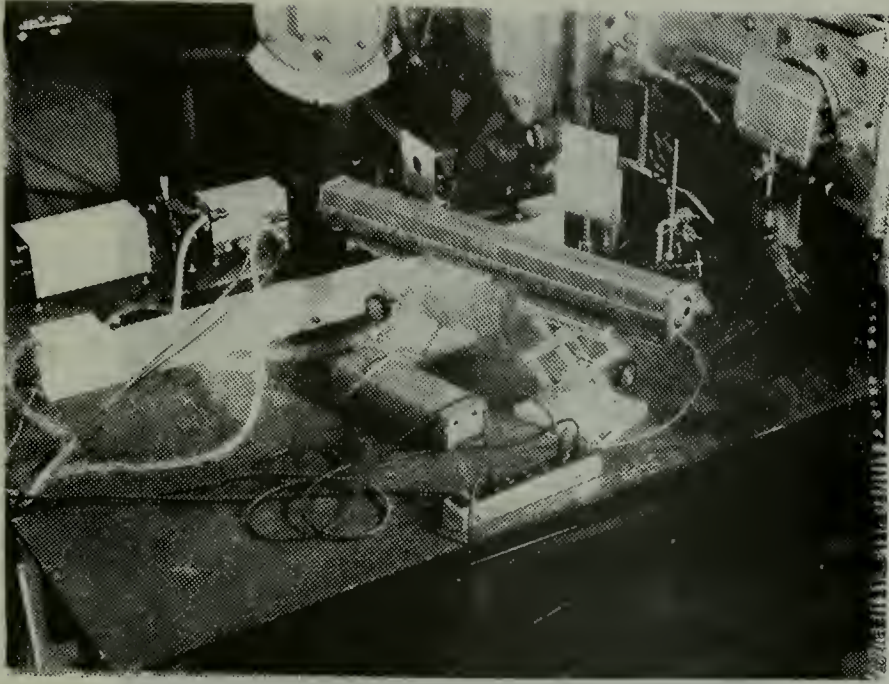


Figure 3. Holographic Arrangement Including the Laser Cooling and Aligning Equipment



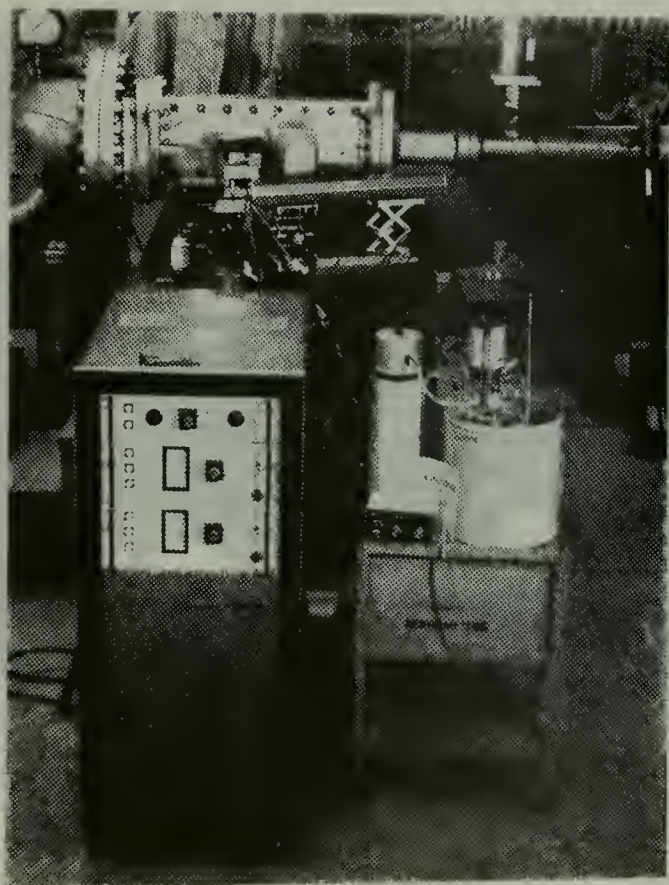


Figure 4. Ruby Laser Power Supply and Cooling Unit





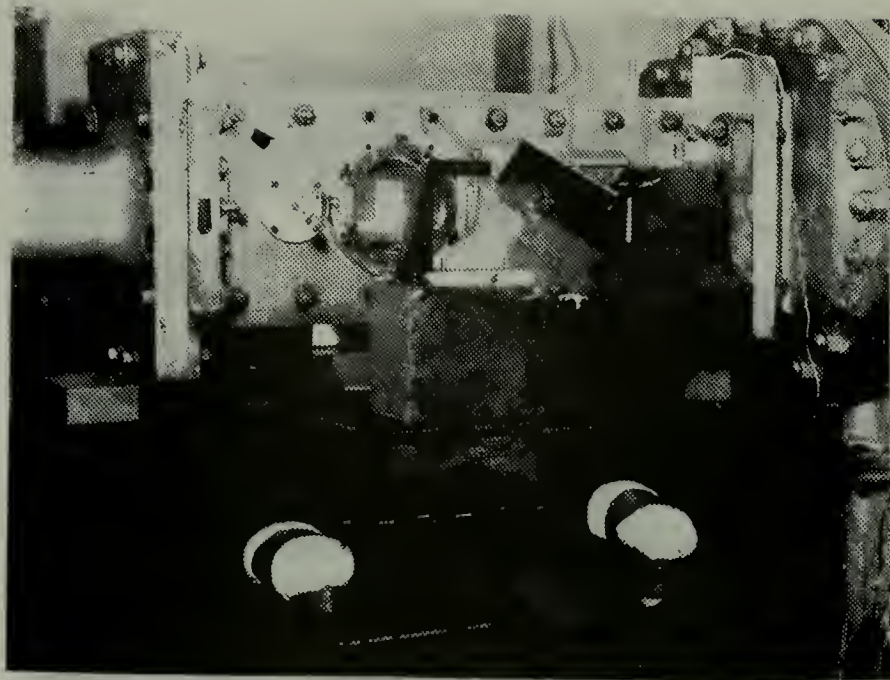


Figure 5. Hologram Plate Holder and Mirrors on the Reverse Side of the Wind Tunnel





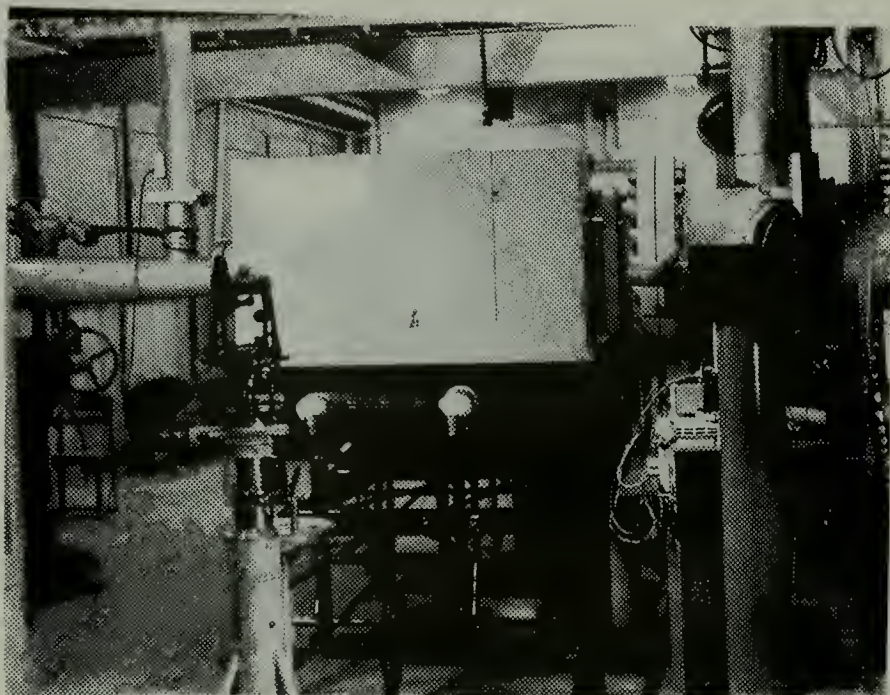


Figure 6. Hologram Platform Box Cover for Daylight Photography



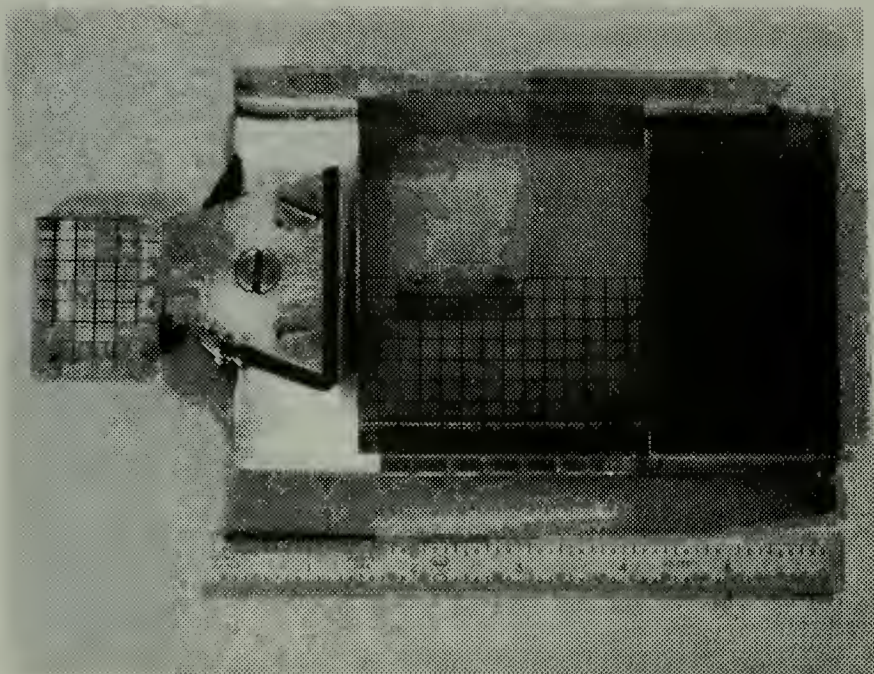


Figure 7(a). Initial Fin-Flat Plate Model Used in the Experiment



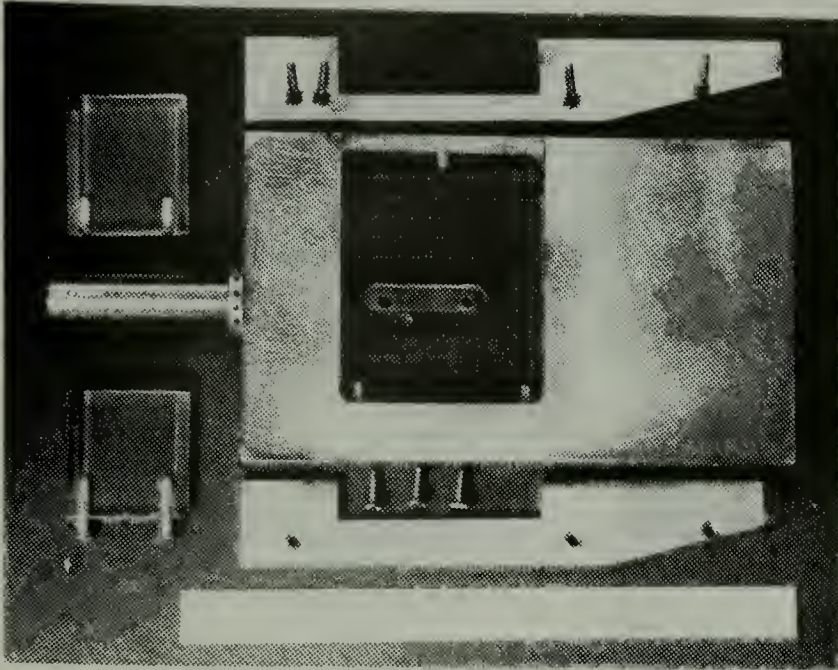
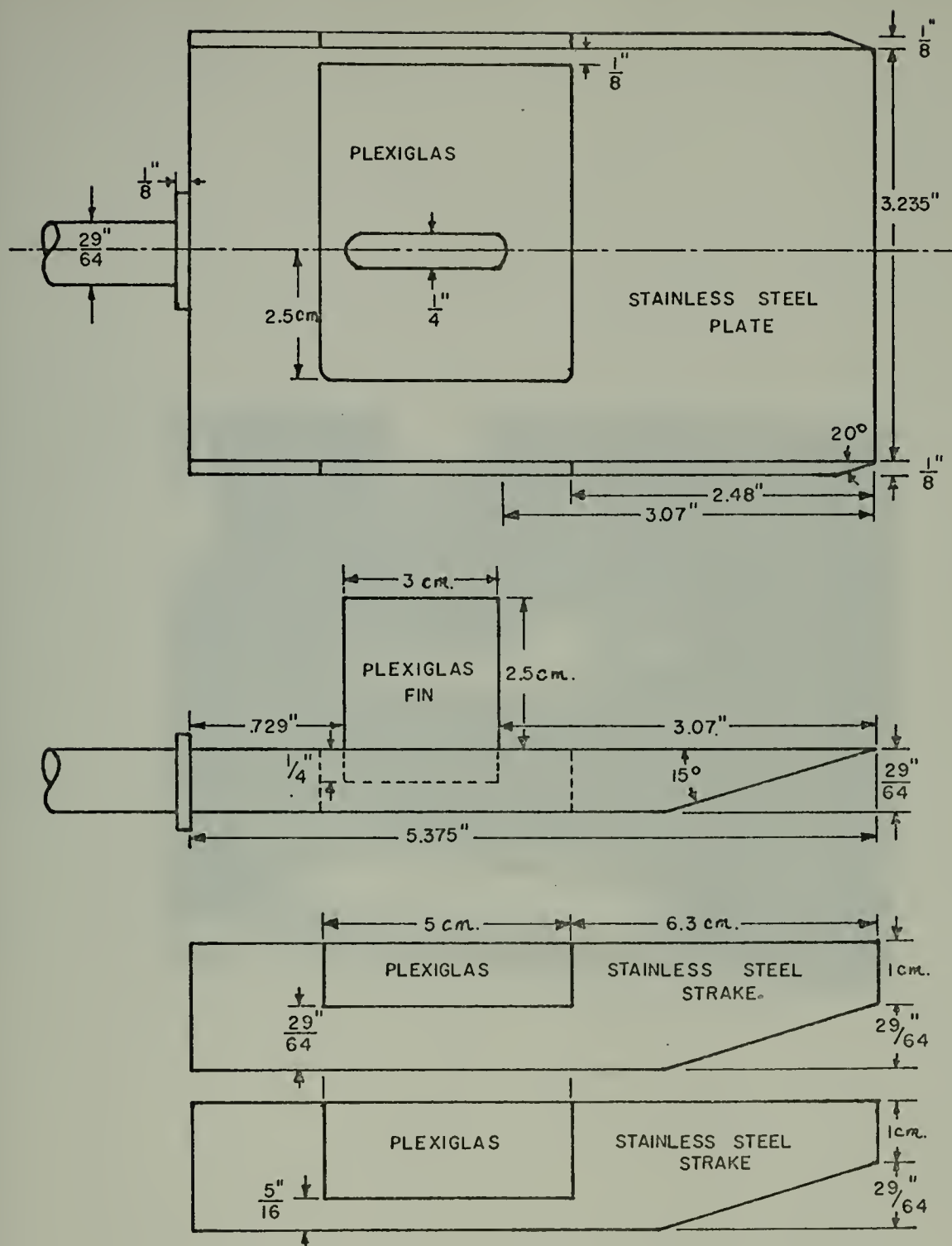


Figure 7(b). Second Fin-Flat Plate Model Used in the Experiment











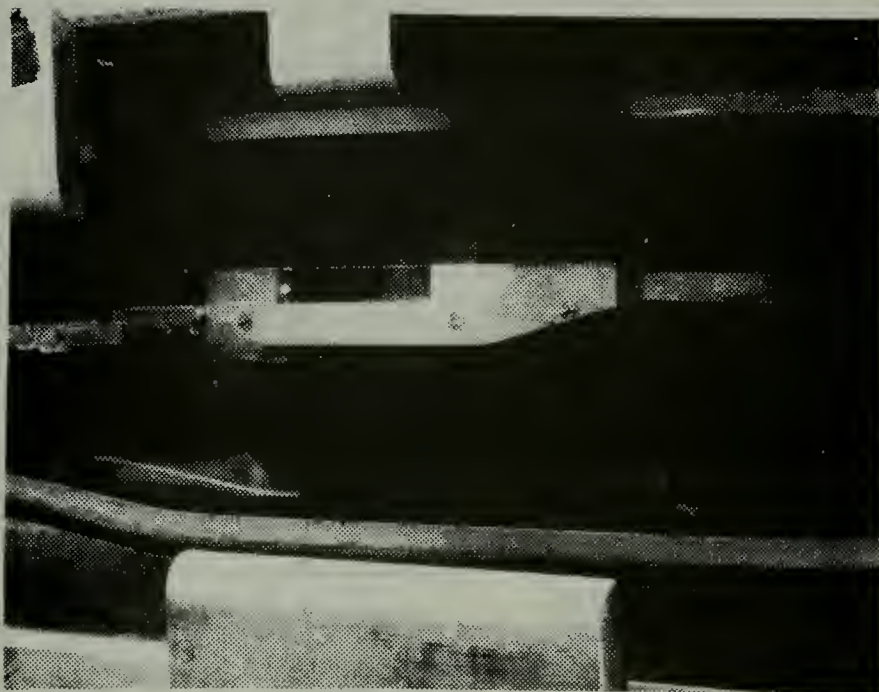


Figure 7(d). Model Mounting in the Wind Tunnel Test Section



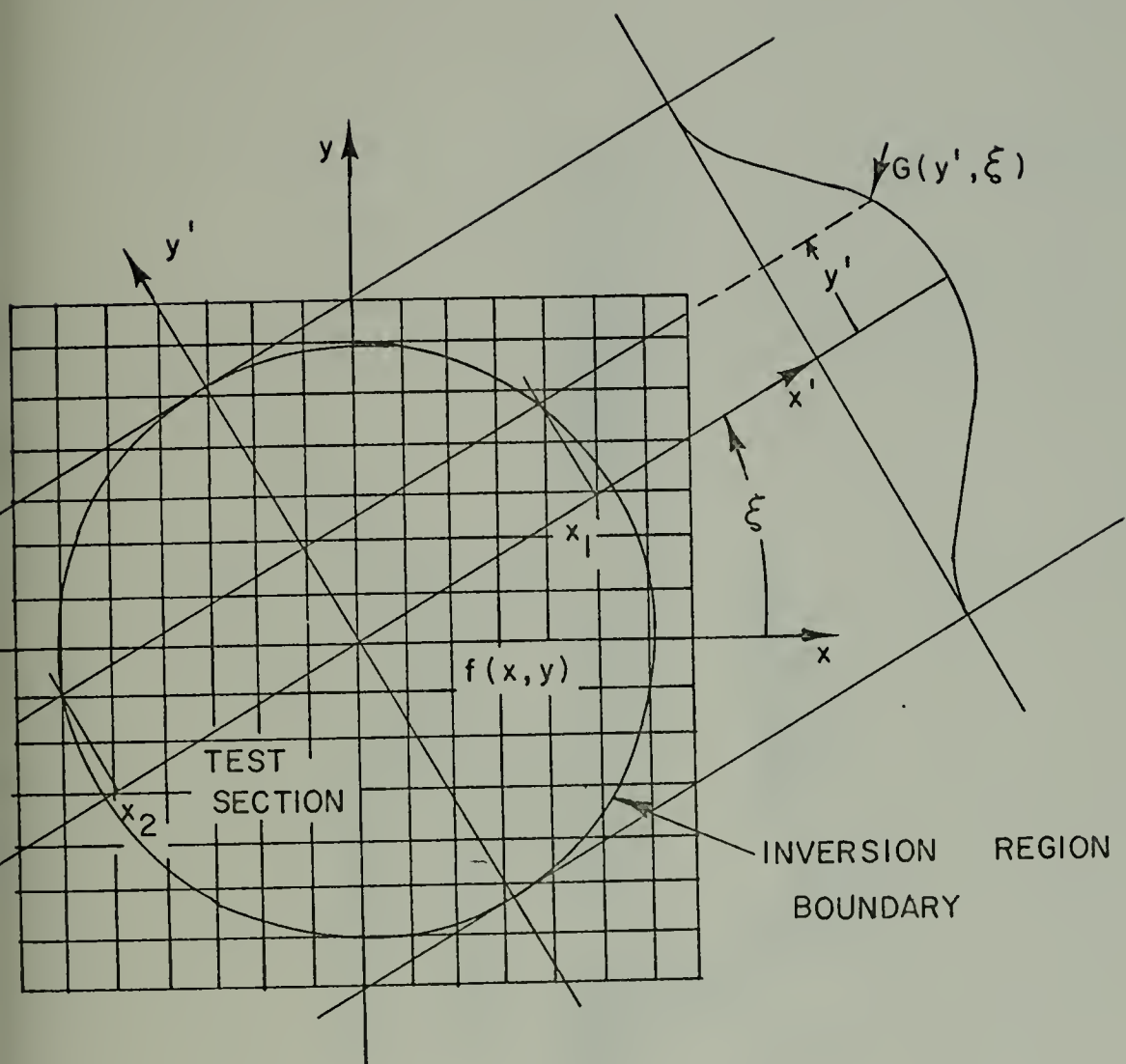


FIGURE 8. CO-ORDINATE SYSTEM USED FOR THE INVERSION OF FRINGE NUMBER TO DENSITY



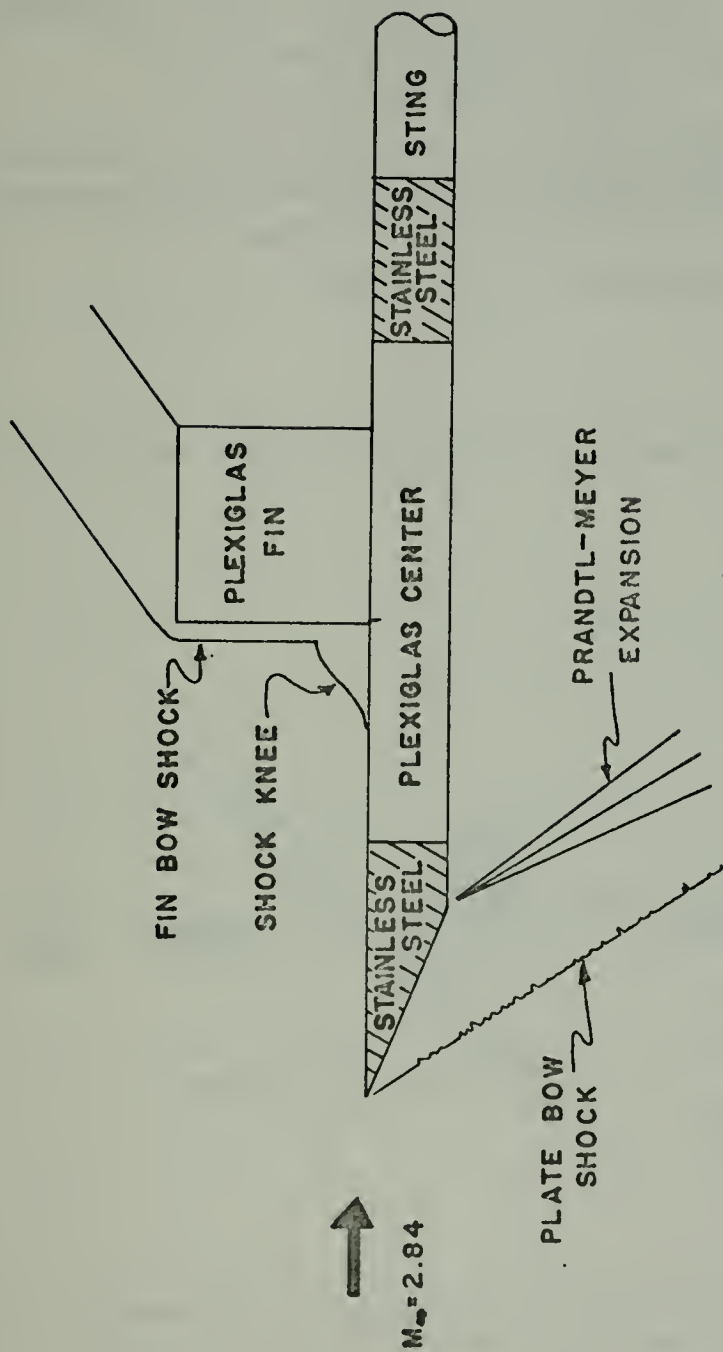


Figure 9. Schematic of the Desired Model Flow Network



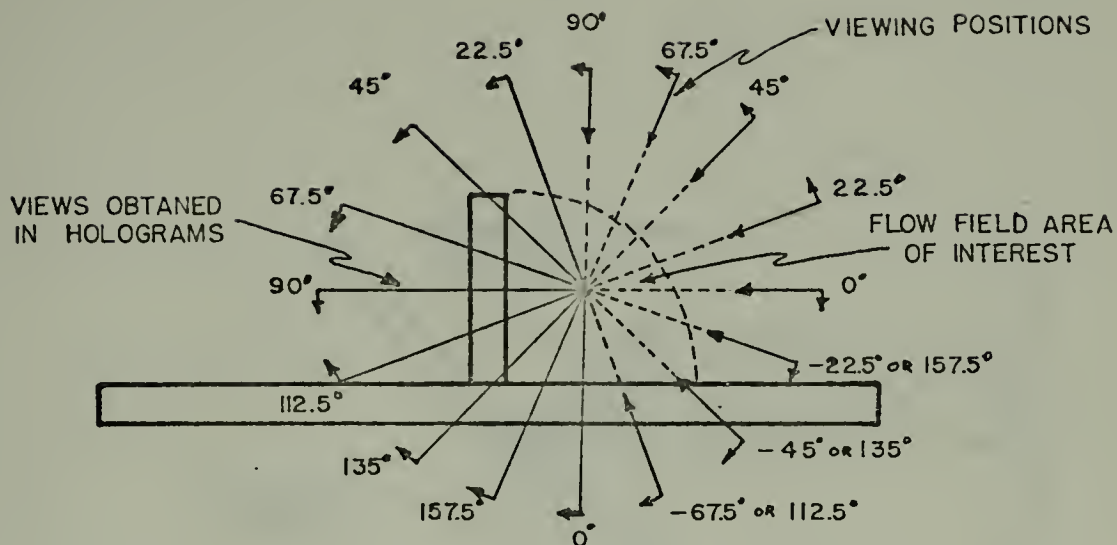


Figure 10. Desired Holographic Views in the Direct Fin-Roof Method

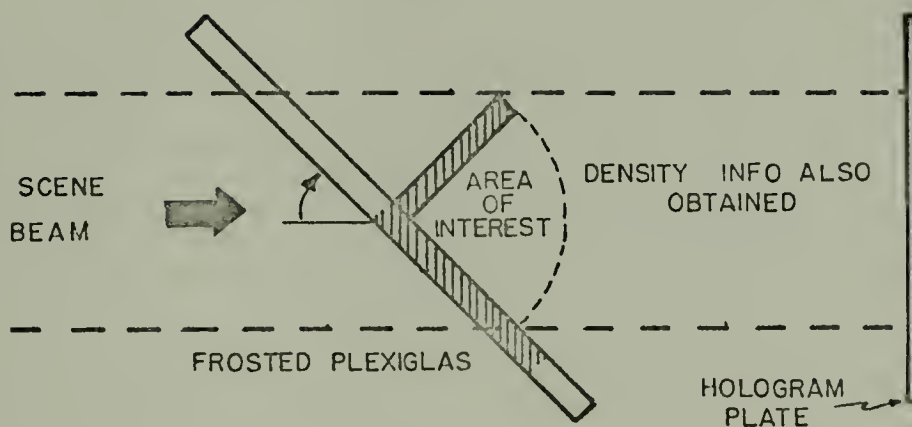


Figure 11. Schematic of the Model Used to Obtain Holographic Views Between 0° and 90° in the Direct Fin-Roof Method





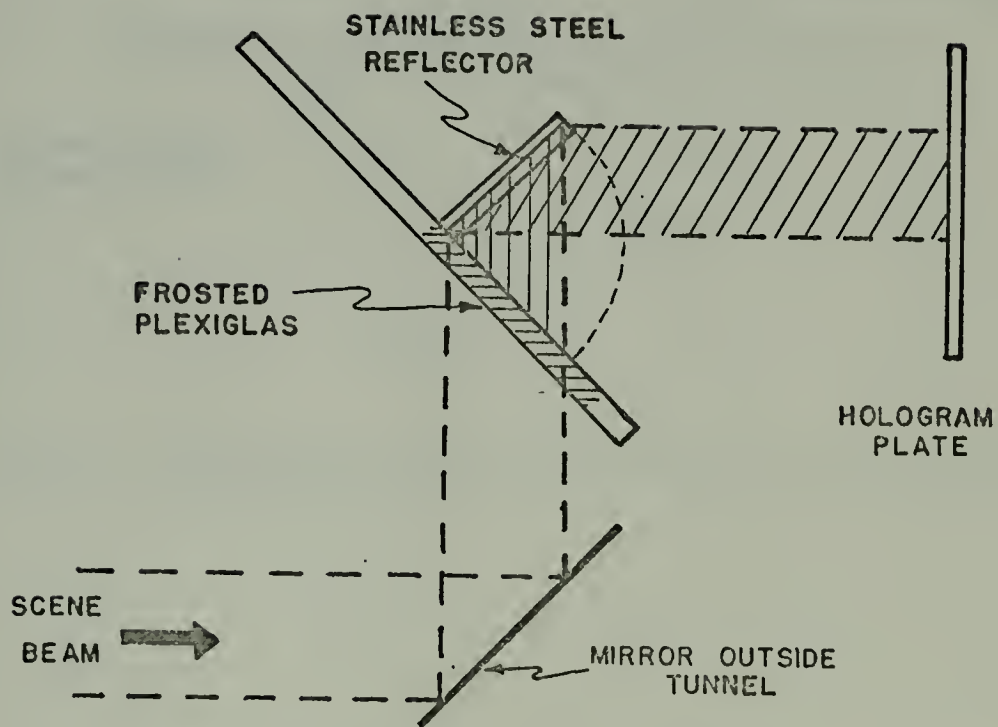


Figure 12. Schematic of the Model Used to Obtain Holographic Views Between  $90^\circ$  and  $180^\circ$  in the Direct Fin-Root Method



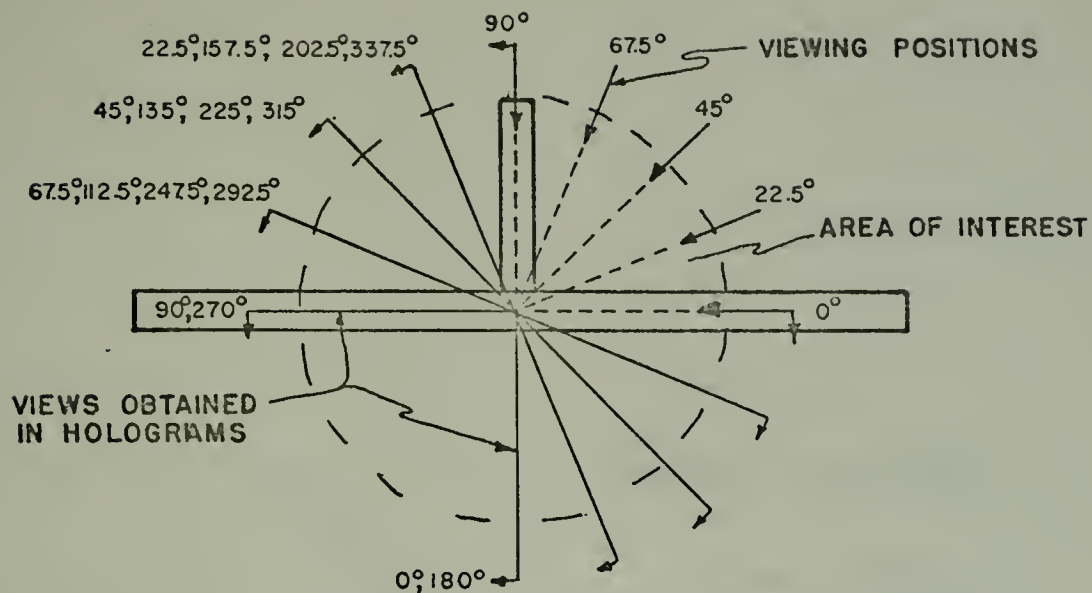


Figure 13. Holographic Viewing Angles Required in the Total Model Flow Method

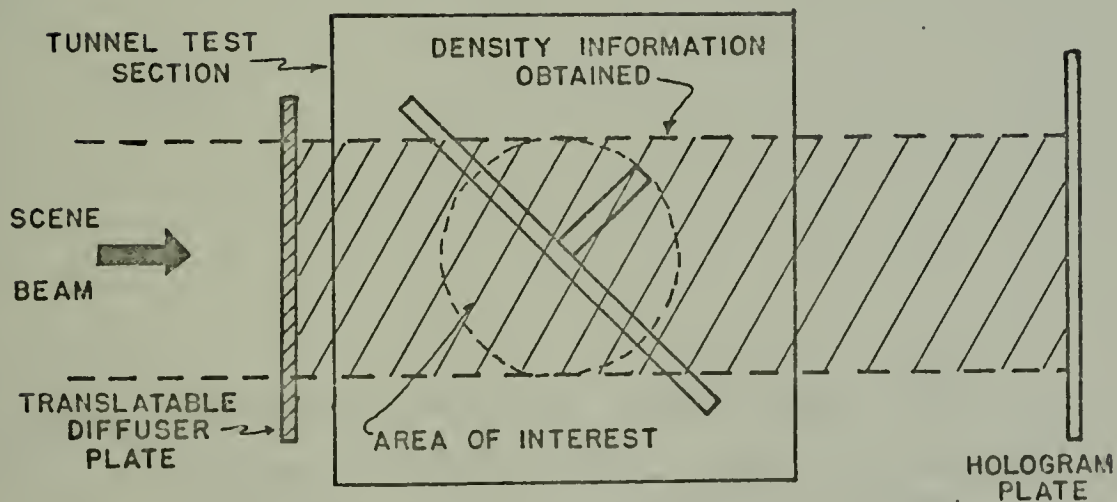


Figure 14. Schematic of the Technique Used to Obtain Holograms in the Total Model Flow Method



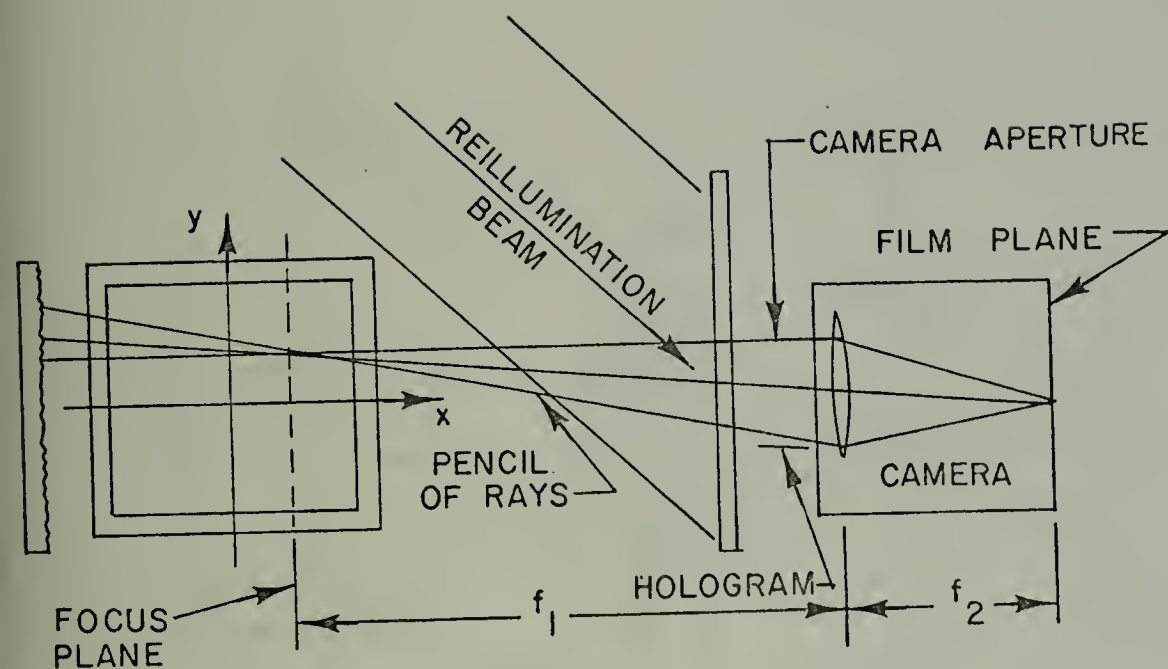


FIGURE 15. EFFECT OF APERTURE SIZE FOCUS PLANE POSITION ON THE PENCIL SIZE OF RAYS ABOUT A LINE OF SIGHT RECORDED BY CAMERA



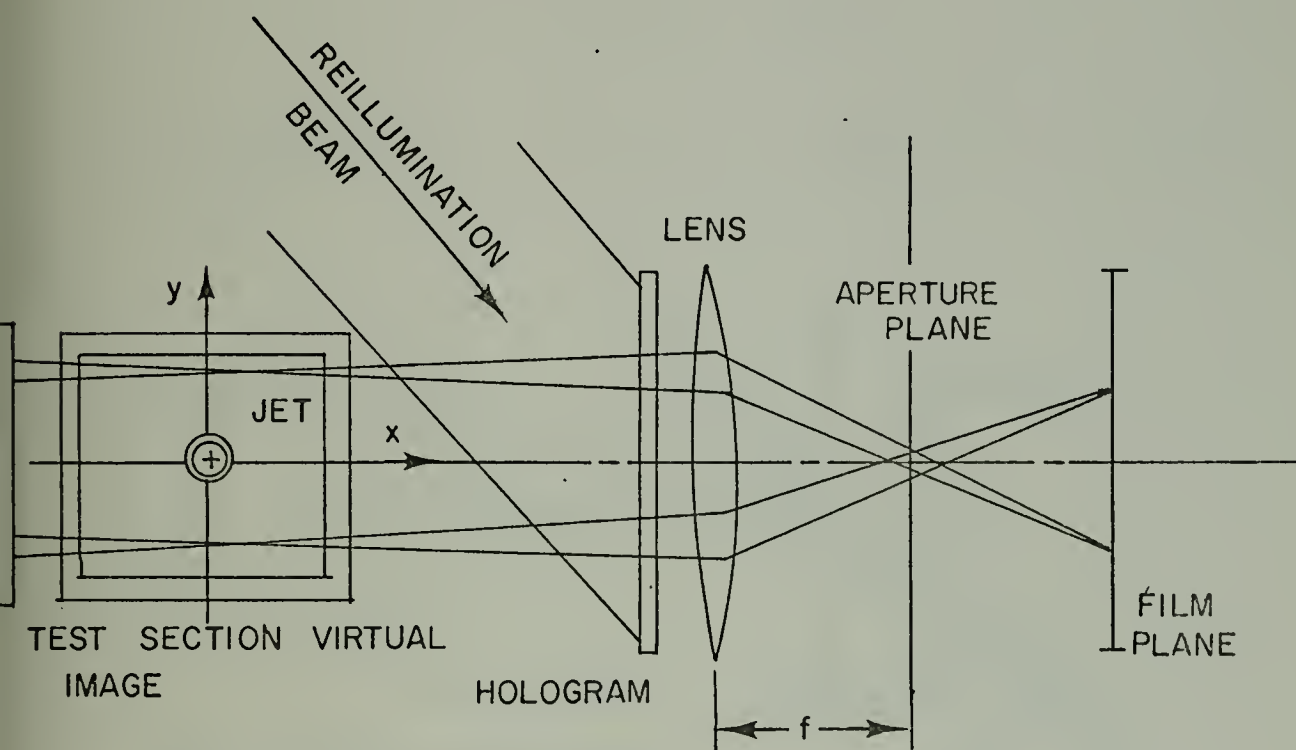


FIGURE 16. SPATIAL FILTERING TECHNIQUE FOR SELECTING PHOTOGRAPH OF CONSTANT ANGLE LINES OF LIGHT





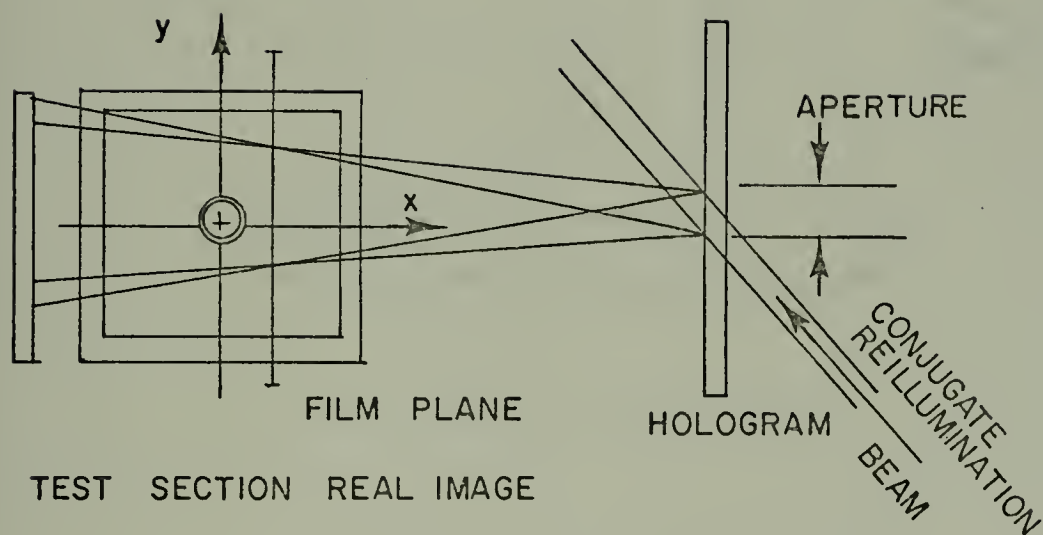


FIGURE 17. LENSLESS PHOTOGRAPHIC TECHNIQUE USING A CONJUGATE REFERENCE BEAM OF SMALL DIAMETER



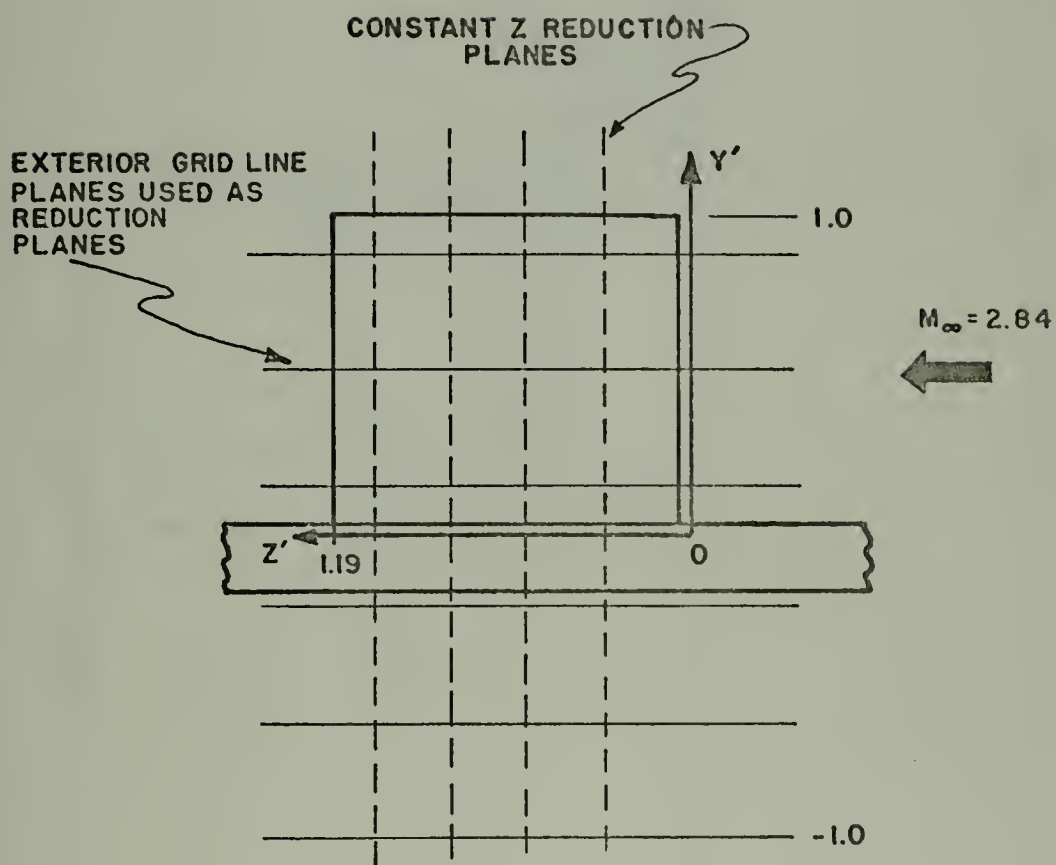


Figure 18. Data Reduction Planes Desired to Describe the Density Field Down the Fin



# PLASTIC GRID AND TUNNEL WALL

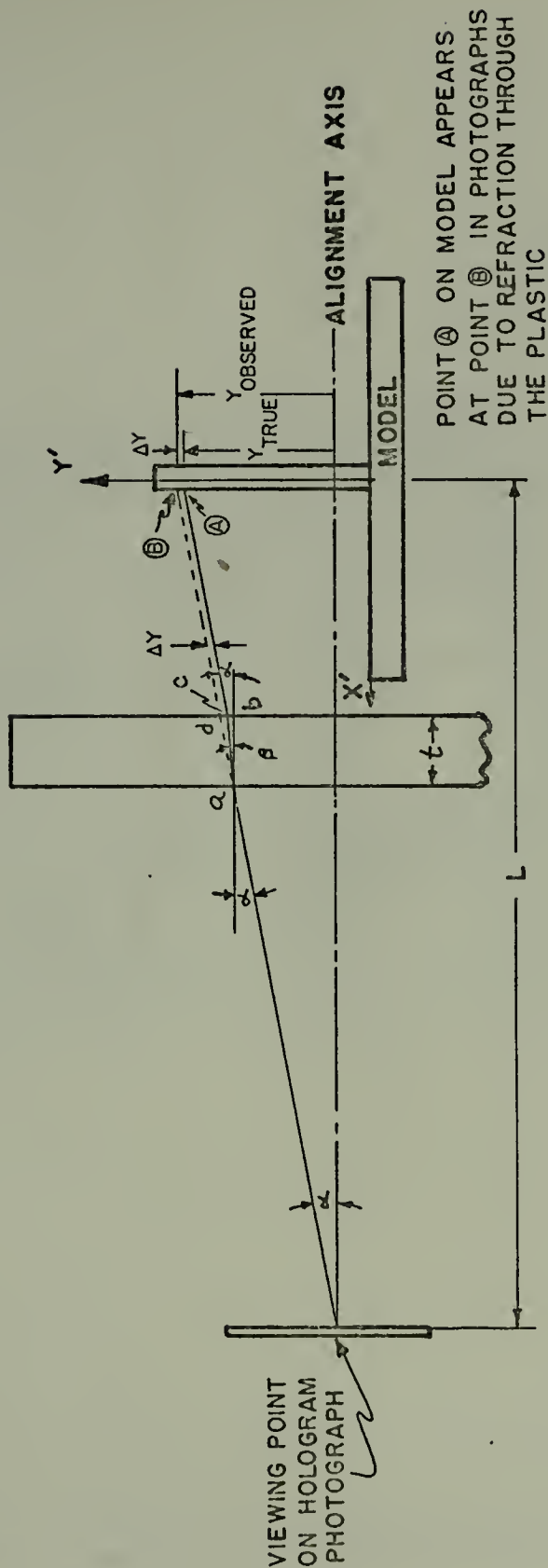


Figure 19. Schematic of the Refraction Displacement in the Photographs of Points Not Located on the Aligned Axis Caused by the Plastic Grid and Tunnel Wall



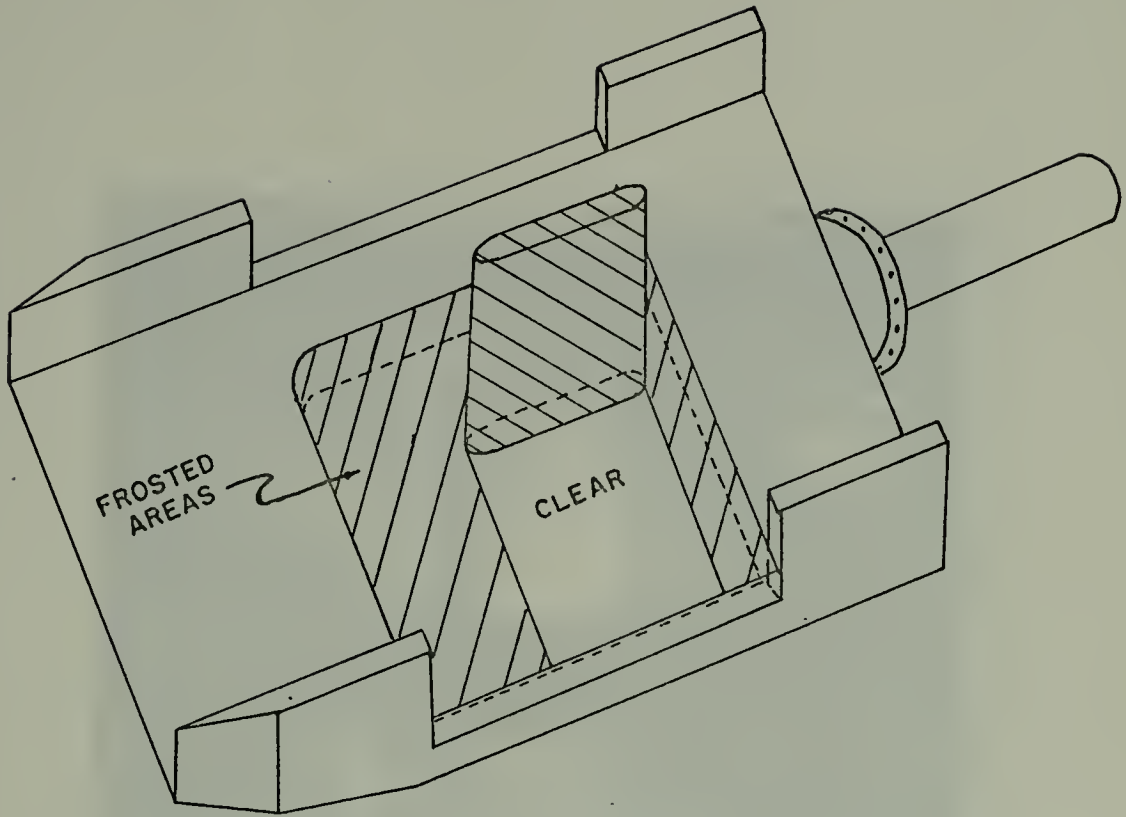


Figure 20. Schematic of the Diffused Plastic Portion of the Model  
Required in the Direct Fin Flow Method





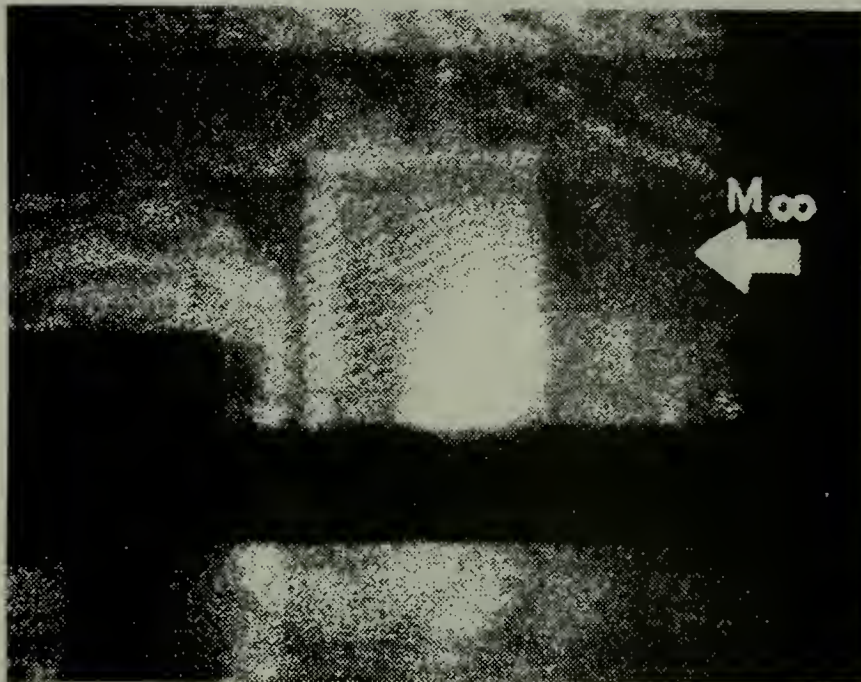


Figure 21. Interferogram Obtained in the Direct Fin Flow Method with the Model at  $0^\circ$  Rotation, Mach 2.84 and no Translation of Mirror,  $M_5$



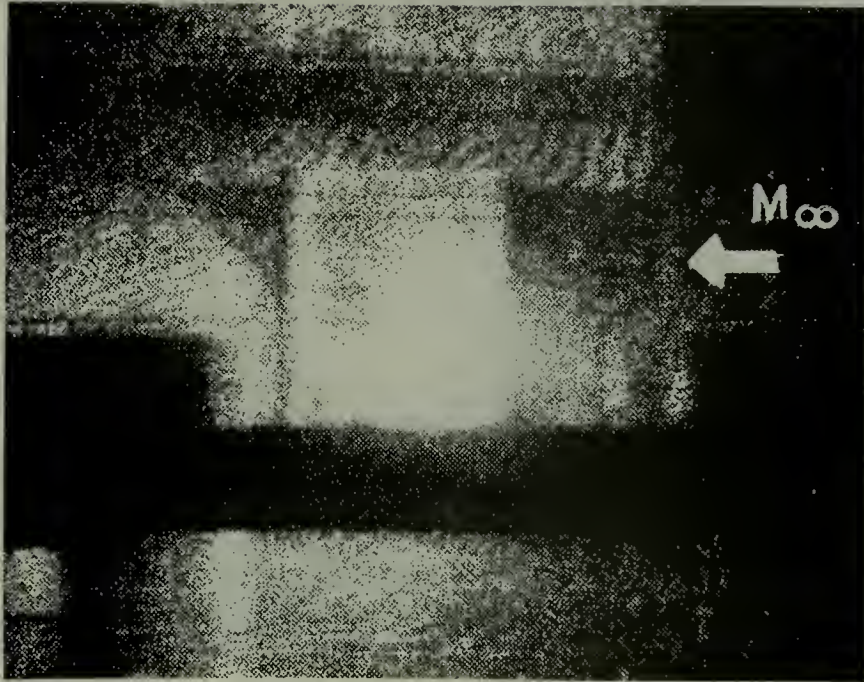


Figure 22. Interferogram Obtained in the Direct Fin Flow Method with the Model at  $0^\circ$  Rotation, Mach 2.84 and a .006 inches Translation of Mirror,  $M_5$ , Parallel to the Test Section



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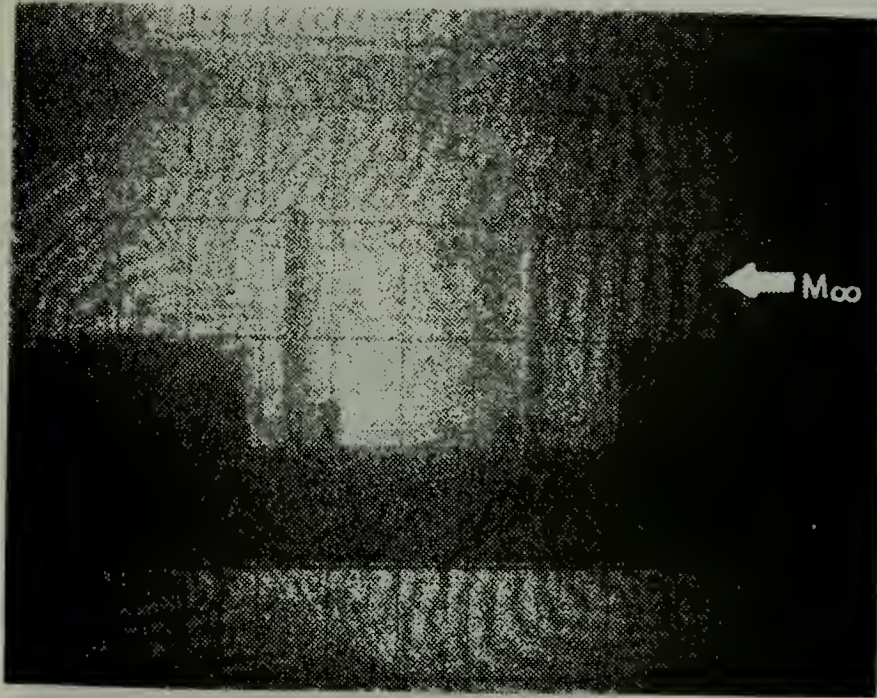


Figure 23. Interferogram Obtained in the Total Flow Method with the Clear Plexiglas Fin Model at  $0^\circ$  Rotation, Mach 2.84 and a .0015 inches Horizontal Translation of the Diffuser Plate Parallel to the Test Section



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Figure 24. Interferogram Obtained in the Total Flow Method with the Clear Plexiglas Fin Model at  $0^\circ$  Rotation, Mach 2.84 and a .0045 Vertical Translation of the Diffuser Plate Parallel to the Test Section







Figure 25(a). Illustration of the Shock Network Around the Fin at Mach 2.84 with a  $90^\circ$  Model Rotation Angle Using a Horizontal Knife Edge Schlieren System



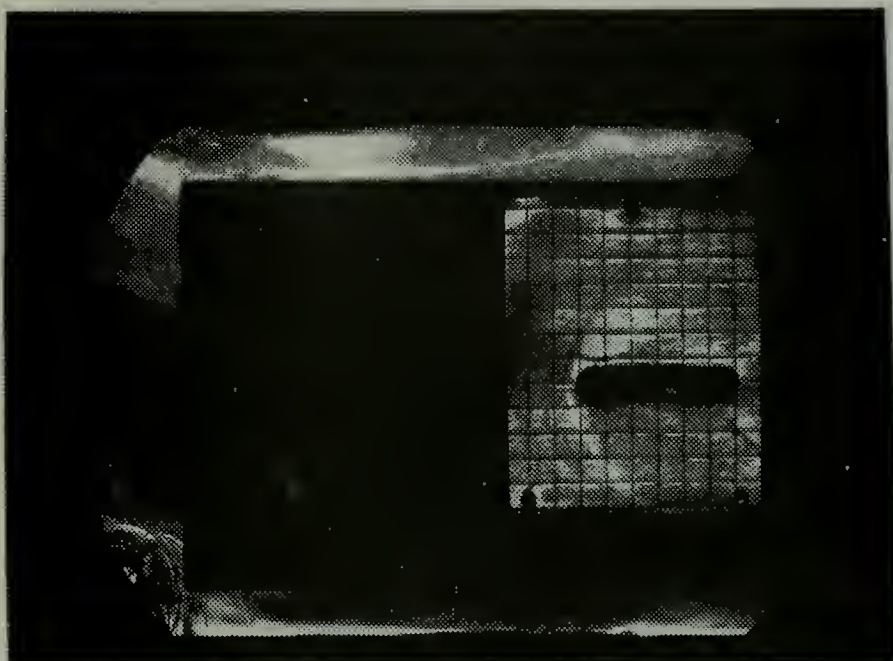


Figure 25(b). Illustration of the Shock Network Around the Fin at Mach 2.84 with a  $90^\circ$  Model Rotation Angle Using a Vertical Knife Edge Schlieren System



Figure 2.10. A photograph of a small, dark, rectangular object, possibly a piece of wood or a small box, resting on a light-colored surface. The object is oriented horizontally and has a slightly irregular shape. The background is a plain, light-colored surface.

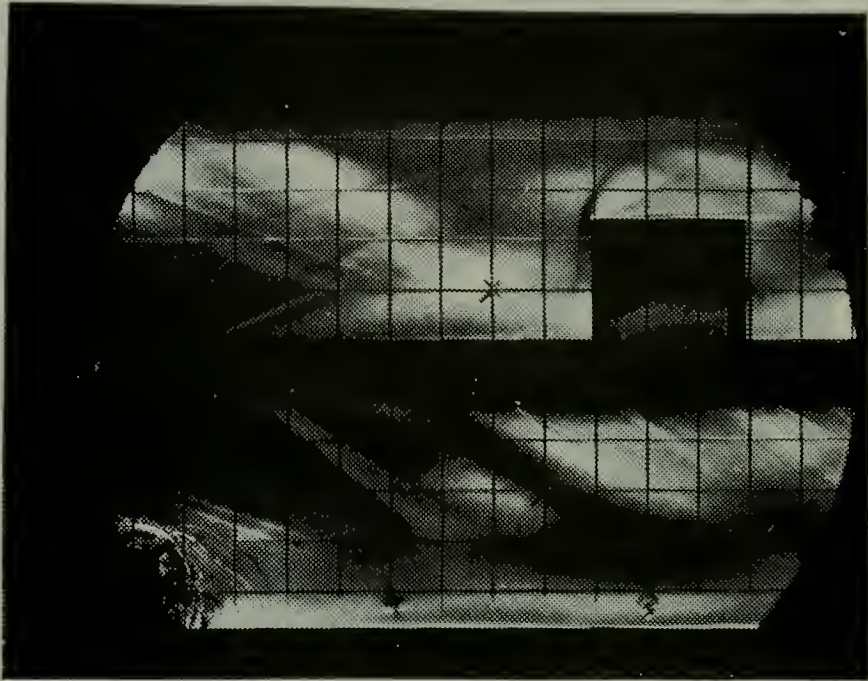


Figure 26. Illustration of the Model Flow Network at Mach 2.84,  $0^\circ$  Model Rotation and a Plate Vibration of  $\pm .05^\circ$  Angle of Attack Using a Schlieren System



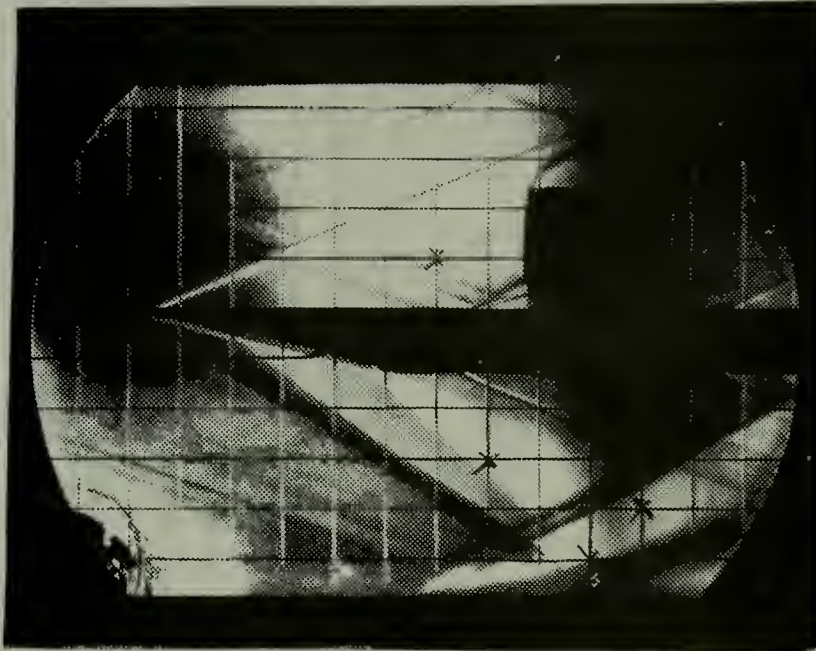


Figure 27. Illustration of the Model Flow Network at Mach 2.84, 0° Model Rotation and No Plate Vibration Using a Schlieren System





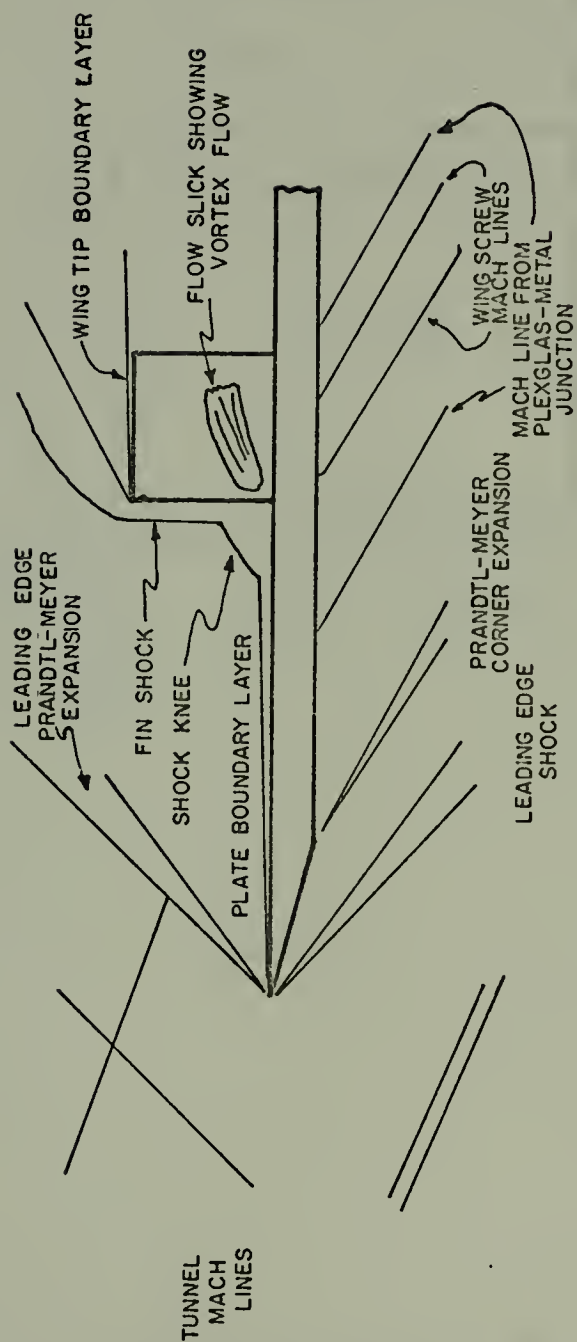


Figure 28. Schematic of the Flow Observed in the Schlieren Photographs



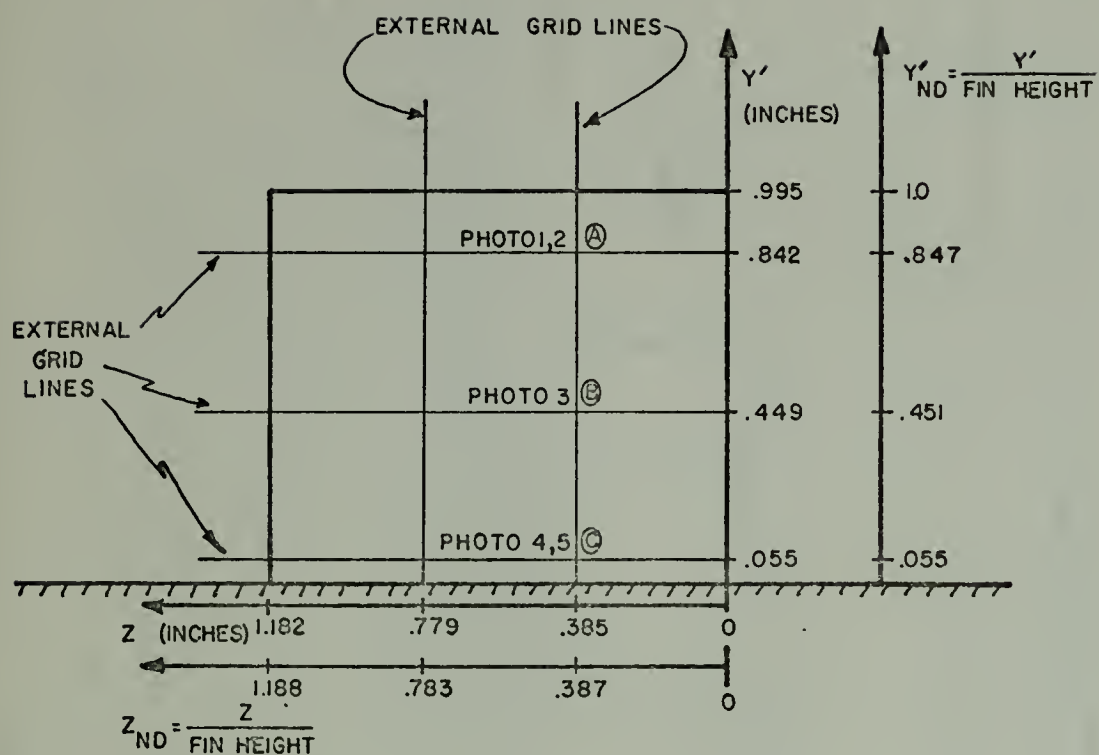


Figure 29. Schematic of the Fin at 0° Rotation Showing the Locations of the External Grid Lines and of the Interferogram Photographic Points



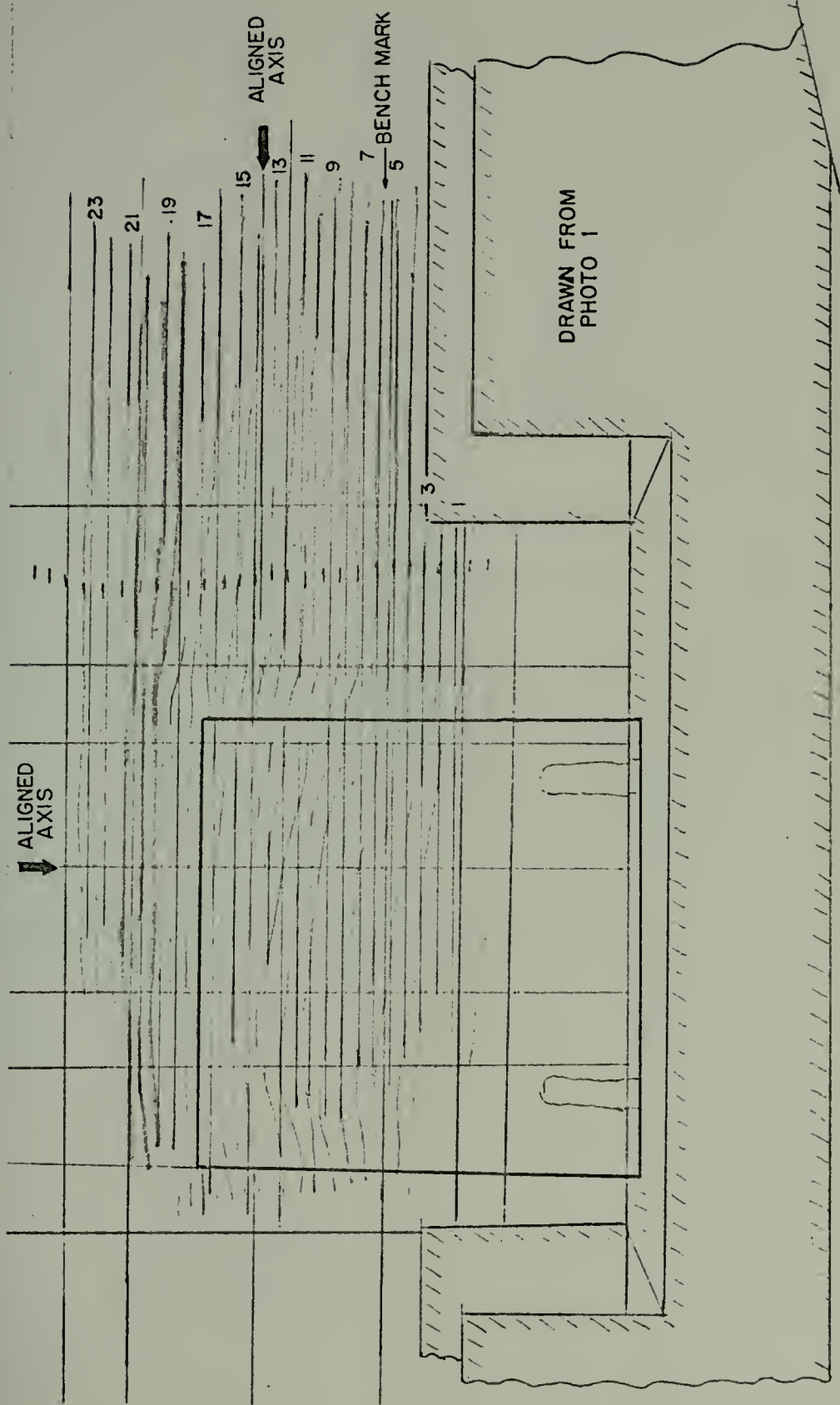


Figure 30. Actual Drawing and Data Reduction of Interferogram Photograph 1 Aligned at  $Z = 0.387$ ,  $Y' = 0.847$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle



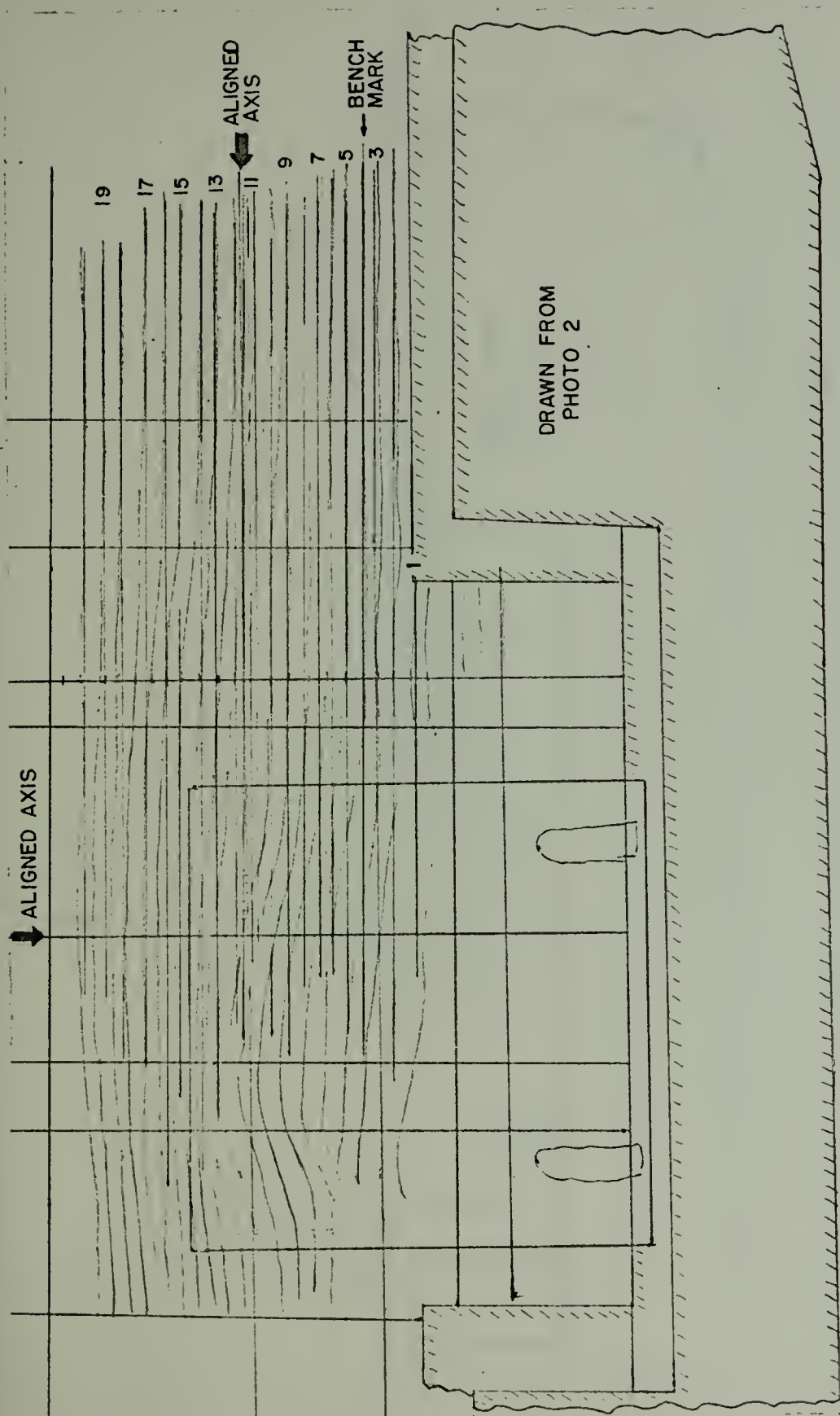


Figure 31. Actual Drawing and Data Reduction of Interferogram Photograph 2 Aligned at  $Z = 0.387$ ,  $Y' = 0.847$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle





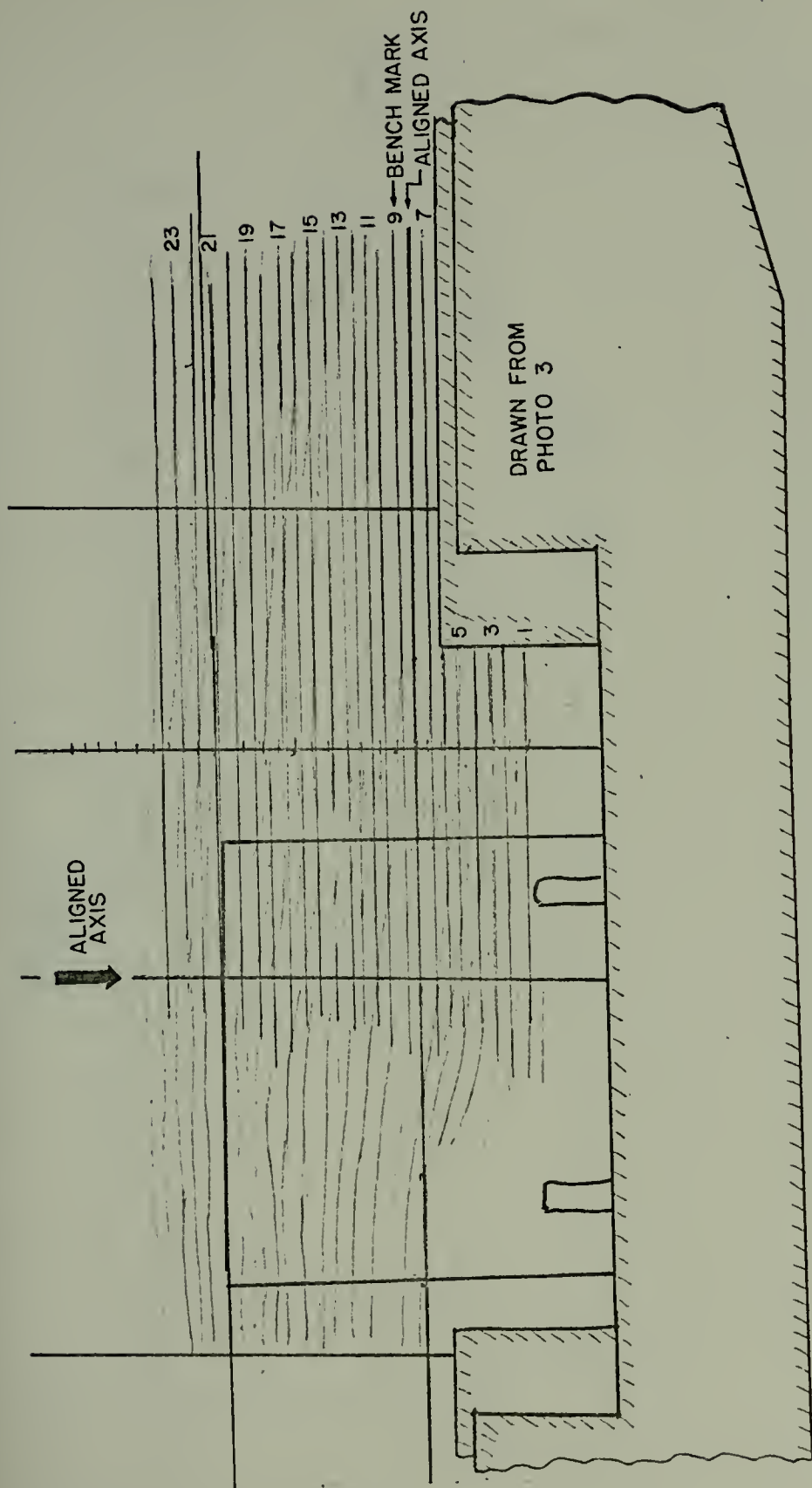


Figure 32. Actual Drawing and Data Reduction of Interferogram Photograph 3 Aligned at  $Z = 0.381$ ,  $y' = 0.451$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle



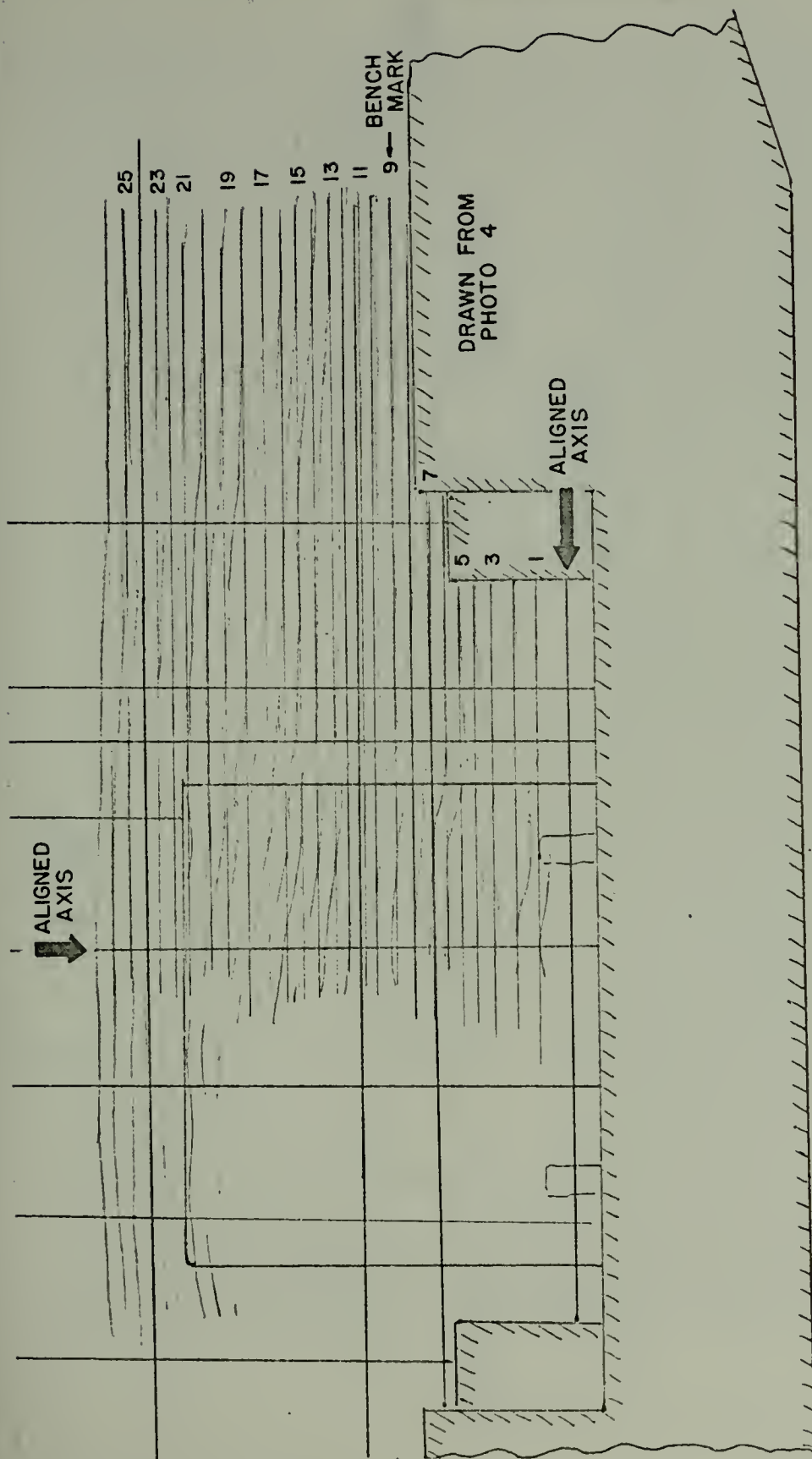


Figure 33. Actual Drawing and Data Reduction of Interferogram Photograph 4 Aligned at  $Z = 0.387$ ,  $y' = 0.055$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle



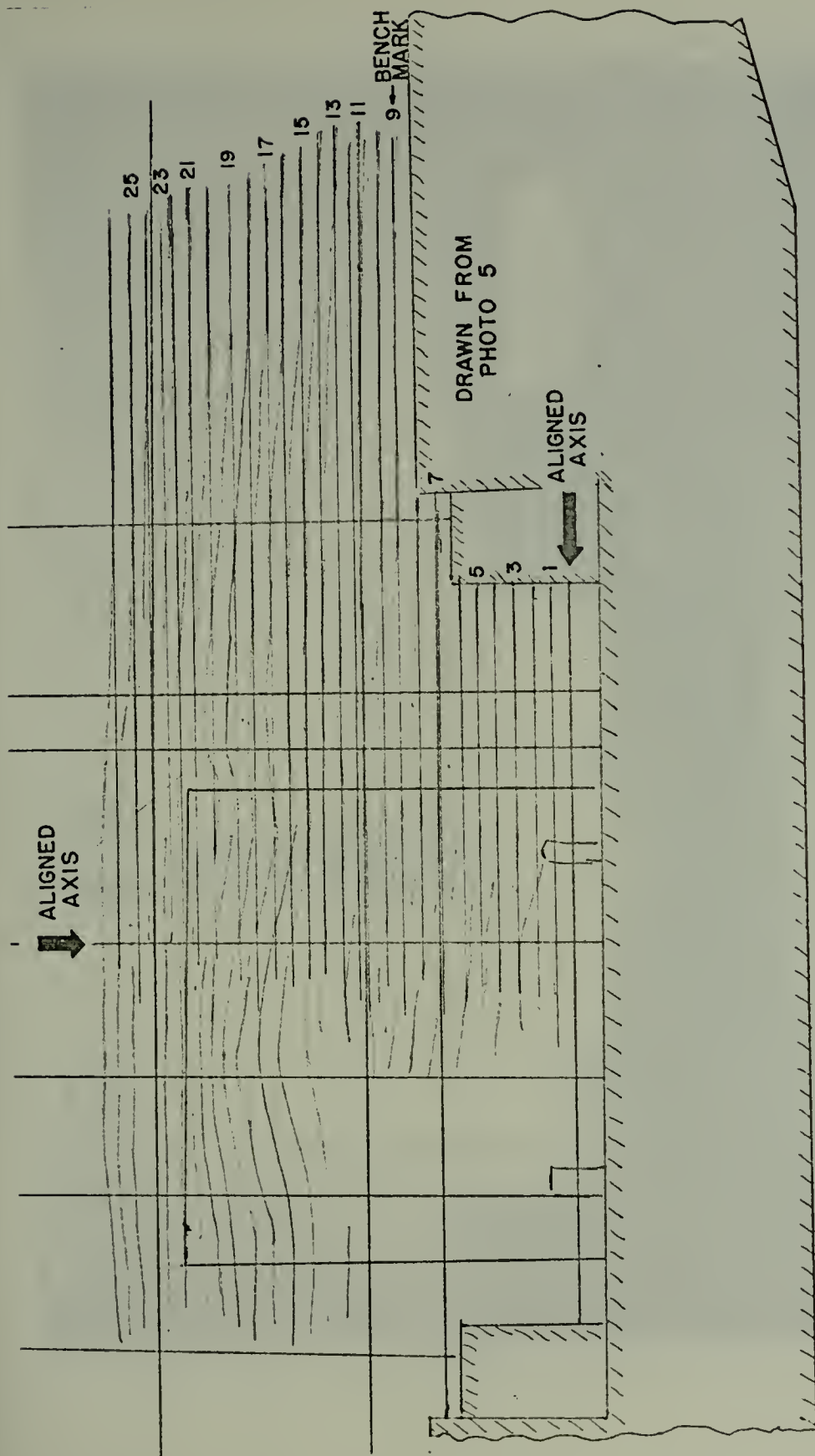


Figure 34. Actual Drawing and Data Reduction of Interferogram Photograph 5 Aligned at  $Z = 0.387$ ,  $y' = 0.055$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle





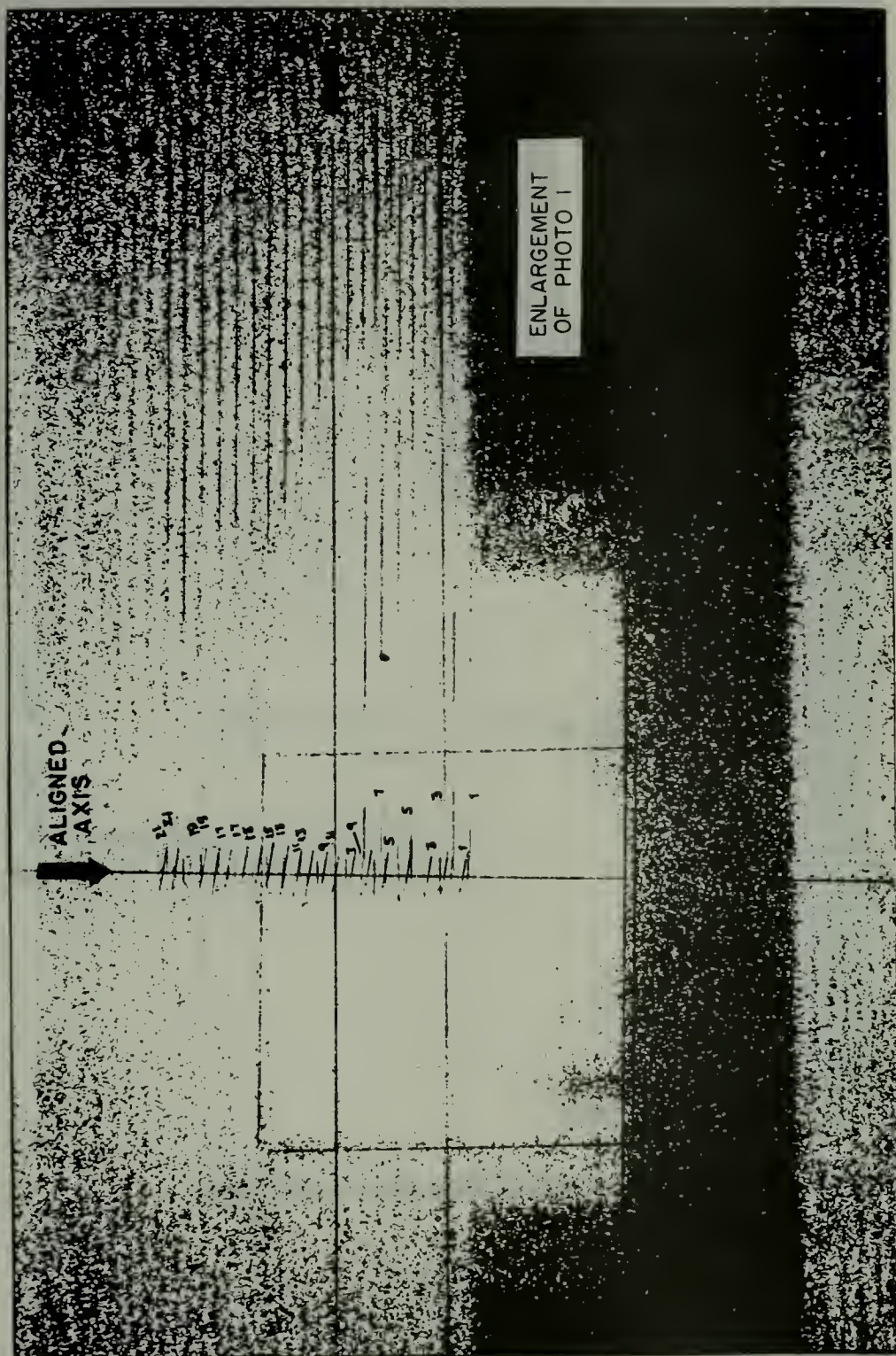


Figure 35. Actual Enlarged Photograph Used for the Data Reduction of Interferogram Photograph 1  
 Aligned at  $Z = 0.387$ ,  $Y' = 0.847$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle





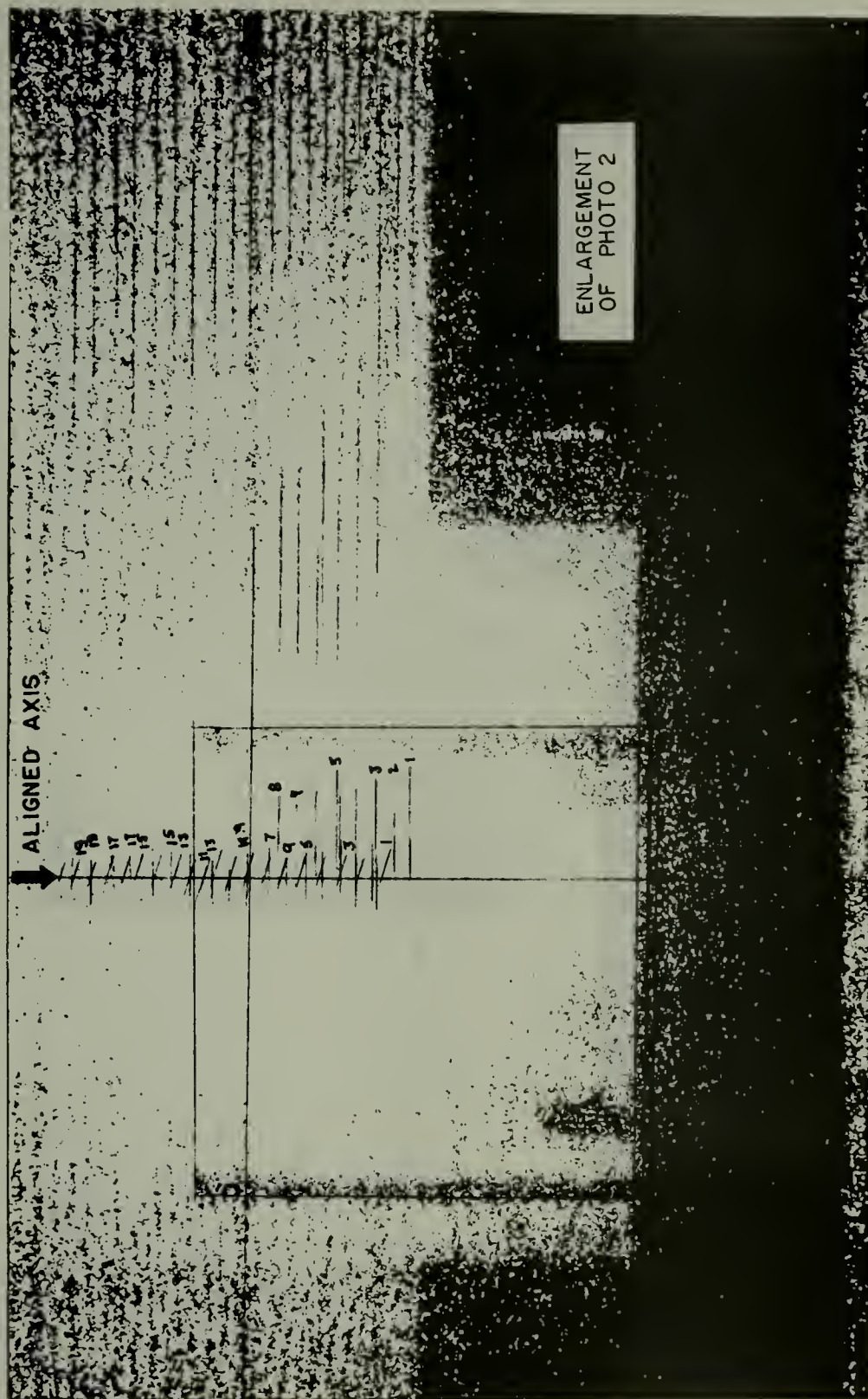


Figure 36. Actual Enlarged Photograph Used for the Data Reduction of Interferogram Photograph 2  
 Aligned at  $Z = 0.387$ ,  $Y' = 0.847$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle



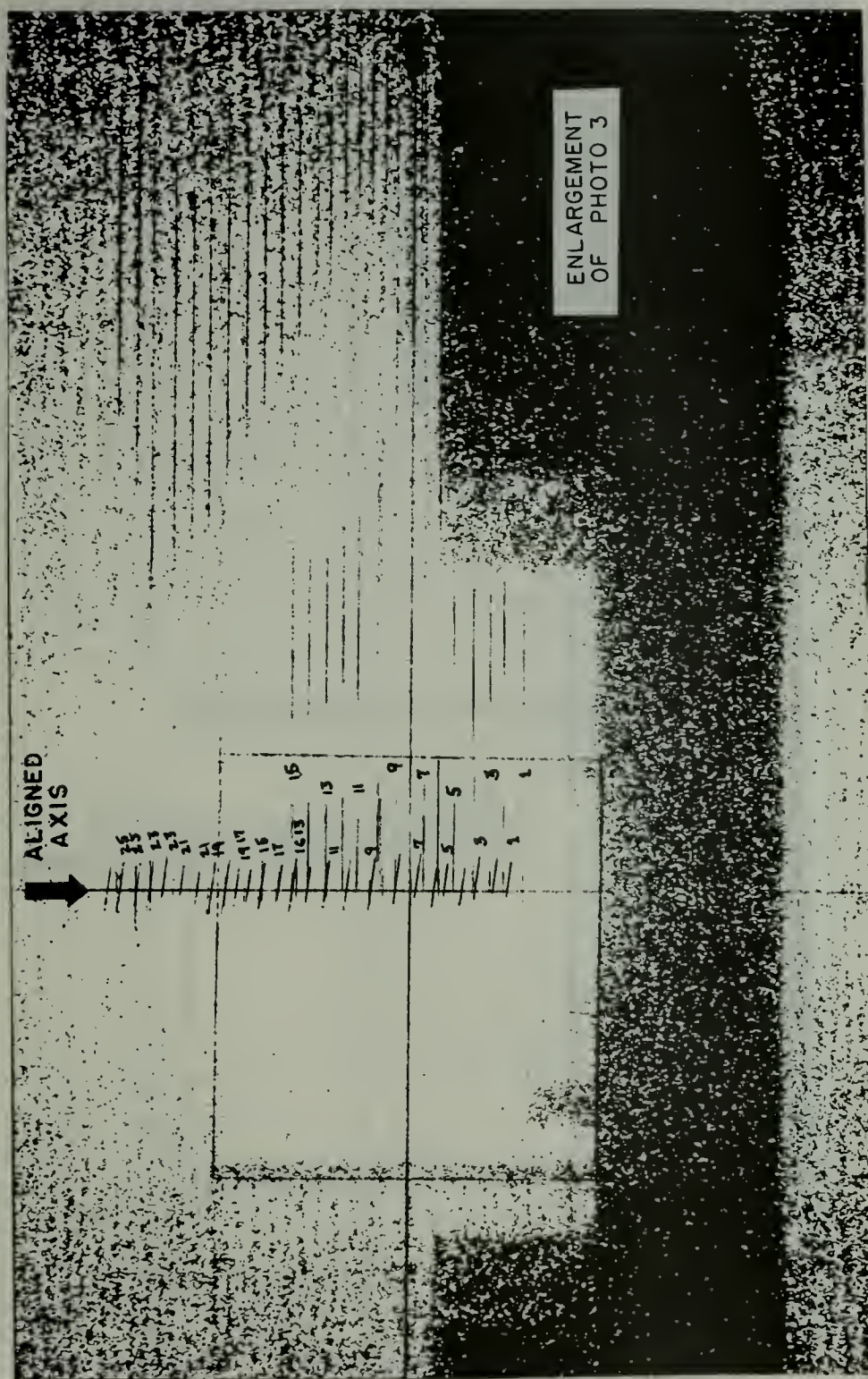
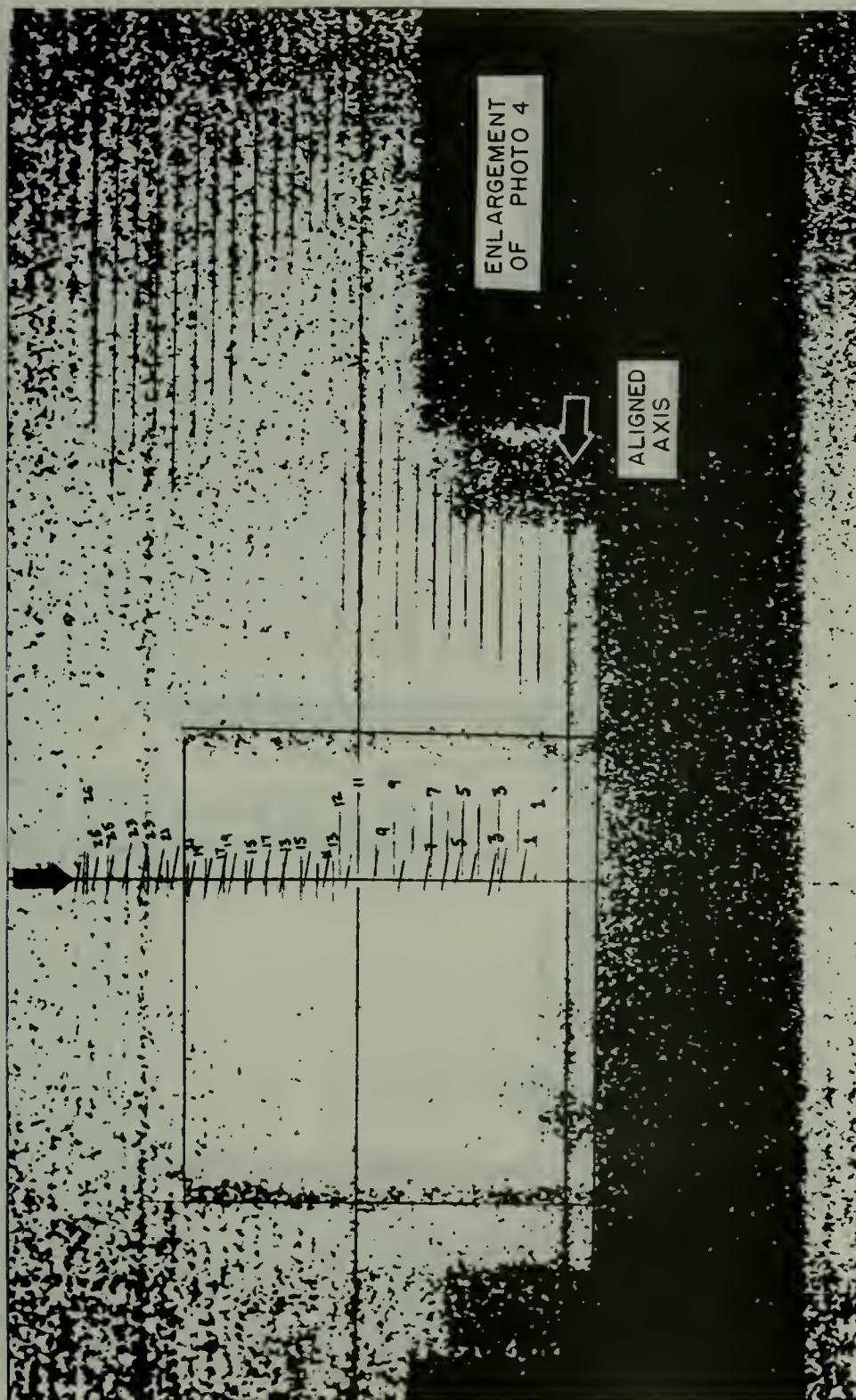


Figure 37. Actual Enlarged Photograph Used for the Data Reduction of Interferogram Photograph 3  
 Aligned at  $Z = 0.387$ ,  $Y' = 0.451$  for Mach 2.84 and 0 Model Potation Angle









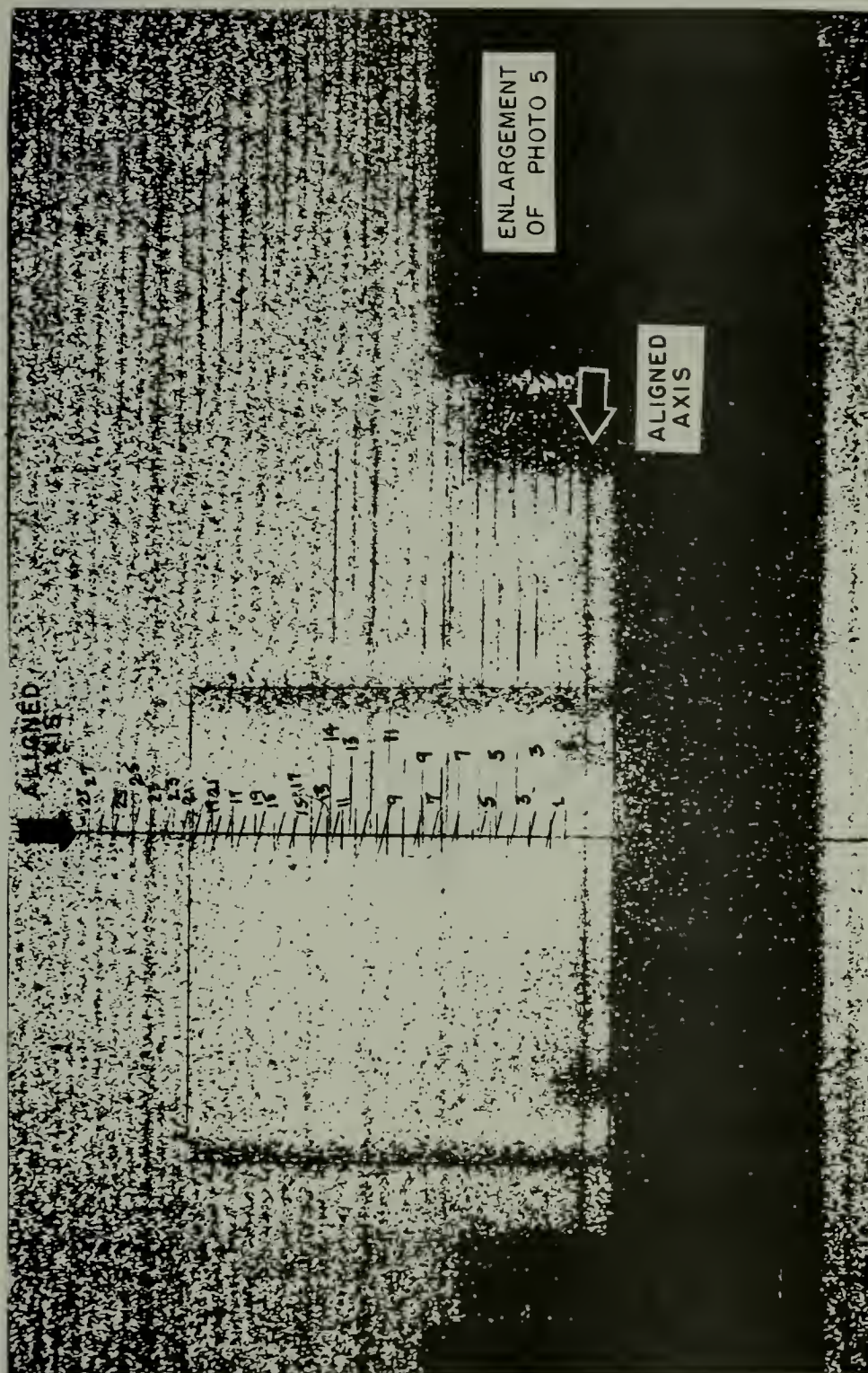


Figure 39. Actual Enlarged Photograph Used for the Data Reduction of Interferogram Photograph. 5  
 Aligned at  $Z = 0.387$ ,  $Y' = 0.055$  for Mach 2.84 and  $0^\circ$  Model Rotation Angle







Figure 40. Comparison of Fringe Data Obtained from Drawings of Interferogram Photographs 1 and 2 Aligned at  $Y' = 0.847$





Figure 41. Comparison of Fringe Data Obtained from Photographic Enlargements of Interferogram Photographs 1 and 2 Aligned at  $Y' = 0.847$



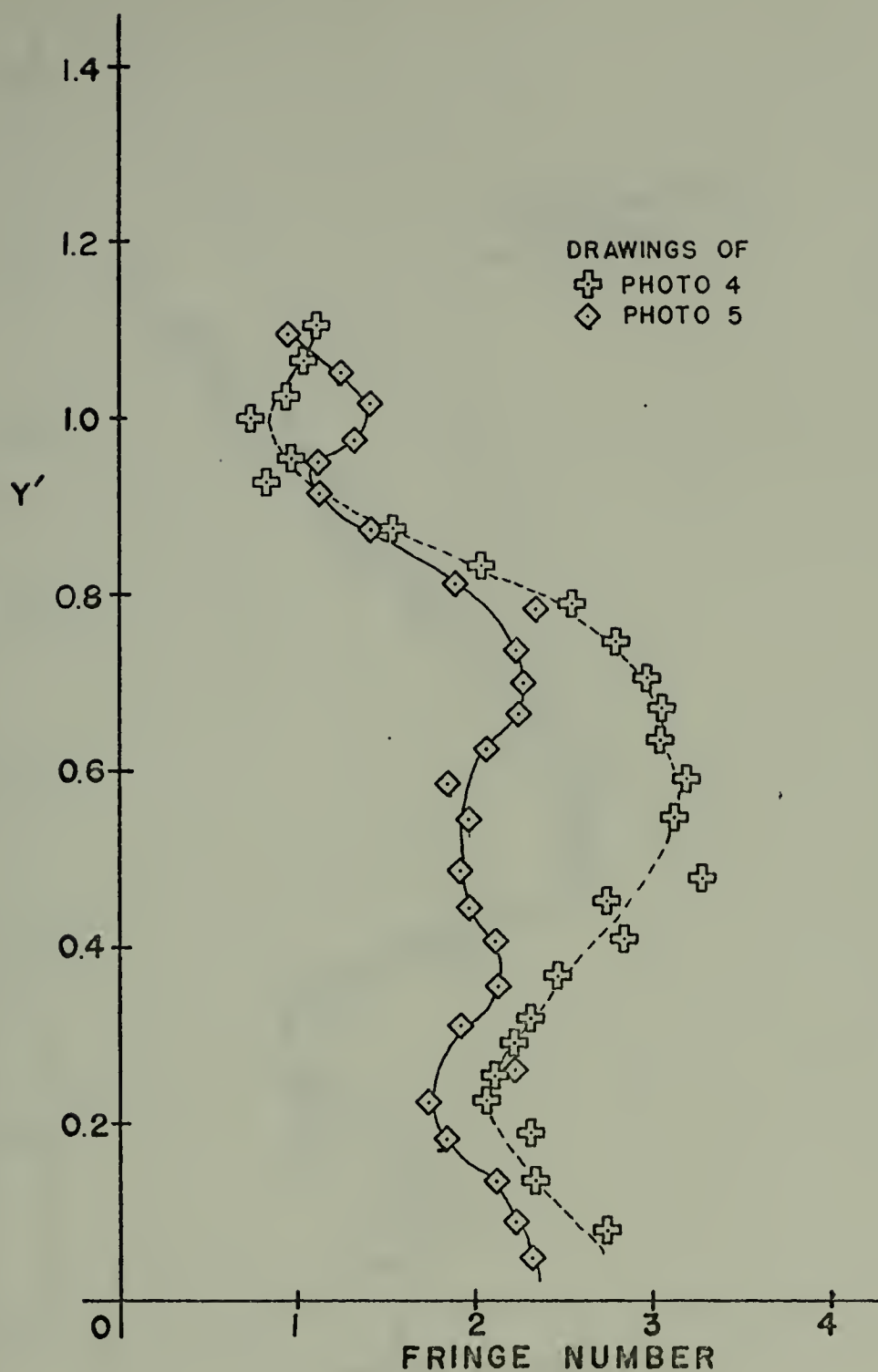


Figure 42. Comparison of Fringe Data Obtained from Drawings of Interferogram Photographs 4 and 5 Aligned at  $Y' = 0.055$



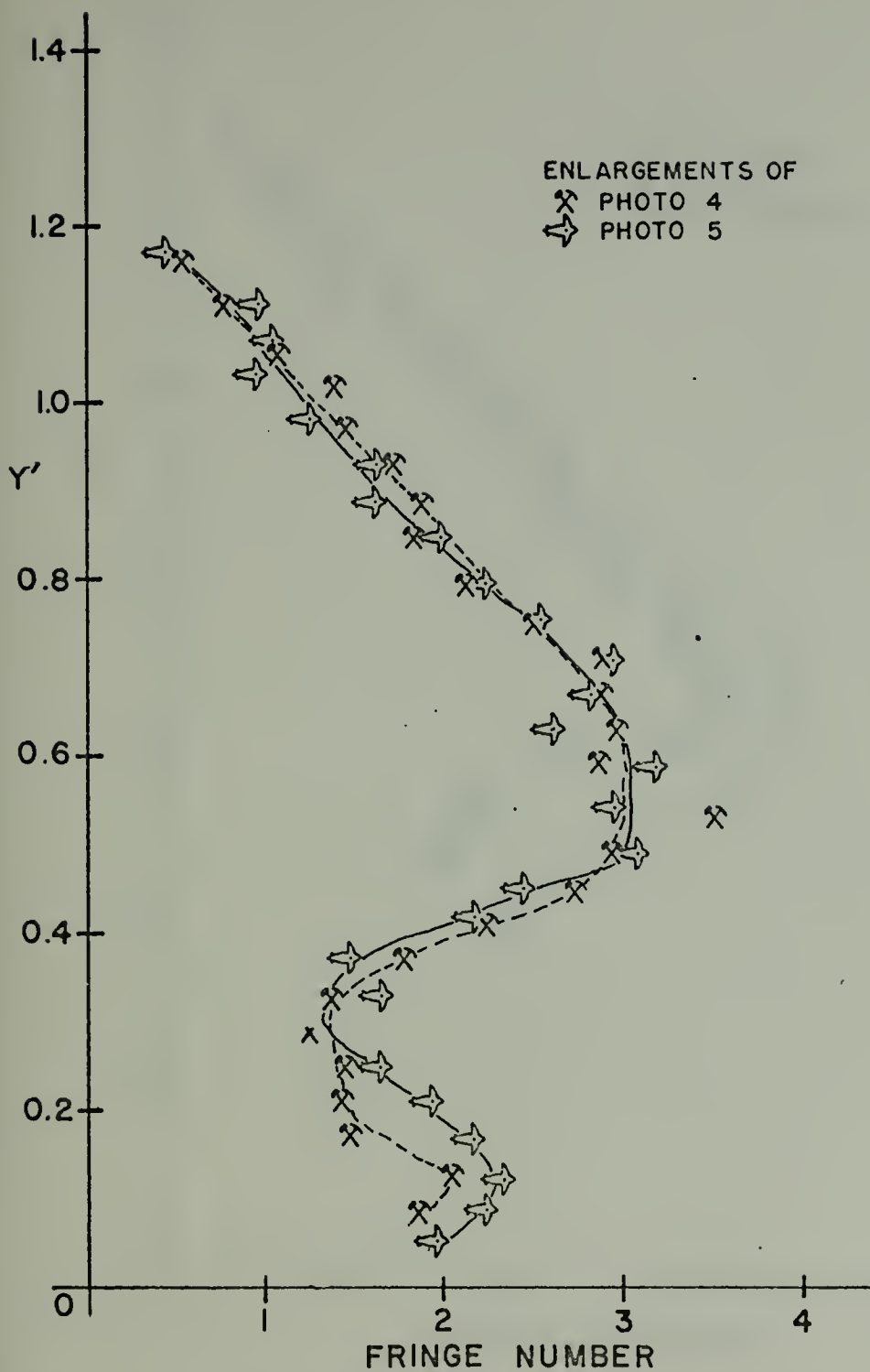


Figure 43. Comparison of Fringe Data Obtained from Photographic Enlargements of Interferogram Photographs 4 and 5 Aligned at  $Y' = 0.055$





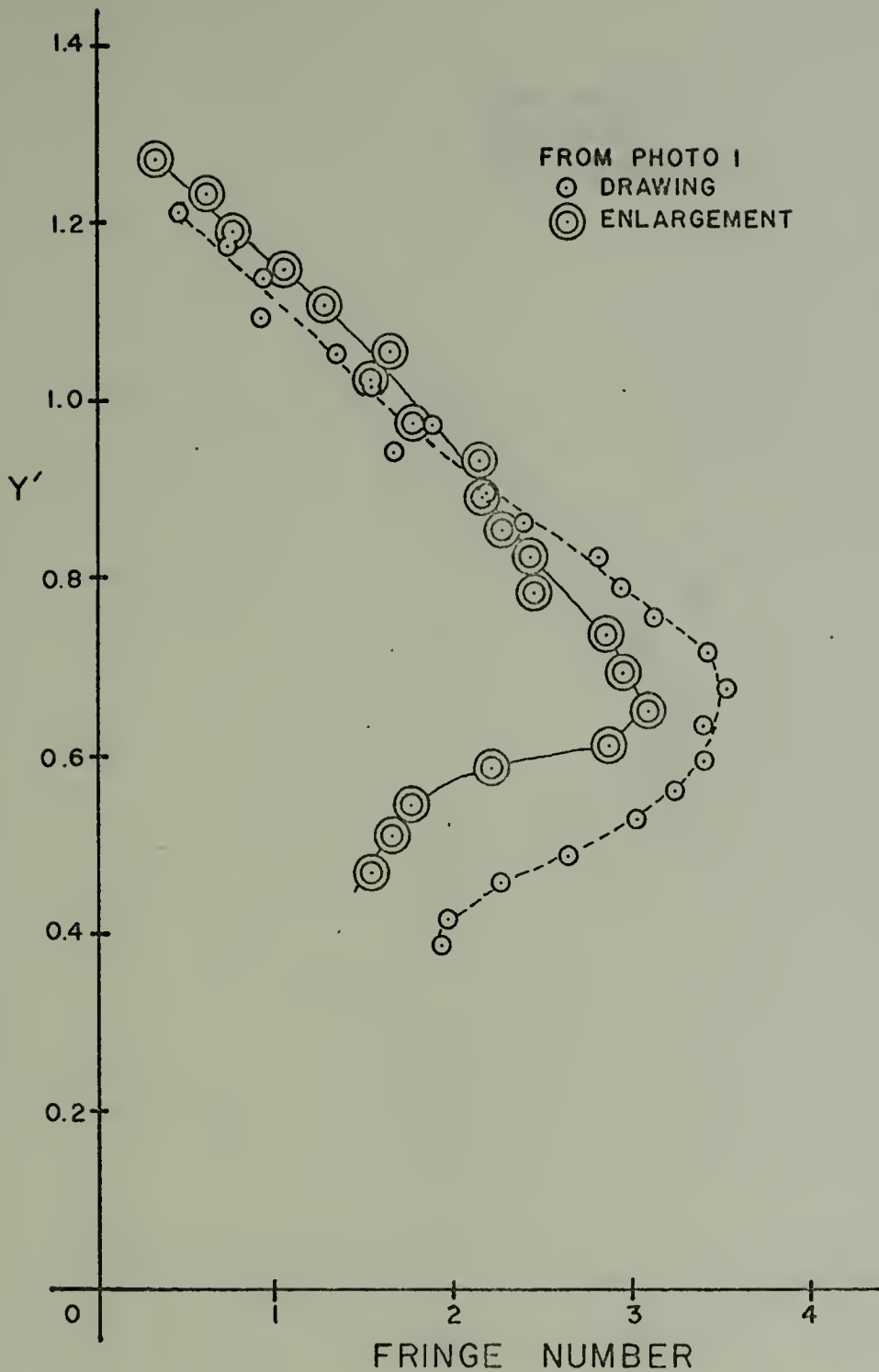


Figure 44. Comparison of Fringe Data Obtained from the Drawing and Photographic Enlargement of Interferogram Photograph 1 Aligned at  $Y' = 0.847$



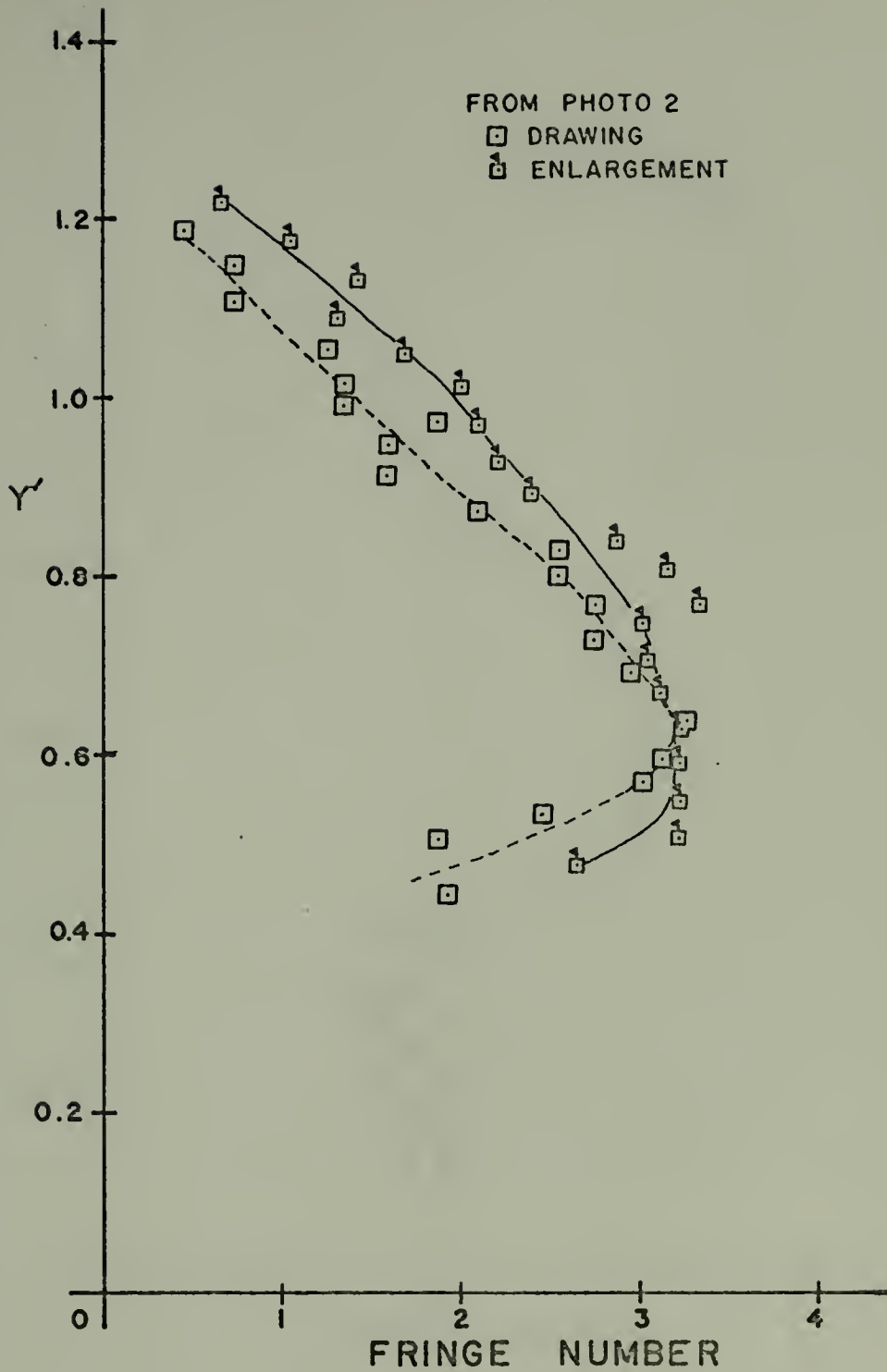


Figure 45. Comparison of Fringe Data Obtained from the Drawing and Photographic Enlargement of Interferogram Photograph 2 Aligned at  $Y' = 0.847$



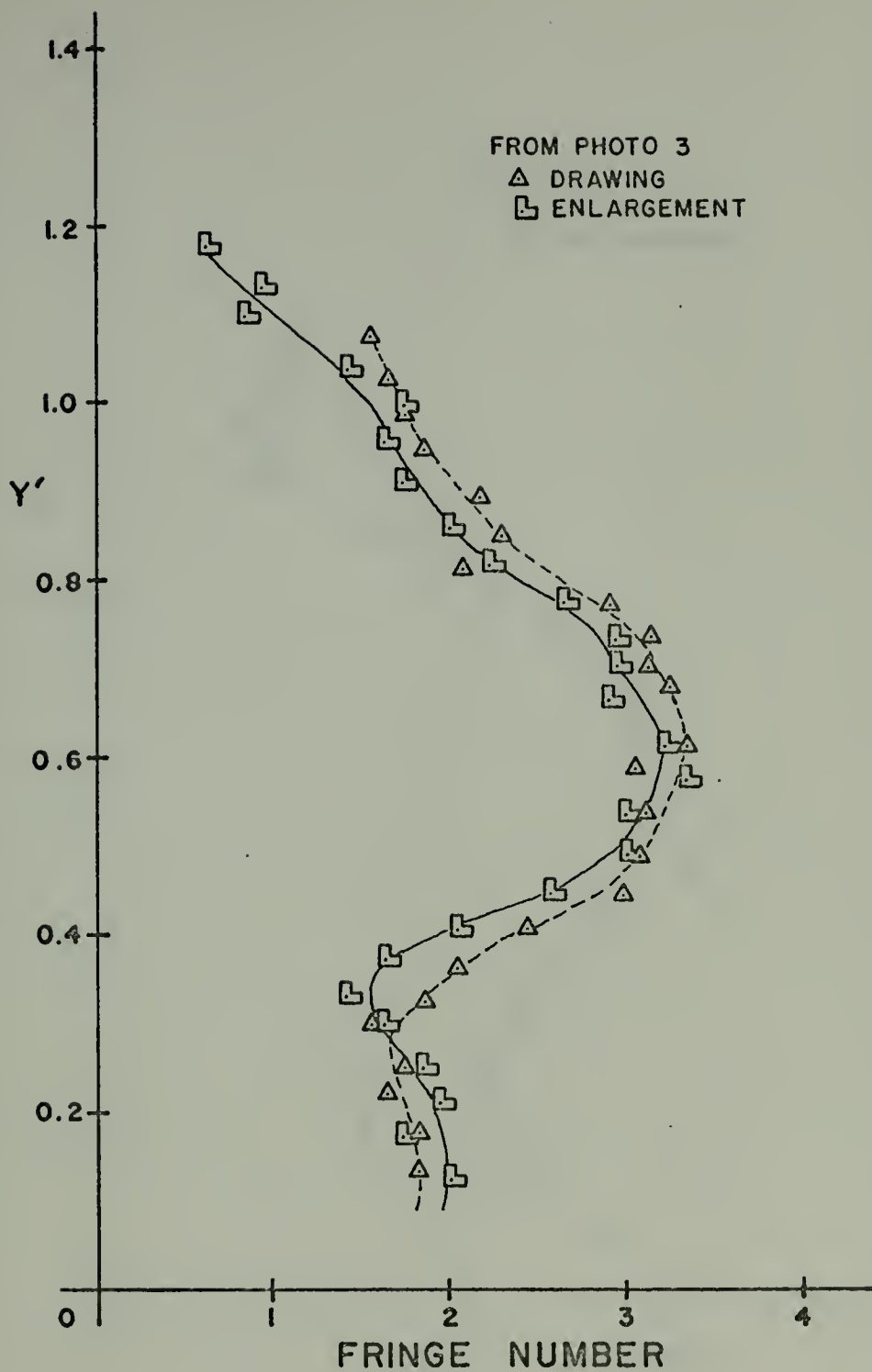


Figure 46. Comparison of Fringe Data Obtained from the Drawing and Photographic Enlargement of Interferogram Photograph 3 Aligned at  $Y' = 0.451$



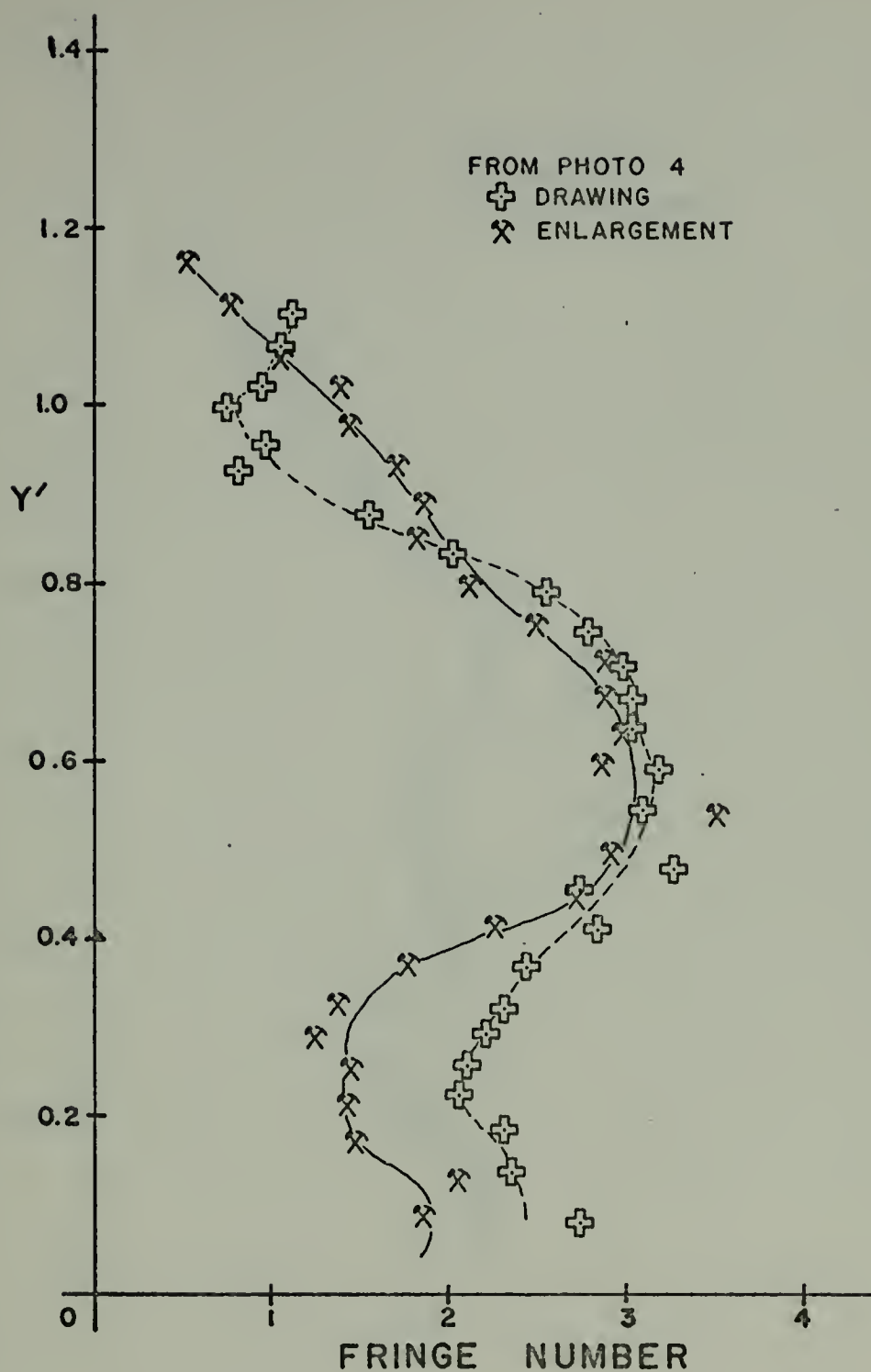


Figure 47. Comparison of Fringe Data Obtained from the Drawing and Photographic Enlargement of Interferogram Photograph 4 Aligned at  $Y' = 0.055$





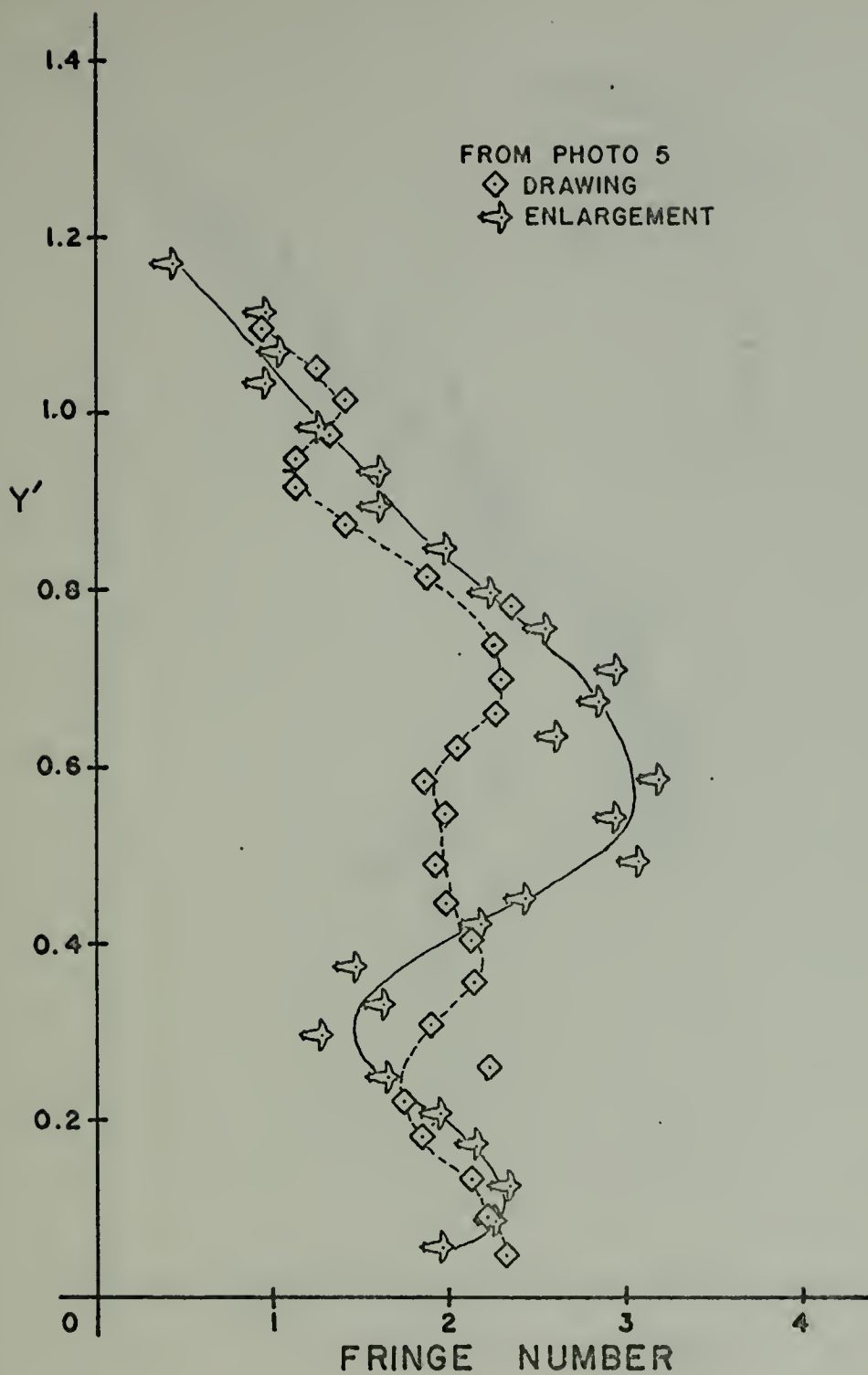


Figure 48. Comparison of Fringe Data Obtained from the Drawing and Photographic Enlargement of Interferogram Photograph 5 Aligned at  $Y' = 0.055$



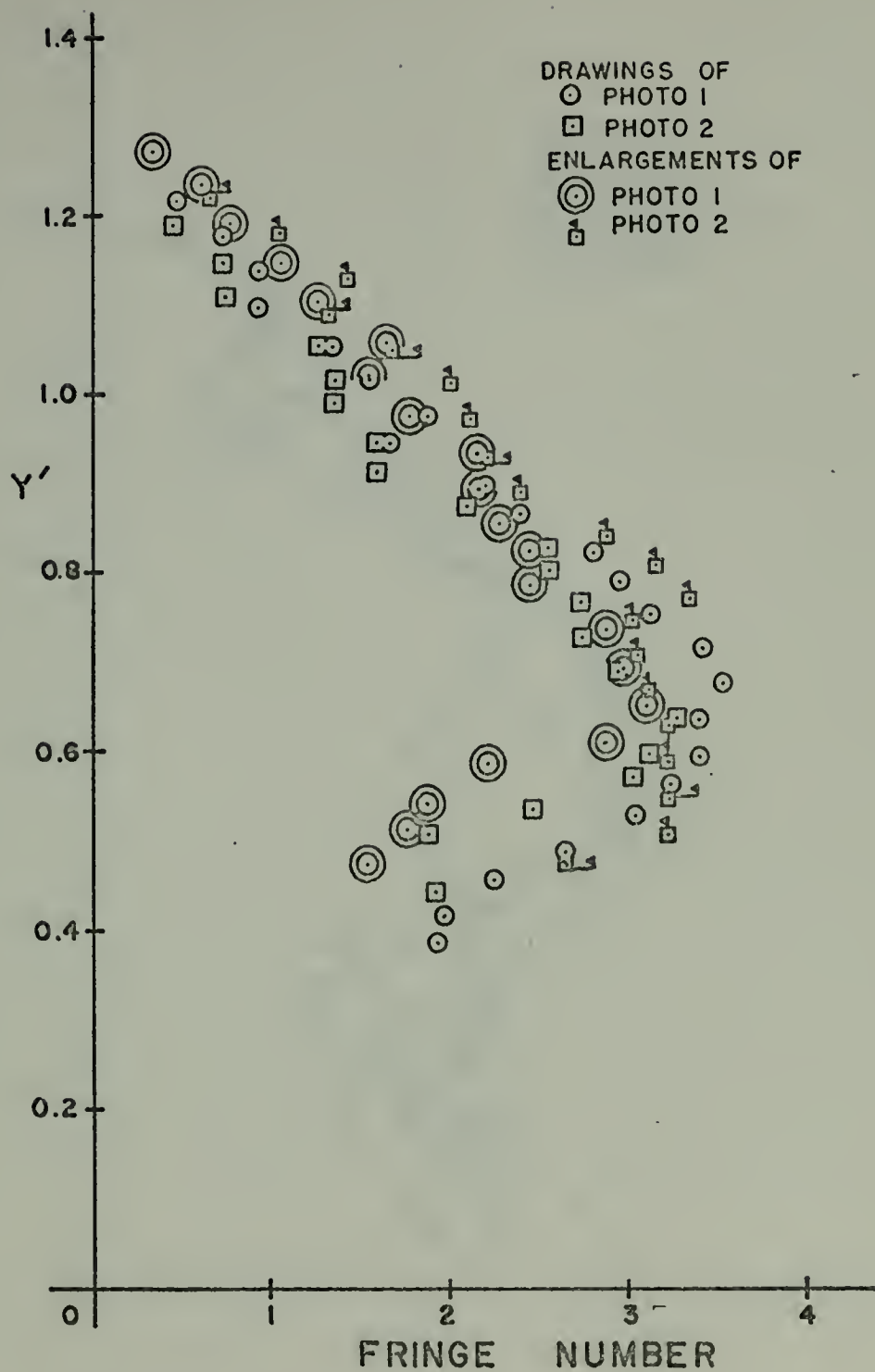


Figure 49. Comparison of Fringe Data Obtained from the Drawings and Photographic Enlargements of Interferogram Photographs 1 and 2 Aligned at  $Y' = 0.847$



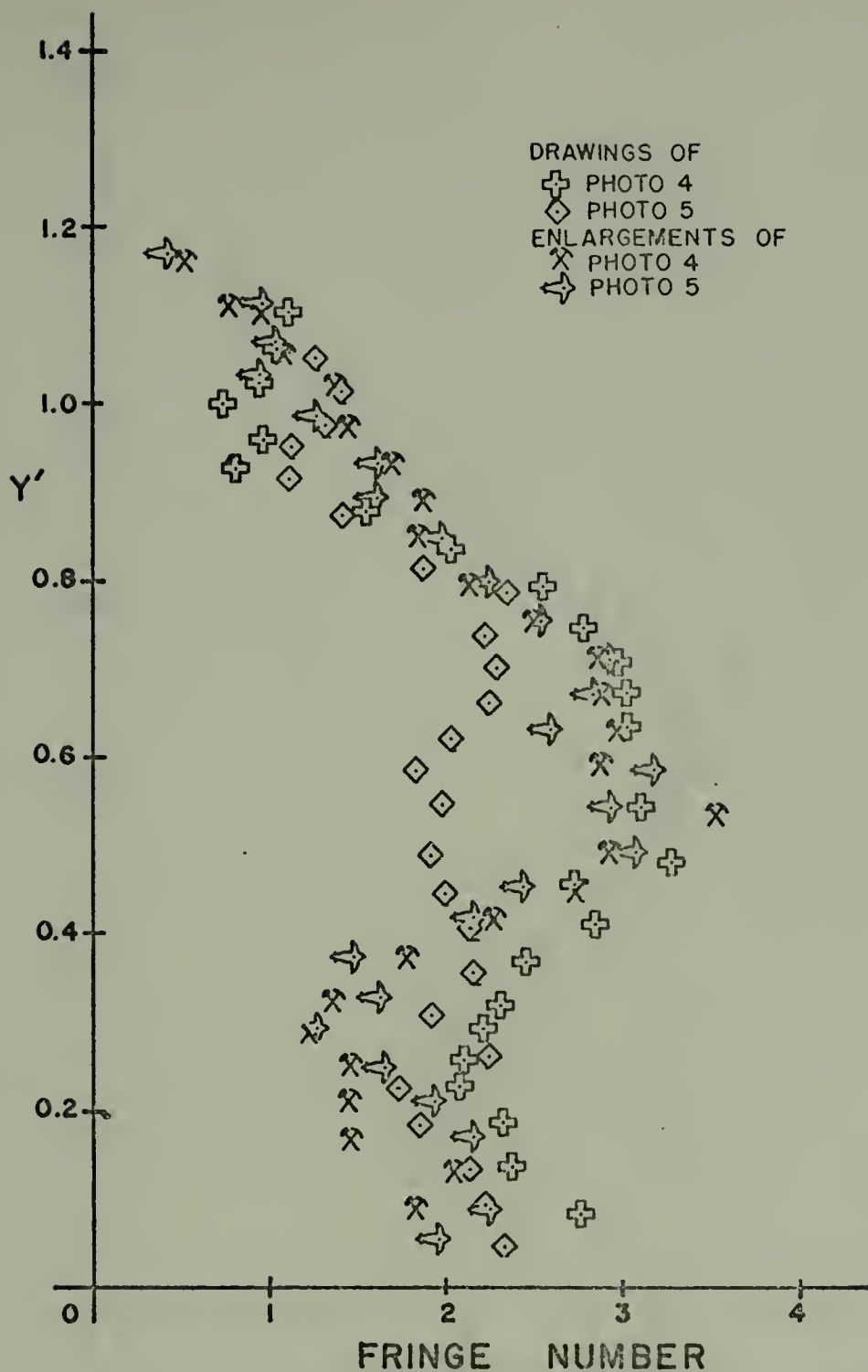


Figure 50. Comparison of Fringe Data Obtained from the Drawings and Photographic Enlargements of Interferogram Photographs 4 and 5 Aligned at  $Y' = 0.055$



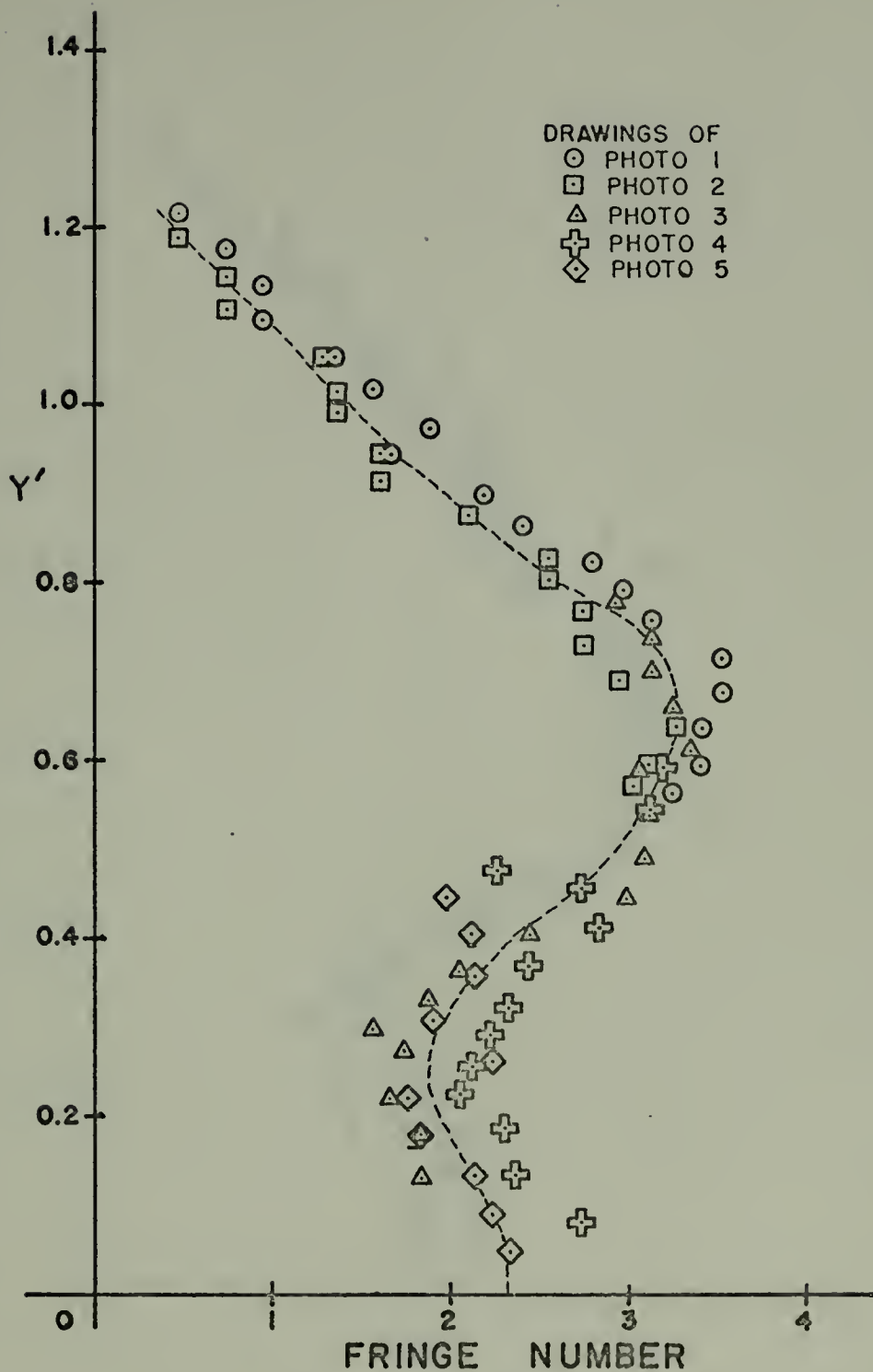


Figure 51. Fringe Number Across the Fin in the  $Z = 0.387$  Plane for  $\beta = 0^\circ$ , Mach 2.84 as Determined from the Drawings of the Interferogram Photographs





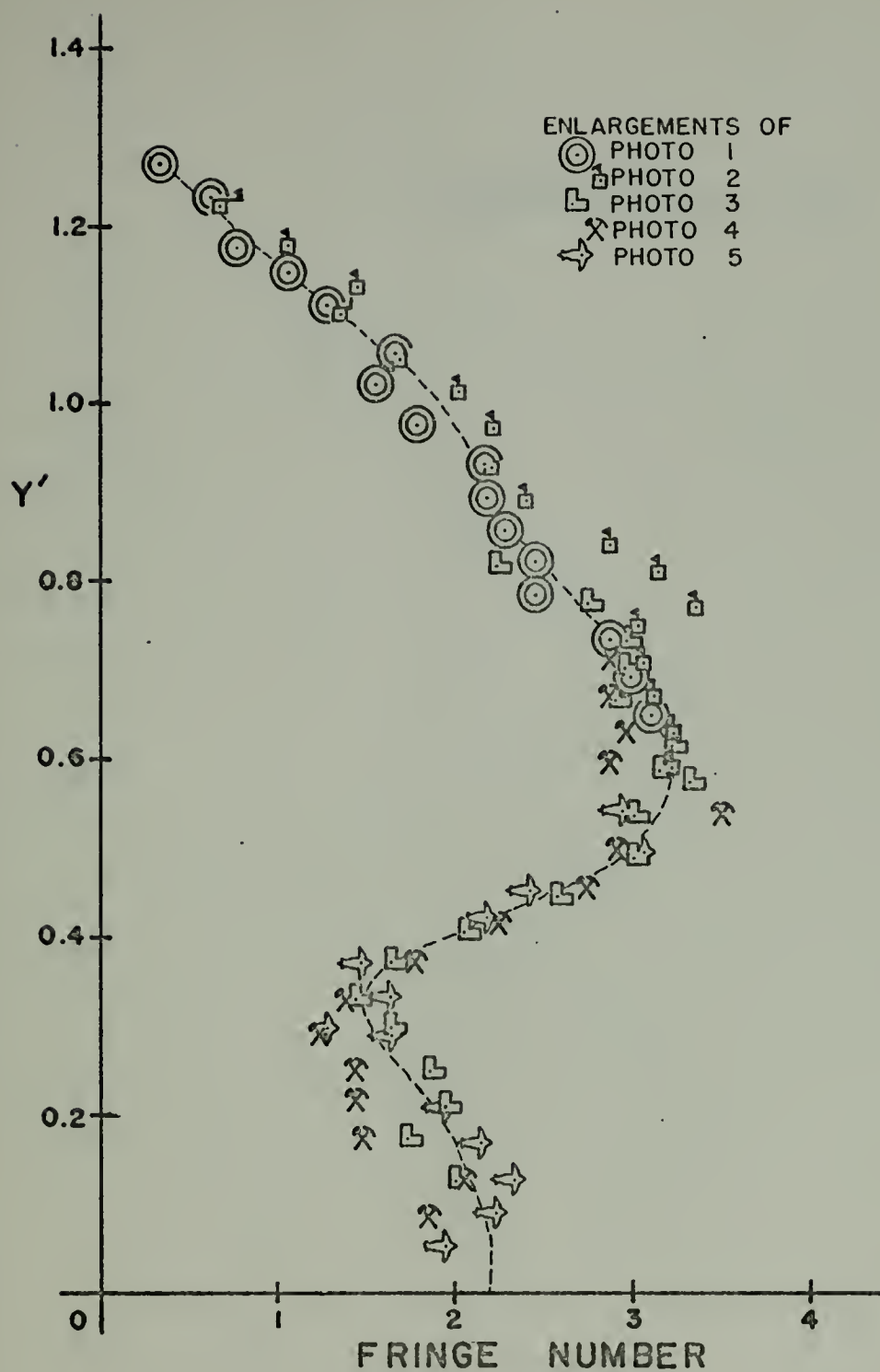


Figure 52. Fringe Number Across the Fin in the  $Z = 0.387$  Plane for  $\beta = 0^\circ$ , Mach 2.84 as Determined from the Photographic Enlargements of the Interferogram Photographs



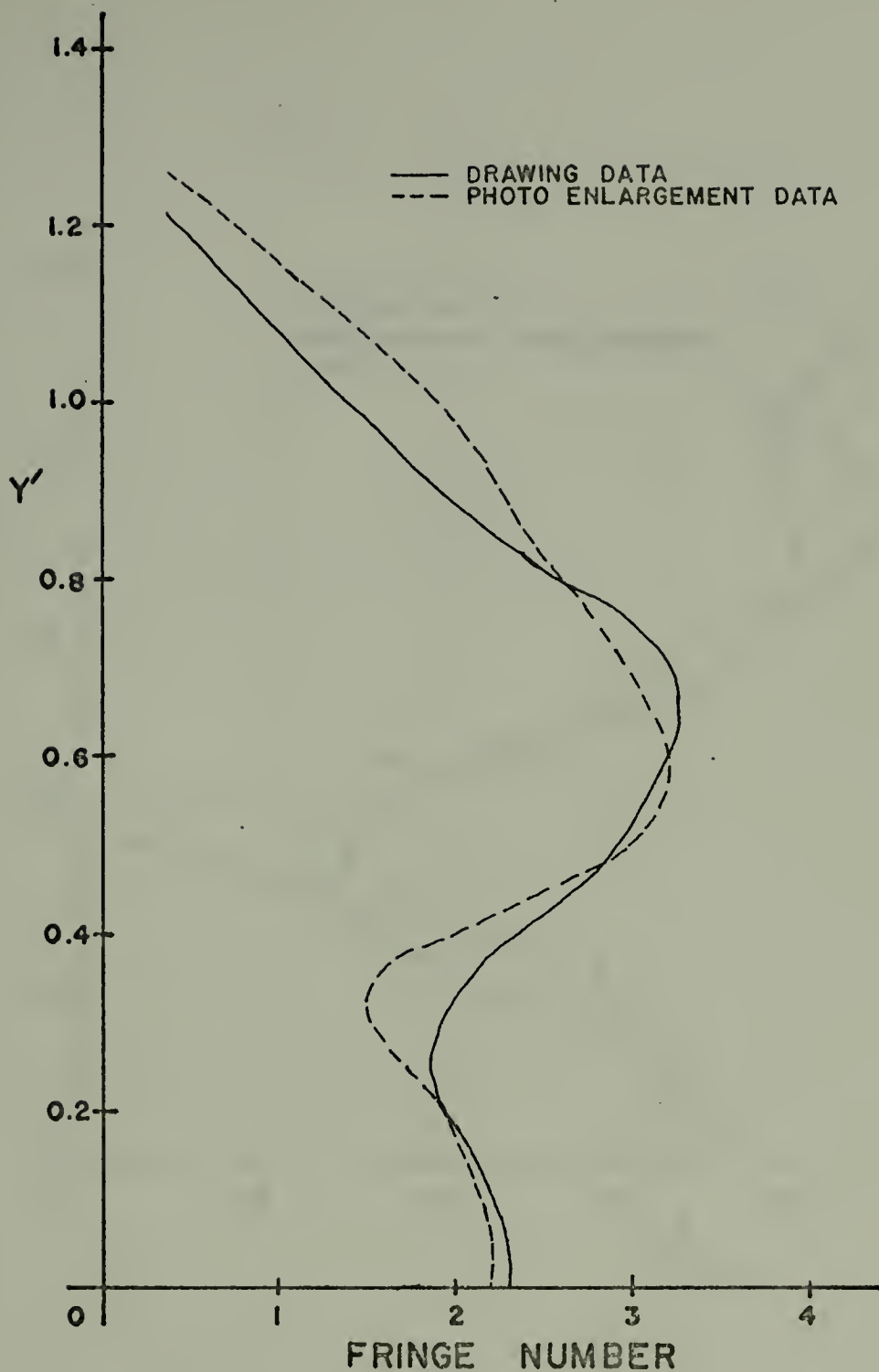


Figure 53. Comparison of the Fringe Numbers Across the Fin in the  $Z = 0.387$  Plane for  $\xi = 0^\circ$ , Mach 2.84 as Determined from the Drawings and Photographic Enlargements of the Interferogram Photographs



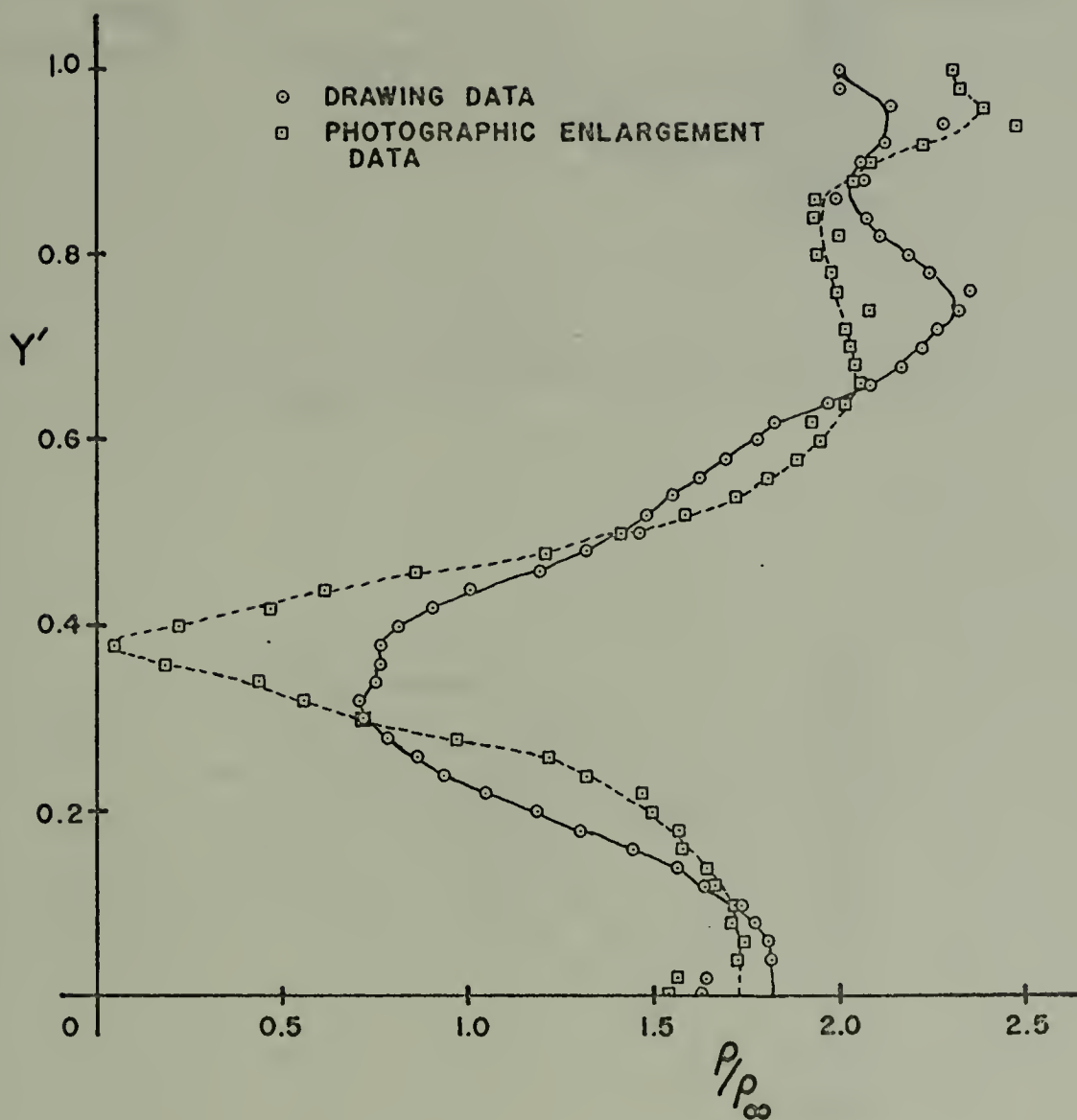


Figure 54. Comparison of the Density Distributions Calculated by HOLOVLR for an Axisymmetric Case



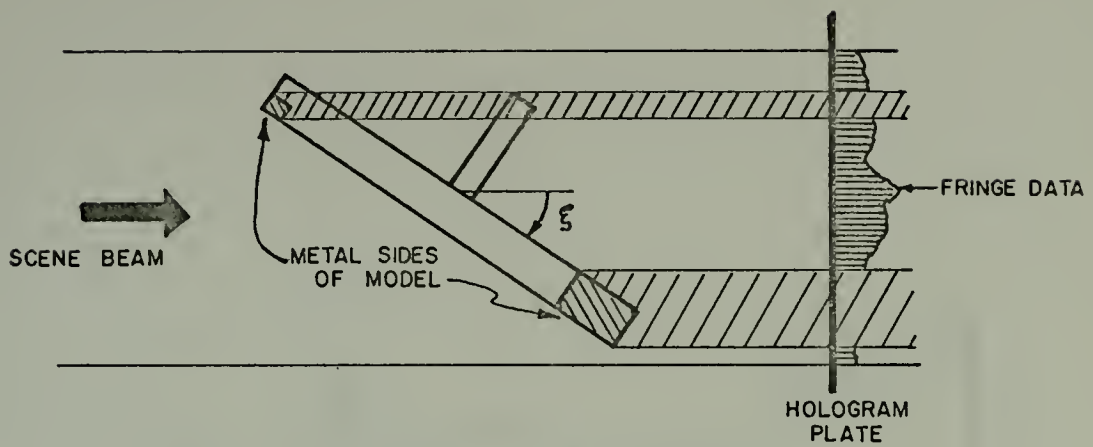


Figure 55. Schematic of the Model Center Section Rotated to Illustrate the Loss of Fringe Information Due to Model Shadows on the Hologram

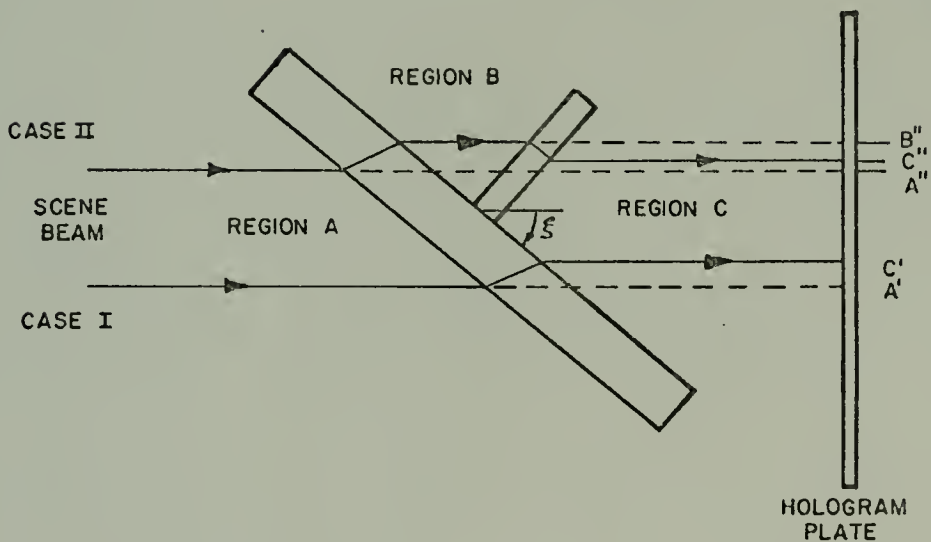


Figure 56. Schematic Illustrating the Problem of Different Fringe Information Being Superimposed on One Beam Caused by the Model Plastic





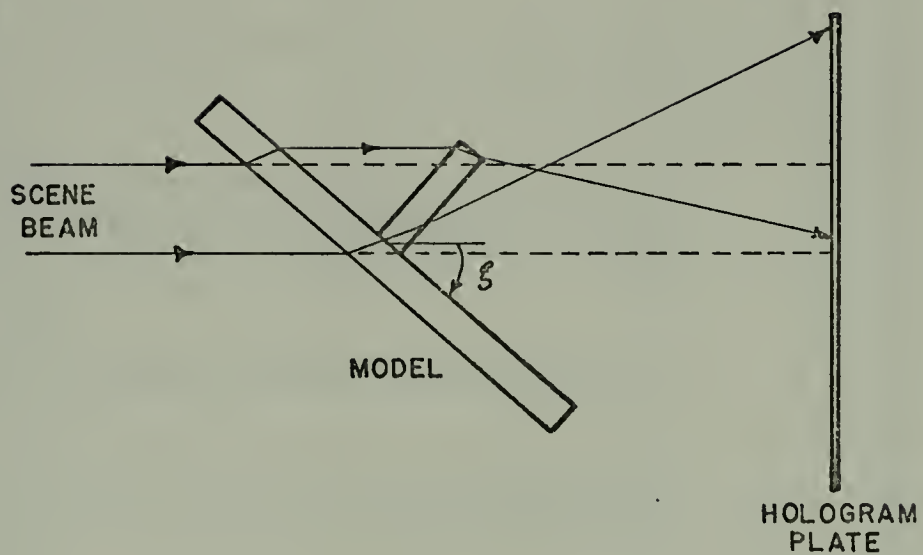


Figure 52. Schematic Illustrating the Scene Beam Refraction in the Fin Root and Tip Areas



Line No.	Distance Above Aligned Plane			Y' $\odot = \odot + .055$	Fringe Location (inches) $\odot$	Fringe Change $f = \odot - \odot$	Fringe Number $g = \frac{f}{h}$	Amount of 1.1 Fringe Correction Required	Corrected Values	
	Measured (inches)	Normalized $\odot$ wing height	Corrected Distance from Tables						Y'	g
1	.190	.075	.071	.126	.360	.170	1.641	1.1	.081	2.74
2	.330	.131	.124	.179	.460	.130	1.255	1.1	.134	2.36
3	.465	.185	.176	.231	.590	.125	1.208	1.1	.186	2.31
4	.570	.226	.215	.270	.670	.100	.965	1.1	.225	2.07
5	.660	.262	.249	.304	.765	.105	1.013	1.1	.259	2.11
6	.745	.296	.281	.336	.860	.115	1.110	1.1	.291	2.21
7	.825	.327	.311	.366	.950	.125	1.208	1.1	.321	2.31
8	.950	.377	.358	.413	1.090	.140	1.352	1.1	.368	2.45
9	1.060	.421	.400	.455	1.240	.180	1.737	1.1	.410	2.84
10	1.180	.468	.445	.500	1.350	.170	1.641	1.1	.455	2.74
11	1.245	.494	.469	.524	1.500	.255	2.172	1.1	.479	3.27
12	1.350	.536	.509	.564	1.620	.270	2.607	.5	.543	3.11
13	1.430	.567	.539	.594	1.750	.320	3.090	.1	.590	3.19
14	1.533	.609	.579	.634	1.850	.315	3.041	0.	.634	3.04
15	1.625	.649	.617	.672	1.940	.315	3.041	0.	.672	3.04
16	1.720	.683	.649	.704	2.030	.310	2.994	0.	.704	2.99
17	1.830	.726	.690	.745	2.120	.290	2.800	0.	.745	2.80
18	1.955	.776	.737	.792	2.220	.265	2.560	0.	.792	2.56
19	2.080	.825	.783	.838	2.290	.210	2.028	0.	.838	2.03
20	2.180	.865	.821	.876	2.340	.160	1.545	0.	.876	1.55
21	2.315	.919	.872	.927	2.400	.085	.820	0.	.927	.82
22	2.400	.952	.904	.959	2.500	.100	.965	0.	.959	.97
23	2.500	.992	.942	.997	2.580	.080	.772	0.	.997	.77
24	2.570	1.020	.969	1.024	2.670	.100	.965	0.	1.024	.97
25	2.680	1.063	1.009	1.064	2.790	.110	1.062	0.	1.064	1.06
26	2.780	1.103	1.048	1.103	2.885	.115	1.110	0.	1.103	1.11

Wing height = 2.52 inches

Wing length = 2.92 inches

Average fringe interval = .1036 inches

Table I. Tabulation of Fringe Shift Data Taken From the Drawing of Photo 4



Line No.	Distance Above Aligned Plane				Y' ④ = ③ + .055	Fringe Location (inches) ⑤	Fringe Change f = ⑤ - ②	Fringe Number $g = \frac{f}{b}$
	Measured (inches) ①	Corrected to Free Stream (inches) ②	Normalized $\frac{②}{(\text{wing height})}$	Corrected Distance from Tables ③				
1	.190*	.076	.032	.030	.085	.270	.194	1.867
2	.290*	.176	.073	.069	.124	.390	.214	2.050
3	.410*	.296	.123	.117	.172	.450	.154	1.482
4	.525*	.411	.171	.162	.217	.560	.149	1.434
5	.615*	.501	.208	.198	.253	.650	.149	1.434
6	.710*	.596	.247	.235	.290	.725	.129	1.242
7	.800*	.686	.285	.271	.326	.830	.144	1.386
8	.910*	.796	.330	.314	.369	.980	.184	1.771
9	1.020*	.906	.376	.357	.412	1.140	.234	2.252
10	1.120*	1.006	.417	.396	.451	1.290	.284	2.738
11	1.230*	1.116	.463	.440	.495	1.420	.304	2.926
12	1.330*	1.216	.505	.480	.535	1.580	.364	3.509
13	1.370	1.370	.568	.540	.595	1.670	.300	2.887
14	1.460	1.460	.606	.576	.631	1.770	.310	2.984
15	1.560	1.560	.647	.615	.670	1.860	.300	2.887
16	1.660	1.660	.689	.655	.710	1.960	.300	2.887
17	1.760	1.760	.730	.694	.749	2.020	.260	2.502
18	1.880	1.880	.780	.741	.796	2.100	.220	2.117
19	2.000	2.000	.830	.788	.843	2.190	.190	1.829
20	2.110	2.110	.876	.832	.887	2.305	.195	1.877
21	2.220	2.220	.921	.875	.930	2.390	.170	1.704
22	2.330	2.330	.967	.919	.974	2.480	.150	1.444
23	2.445	2.445	1.015	.964	1.019	2.590	.145	1.396
24	2.570	2.570	1.066	1.013	1.068	2.680	.110	1.059
25	2.680	2.680	1.112	1.056	1.111	2.760	.080	.770
26	2.805	2.805	1.164	1.106	1.161	2.860	.055	.529

\* Locations to be corrected to free stream conditions

Table II. Tabulation of Fringe Shift Data Taken From the Photographic Enlargement of Photo 4



## APPENDIX A

### REDUCTION OF AN INTERFEROGRAM TO OBTAIN FRINGE SHIFT DATA

The fringe shift reduction process was accomplished using two techniques. The first involved projecting the interferogram negative onto a sheet of white paper using a photo-enlarger. The light fringes offered the best contrast and were therefore traced out in Figures 30-34. In each drawing it was necessary to begin tracing the fringes above the fin and work towards the fin root since the transition across the fin tip determined the correct connection of the fringes across the fin leading edge shock. In order to determine the fringe change, one fringe line in the free stream region forward of the fin which appeared the straightest and paralleled the majority of other fringe lines was selected. A straight line, called the fringe reference line, was drawn over its centerline and extended to cross the  $y'$  axis. The remaining reference lines were then drawn parallel to the first and along the centerlines of the remaining free stream fringes. In reducing the drawings it was not realized until later that the free stream fringe patterns before and after the plate leading edge Prandtl-Meyer expansion differed considerably. This effect was taken into account later.

A ruler scaled to 0.01 inches was then placed along the  $y'$  axis and the distances of the fringes and reference lines above and below the aligned  $y'$  plane were then recorded in Table I for Photograph 4. The fin width and length were then measured and the average values were recorded. The average fringe interval was determined by measuring the distance between the first and last fringes used and dividing by the number of intervals. The fringe change was found by subtracting the fringe crossing





point. The fringe number was calculated by dividing the change by the average fringe interval. The reference line location was then normalized with respect to the measured fin height. Since the actual fin height was very close to one inch, the above number was considered to be the number of inches above or below the aligned plane. Thus the table computing the tunnel wall and grid refraction displacements in Appendix B could be entered to determine the actual reference line location with respect to the aligned plane of the drawing (see Figure 29). The locations were then converted to the  $y'$  axis system by adding the normalized location of the aligned plane. After realizing that not all the reference lines were referenced to the free stream density forward of the plate leading edge, the enlarged photographs (Figures 35-39) were checked against the drawings. It was found that the Prandtl-Meyer expansion caused a fringe number change of approximately 1.1. Consequently each reference line in the drawing was compared with the enlarged photograph and an appropriate percentage of the 1.1 fringe number was used as a correction (see Table I). The calculated fringe numbers had the correction fringe number added to them while the  $y'$  locations were corrected by

$$y'_{\text{corr}} = y'_{\text{orig}} - \frac{(1.1 \text{ Fringe no.}) \times (\text{Fringe Interval})}{(\text{Fin Height})} \quad (\text{A-1})$$

The second reduction technique was to use enlarged photographs made from the interferogram negatives to obtain the fringe change. The fringe lines were first traced over lightly with a pencil and then verified against the other photographs to ensure correct tracing. A datum fringe reference line was chosen as before and drawn. The remaining reference lines for the upper free stream fringe lines forward of the Prandtl-Meyer expansion were then drawn parallel to the datum. When it became impossible



to use fringe lines forward of the expansion, the reference lines were then drawn along the centerline of the fringes between the expansion and the fin. The same ruler was used to obtain the fringe crossing points, the reference crossing points, fin measurements, and the average fringe interval and they were recorded in Table 2. The reference line crossing points were then corrected to free stream conditions, using a 1.1 Fringe number correction for those referenced to the pattern between the expansion and fin leading edge. The fringe number and reference fin locations were calculated as before. This method was considered more accurate because the reference lines and fringe lines could easily be rechecked for accuracy and corrected in the event that a fringe line was traced incorrectly or a reference line was misaligned. The accuracy was directly proportional to the hologram resolution which was not true for the drawings since the reference lines are drawn parallel to hand-drawn fringe lines.



## APPENDIX B

### CALCULATION OF TUNNEL WALL AND GRID PLASTIC REFRACTION CORRECTION

In the interferogram photographs used to obtain the fringe change, every point off the alignment axis will be slightly distorted due to the plastic tunnel wall and grid. This effect is illustrated in Figure 19. From Snells Law of Refraction, the angles of incidence and refraction are related by

$$\sin \alpha = n \sin \beta \quad (B-1)$$

where  $n$  is the index of refraction between plastic and air. Then can be written

$$\beta = \sin^{-1} \left( \frac{\sin \alpha}{n} \right) \quad (B-2)$$

Since the tunnel wall and grid have the same index of refraction, they can be considered on material with thickness  $t = ab$  in Figure 19. Then consider the height  $bd$  which can be written

$$bd = t \tan \alpha \quad (B-3)$$

The beam displacement,  $\Delta y$ , is then

$$\Delta Y = bd - bc = t \tan \alpha - t \tan \beta \quad (B-4)$$

but  $\tan \alpha$  is

$$\tan \alpha = \frac{Y_{\text{observed}}}{L} \quad (B-5)$$

Combining Equation (B-2), (B-4), and (B-5) the beam displacement becomes

$$\Delta Y = t \left[ \frac{Y_{\text{observed}}}{L} - \tan \left[ \sin^{-1} \left( \frac{\sin \alpha}{n} \right) \right] \right] \quad (B-6)$$

The true location of the observed point is then

$$Y_{\text{true}} = Y_{\text{observed}} - \Delta Y \quad (B-7)$$



A FORTRAN computer program was written to generate a table giving the true locations versus the observed locations. In the program the constants and variables from the above equations were defined as

$Y_{\text{observed}}$	=	Y
$Y_{\text{true}}$	=	YTRUE
$\Delta Y$	=	DY
$\alpha$	=	ALFA
L	=	L
n	=	N

Since the computer cannot calculate Equations (B-2) and (B-6) as written, they were constructed by parts using such letters as AA, AB, etc. The program and tables are included in the next few pages.





```

*****
**
** THIS PROGRAM COMPUTES THE CORRECTION FOR THE
** OFF AXIS TUNNEL WALL PARALLAX. THE FOLLOWING
** INPUT PARAMETERS ARE REQUIRED:
** L = DISTANCE FROM HOLOGRAM PLANE TO THE MODEL
** CENTER LINE
** N = INDEX OF REFRACTION FOR TUNNEL WALL
** T = THICKNESS OF TUNNEL AND GRID PLEXIGLAS
** Y = OBSERVED DISTANCE OF THE POINT IN THE
** PHOTOGRAPH FROM THE ALIGNED POINT
** YTRUE = TRUE DISTANCE FROM ALIGNED POINT TO
** THE POINT ON THE MODEL PLANE
** YMAX = MAXIMUM DISTANCE OF INTEREST FROM THE
** ALIGNED POINT
**
*****

```

```

REAL*4 L,N

```

# INPUT PARAMETERS

```

L=15.0
N=1.5
T=2.25
YMAX=2.5
PI=3.141593

```

# OUTPUT FORMAT

```

WRITE(6,100)L,N,T
100 FORMAT('1',T25,'THE FOLLOWING TABLE IS TO ACCOUNT FOR'
1 'THE PARALLAX'/T25,'ERROR IN VIEWING THE MODEL '
2 'THROUGH GLASS WALLS AT'/T25,'ANY OTHER POINT THAN '
3 'AT THE ALIGNMENT POINT. INPUT'/T25,'PARAMETERS ARE:'
4 '/T35,'(1) L = ',F6.3,' INCHES'/T35,'(2) N = ',F6.3,
5 ' INCHES'/T35,'(3) T = ',F6.3,' INCHES')
WRITE(6,110)
110 FORMAT(' ',T24,'OBSERVED',T40,'TRUE',T54,'ERROR',T70,
1 'RAY'/T24,'DISTANCE',T38,'DISTANCE',T69,'ANGLE'/T24,
2 '(INCHES)',T38,'(INCHES)',T53,'(INCHES)',T67,
3 '(DEGREES)'/)

```

# CALCULATIONS

```

Y=0.0
10 DO 15 I=1,2000
AA=Y/L
AB=ATAN(AA)
ALFA=AB*180.0/PI
AC=SIN(AB)/N
BETA=ARSIN(AC)
AD=TAN(BETA)
DY=T*(AA-AD)
YTRUE=Y-DY
C
WRITE(6,200)Y,YTRUE,DY,ALFA
200 FORMAT(' ',T24,F7.4,T38,F7.4,T53,F8.6,T67,F8.4)
IF(MOD(I,10).EQ.0)WRITE(6,201)
201 FORMAT(' ')
IF(MOD(I,60).EQ.0)GO TO 11
GO TO 12
11 WRITE(6,202)
202 FORMAT('1')
WRITE(6,110)
12 Y=.002*FLOAT(I)
IF(Y.GE.YMAX)GO TO 20
15 CONTINUE
20 STOP
END

```



THE FOLLOWING TABLE IS TO ACCOUNT FOR THE PARALLAX ERROR IN VIEWING THE MODEL THROUGH GLASS WALLS AT ANY OTHER POINT THAN AT THE ALIGNMENT POINT. INPUT PARAMETERS ARE:

- (1) L = 15.000 INCHES
- (2) N = 1.500 INCHES
- (3) T = 2.250 INCHES

OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.0	0.0	0.0	0.0
0.0020	0.0019	0.000100	0.0076
0.0040	0.0038	0.000200	0.0153
0.0060	0.0057	0.000300	0.0229
0.0080	0.0076	0.000400	0.0306
0.0100	0.0095	0.000500	0.0382
0.0120	0.0114	0.000600	0.0458
0.0140	0.0133	0.000700	0.0535
0.0160	0.0152	0.000800	0.0611
0.0180	0.0171	0.000900	0.0688
0.0200	0.0190	0.001000	0.0764
0.0220	0.0209	0.001100	0.0840
0.0240	0.0228	0.001200	0.0917
0.0260	0.0247	0.001300	0.0993
0.0280	0.0266	0.001400	0.1070
0.0300	0.0285	0.001500	0.1146
0.0320	0.0304	0.001600	0.1222
0.0340	0.0323	0.001700	0.1299
0.0360	0.0342	0.001800	0.1375
0.0380	0.0361	0.001900	0.1451
0.0400	0.0380	0.002000	0.1528
0.0420	0.0399	0.002100	0.1604
0.0440	0.0418	0.002200	0.1681
0.0460	0.0437	0.002300	0.1757
0.0480	0.0456	0.002400	0.1833
0.0500	0.0475	0.002500	0.1910
0.0520	0.0494	0.002600	0.1986
0.0540	0.0513	0.002700	0.2063
0.0560	0.0532	0.002800	0.2139
0.0580	0.0551	0.002900	0.2215
0.0600	0.0570	0.003000	0.2292
0.0620	0.0589	0.003100	0.2368
0.0640	0.0608	0.003200	0.2445
0.0660	0.0627	0.003300	0.2521
0.0680	0.0646	0.003400	0.2597
0.0700	0.0665	0.003500	0.2674
0.0720	0.0684	0.003600	0.2750
0.0740	0.0703	0.003700	0.2827
0.0760	0.0722	0.003800	0.2903
0.0780	0.0741	0.003900	0.2979
0.0800	0.0760	0.004000	0.3056
0.0820	0.0779	0.004100	0.3132
0.0840	0.0798	0.004200	0.3209
0.0860	0.0817	0.004300	0.3285
0.0880	0.0836	0.004400	0.3361
0.0900	0.0855	0.004500	0.3438
0.0920	0.0874	0.004600	0.3514
0.0940	0.0893	0.004700	0.3590
0.0960	0.0912	0.004800	0.3667
0.0980	0.0931	0.004900	0.3743
0.1000	0.0950	0.005000	0.3820
0.1020	0.0969	0.005100	0.3896
0.1040	0.0988	0.005200	0.3972
0.1060	0.1007	0.005300	0.4049
0.1080	0.1026	0.005400	0.4125
0.1100	0.1045	0.005500	0.4202
0.1120	0.1064	0.005600	0.4278
0.1140	0.1083	0.005700	0.4354
0.1160	0.1102	0.005800	0.4431
0.1180	0.1121	0.005900	0.4507



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.1200	0.1140	0.006000	0.4584
0.1220	0.1159	0.006100	0.4660
0.1240	0.1178	0.006200	0.4736
0.1260	0.1197	0.006300	0.4813
0.1280	0.1216	0.006400	0.4889
0.1300	0.1235	0.006500	0.4966
0.1320	0.1254	0.006600	0.5042
0.1340	0.1273	0.006700	0.5118
0.1360	0.1292	0.006800	0.5195
0.1380	0.1311	0.006900	0.5271
0.1400	0.1330	0.007000	0.5347
0.1420	0.1349	0.007100	0.5424
0.1440	0.1368	0.007200	0.5500
0.1460	0.1387	0.007300	0.5577
0.1480	0.1406	0.007400	0.5653
0.1500	0.1425	0.007500	0.5729
0.1520	0.1444	0.007600	0.5806
0.1540	0.1463	0.007700	0.5882
0.1560	0.1482	0.007800	0.5959
0.1580	0.1501	0.007901	0.6035
0.1600	0.1520	0.008001	0.6111
0.1620	0.1539	0.008101	0.6188
0.1640	0.1558	0.008201	0.6264
0.1660	0.1577	0.008301	0.6340
0.1680	0.1596	0.008401	0.6417
0.1700	0.1615	0.008501	0.6493
0.1720	0.1634	0.008601	0.6570
0.1740	0.1653	0.008701	0.6646
0.1760	0.1672	0.008801	0.6722
0.1780	0.1691	0.008901	0.6799
0.1800	0.1710	0.009001	0.6875
0.1820	0.1729	0.009101	0.6952
0.1840	0.1748	0.009201	0.7028
0.1860	0.1767	0.009301	0.7104
0.1880	0.1786	0.009401	0.7181
0.1900	0.1805	0.009501	0.7257
0.1920	0.1824	0.009601	0.7333
0.1940	0.1843	0.009701	0.7410
0.1960	0.1862	0.009801	0.7486
0.1980	0.1881	0.009901	0.7563
0.2000	0.1900	0.010001	0.7639
0.2020	0.1919	0.010101	0.7715
0.2040	0.1938	0.010201	0.7792
0.2060	0.1957	0.010301	0.7868
0.2080	0.1976	0.010401	0.7944
0.2100	0.1995	0.010501	0.8021
0.2120	0.2014	0.010601	0.8097
0.2140	0.2033	0.010701	0.8174
0.2160	0.2052	0.010801	0.8250
0.2180	0.2071	0.010901	0.8326
0.2200	0.2090	0.011001	0.8403
0.2220	0.2109	0.011101	0.8479
0.2240	0.2128	0.011201	0.8556
0.2260	0.2147	0.011301	0.8632
0.2280	0.2166	0.011401	0.8708
0.2300	0.2185	0.011502	0.8785
0.2320	0.2204	0.011602	0.8861
0.2340	0.2223	0.011702	0.8937
0.2360	0.2242	0.011802	0.9014
0.2380	0.2261	0.011902	0.9090





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.2400	0.2280	0.012002	0.9167
0.2420	0.2299	0.012102	0.9243
0.2440	0.2318	0.012202	0.9319
0.2460	0.2337	0.012302	0.9396
0.2480	0.2356	0.012402	0.9472
0.2500	0.2375	0.012502	0.9548
0.2520	0.2394	0.012602	0.9625
0.2540	0.2413	0.012702	0.9701
0.2560	0.2432	0.012802	0.9778
0.2580	0.2451	0.012902	0.9854
0.2600	0.2470	0.013002	0.9930
0.2620	0.2489	0.013102	1.0007
0.2640	0.2508	0.013202	1.0083
0.2660	0.2527	0.013302	1.0159
0.2680	0.2546	0.013402	1.0236
0.2700	0.2565	0.013502	1.0312
0.2720	0.2584	0.013602	1.0388
0.2740	0.2603	0.013703	1.0465
0.2760	0.2622	0.013803	1.0541
0.2780	0.2641	0.013903	1.0618
0.2800	0.2660	0.014003	1.0694
0.2820	0.2679	0.014103	1.0770
0.2840	0.2698	0.014203	1.0847
0.2860	0.2717	0.014303	1.0923
0.2880	0.2736	0.014403	1.0999
0.2900	0.2755	0.014503	1.1076
0.2920	0.2774	0.014603	1.1152
0.2940	0.2793	0.014703	1.1229
0.2960	0.2812	0.014803	1.1305
0.2980	0.2831	0.014903	1.1381
0.3000	0.2850	0.015003	1.1458
0.3020	0.2869	0.015103	1.1534
0.3040	0.2888	0.015203	1.1610
0.3060	0.2907	0.015304	1.1687
0.3080	0.2926	0.015404	1.1763
0.3100	0.2945	0.015504	1.1839
0.3120	0.2964	0.015604	1.1916
0.3140	0.2983	0.015704	1.1992
0.3160	0.3002	0.015804	1.2069
0.3180	0.3021	0.015904	1.2145
0.3200	0.3040	0.016004	1.2221
0.3220	0.3059	0.016104	1.2298
0.3240	0.3078	0.016204	1.2374
0.3260	0.3097	0.016304	1.2450
0.3280	0.3116	0.016404	1.2527
0.3300	0.3135	0.016504	1.2603
0.3320	0.3154	0.016605	1.2679
0.3340	0.3173	0.016705	1.2756
0.3360	0.3192	0.016805	1.2832
0.3380	0.3211	0.016905	1.2908
0.3400	0.3230	0.017005	1.2985
0.3420	0.3249	0.017105	1.3061
0.3440	0.3268	0.017205	1.3138
0.3460	0.3287	0.017305	1.3214
0.3480	0.3306	0.017405	1.3290
0.3500	0.3325	0.017505	1.3367
0.3520	0.3344	0.017605	1.3443
0.3540	0.3363	0.017705	1.3519
0.3560	0.3382	0.017806	1.3596
0.3580	0.3401	0.017906	1.3672





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.3600	0.3420	0.018006	1.3748
0.3620	0.3439	0.018106	1.3825
0.3640	0.3458	0.018206	1.3901
0.3660	0.3477	0.018306	1.3977
0.3680	0.3496	0.018406	1.4054
0.3700	0.3515	0.018506	1.4130
0.3720	0.3534	0.018606	1.4206
0.3740	0.3553	0.018706	1.4283
0.3760	0.3572	0.018807	1.4359
0.3780	0.3591	0.018907	1.4435
0.3800	0.3610	0.019007	1.4512
0.3820	0.3629	0.019107	1.4588
0.3840	0.3648	0.019207	1.4665
0.3860	0.3667	0.019307	1.4741
0.3880	0.3686	0.019407	1.4817
0.3900	0.3705	0.019507	1.4894
0.3920	0.3724	0.019607	1.4970
0.3940	0.3743	0.019708	1.5046
0.3960	0.3762	0.019808	1.5123
0.3980	0.3781	0.019908	1.5199
0.4000	0.3800	0.020008	1.5275
0.4020	0.3819	0.020108	1.5352
0.4040	0.3838	0.020208	1.5428
0.4060	0.3857	0.020308	1.5504
0.4080	0.3876	0.020408	1.5581
0.4100	0.3895	0.020509	1.5657
0.4120	0.3914	0.020609	1.5733
0.4140	0.3933	0.020709	1.5810
0.4160	0.3952	0.020809	1.5886
0.4180	0.3971	0.020909	1.5962
0.4200	0.3990	0.021009	1.6039
0.4220	0.4009	0.021109	1.6115
0.4240	0.4028	0.021209	1.6191
0.4260	0.4047	0.021310	1.6268
0.4280	0.4066	0.021410	1.6344
0.4300	0.4085	0.021510	1.6420
0.4320	0.4104	0.021610	1.6497
0.4340	0.4123	0.021710	1.6573
0.4360	0.4142	0.021810	1.6649
0.4380	0.4161	0.021910	1.6726
0.4400	0.4180	0.022011	1.6802
0.4420	0.4199	0.022111	1.6878
0.4440	0.4218	0.022211	1.6955
0.4460	0.4237	0.022311	1.7031
0.4480	0.4256	0.022411	1.7107
0.4500	0.4275	0.022511	1.7184
0.4520	0.4294	0.022611	1.7260
0.4540	0.4313	0.022712	1.7336
0.4560	0.4332	0.022812	1.7413
0.4580	0.4351	0.022912	1.7489
0.4600	0.4370	0.023012	1.7565
0.4620	0.4389	0.023112	1.7642
0.4640	0.4408	0.023212	1.7718
0.4660	0.4427	0.023312	1.7794
0.4680	0.4446	0.023413	1.7870
0.4700	0.4465	0.023513	1.7947
0.4720	0.4484	0.023613	1.8023
0.4740	0.4503	0.023713	1.8099
0.4760	0.4522	0.023813	1.8176
0.4780	0.4541	0.023913	1.8252



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.4800	0.4560	0.024014	1.8328
0.4820	0.4579	0.024114	1.8405
0.4840	0.4598	0.024214	1.8481
0.4860	0.4617	0.024314	1.8557
0.4880	0.4636	0.024414	1.8634
0.4900	0.4655	0.024515	1.8710
0.4920	0.4674	0.024615	1.8786
0.4940	0.4693	0.024715	1.8863
0.4960	0.4712	0.024815	1.8939
0.4980	0.4731	0.024915	1.9015
0.5000	0.4750	0.025015	1.9092
0.5020	0.4769	0.025116	1.9168
0.5040	0.4788	0.025216	1.9244
0.5060	0.4807	0.025316	1.9320
0.5080	0.4826	0.025416	1.9397
0.5100	0.4845	0.025516	1.9473
0.5120	0.4864	0.025617	1.9549
0.5140	0.4883	0.025717	1.9626
0.5160	0.4902	0.025817	1.9702
0.5180	0.4921	0.025917	1.9778
0.5200	0.4940	0.026017	1.9855
0.5220	0.4959	0.026118	1.9931
0.5240	0.4978	0.026218	2.0007
0.5260	0.4997	0.026318	2.0083
0.5280	0.5016	0.026418	2.0160
0.5300	0.5035	0.026518	2.0236
0.5320	0.5054	0.026619	2.0312
0.5340	0.5073	0.026719	2.0389
0.5360	0.5092	0.026819	2.0465
0.5380	0.5111	0.026919	2.0541
0.5400	0.5130	0.027019	2.0618
0.5420	0.5149	0.027120	2.0694
0.5440	0.5168	0.027220	2.0770
0.5460	0.5187	0.027320	2.0846
0.5480	0.5206	0.027420	2.0923
0.5500	0.5225	0.027521	2.0999
0.5520	0.5244	0.027621	2.1075
0.5540	0.5263	0.027721	2.1152
0.5560	0.5282	0.027821	2.1228
0.5580	0.5301	0.027921	2.1304
0.5600	0.5320	0.028022	2.1380
0.5620	0.5339	0.028122	2.1457
0.5640	0.5358	0.028222	2.1533
0.5660	0.5377	0.028322	2.1609
0.5680	0.5396	0.028423	2.1686
0.5700	0.5415	0.028523	2.1762
0.5720	0.5434	0.028623	2.1838
0.5740	0.5453	0.028723	2.1914
0.5760	0.5472	0.028824	2.1991
0.5780	0.5491	0.028924	2.2067
0.5800	0.5510	0.029024	2.2143
0.5820	0.5529	0.029124	2.2220
0.5840	0.5548	0.029225	2.2296
0.5860	0.5567	0.029325	2.2372
0.5880	0.5586	0.029425	2.2448
0.5900	0.5605	0.029525	2.2525
0.5920	0.5624	0.029626	2.2601
0.5940	0.5643	0.029726	2.2677
0.5960	0.5662	0.029826	2.2754
0.5980	0.5681	0.029926	2.2830



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.6000	0.5700	0.030027	2.2906
0.6020	0.5719	0.030127	2.2982
0.6040	0.5738	0.030227	2.3059
0.6060	0.5757	0.030327	2.3135
0.6080	0.5776	0.030428	2.3211
0.6100	0.5795	0.030528	2.3287
0.6120	0.5814	0.030628	2.3364
0.6140	0.5833	0.030729	2.3440
0.6160	0.5852	0.030829	2.3516
0.6180	0.5871	0.030929	2.3593
0.6200	0.5890	0.031029	2.3669
0.6220	0.5909	0.031130	2.3745
0.6240	0.5928	0.031230	2.3821
0.6260	0.5947	0.031330	2.3898
0.6280	0.5966	0.031431	2.3974
0.6300	0.5985	0.031531	2.4050
0.6320	0.6004	0.031631	2.4126
0.6340	0.6023	0.031731	2.4203
0.6360	0.6042	0.031832	2.4279
0.6380	0.6061	0.031932	2.4355
0.6400	0.6080	0.032032	2.4431
0.6420	0.6099	0.032133	2.4508
0.6440	0.6118	0.032233	2.4584
0.6460	0.6137	0.032333	2.4660
0.6480	0.6156	0.032434	2.4736
0.6500	0.6175	0.032534	2.4813
0.6520	0.6194	0.032634	2.4889
0.6540	0.6213	0.032735	2.4965
0.6560	0.6232	0.032835	2.5041
0.6580	0.6251	0.032935	2.5118
0.6600	0.6270	0.033035	2.5194
0.6620	0.6289	0.033136	2.5270
0.6640	0.6308	0.033236	2.5346
0.6660	0.6327	0.033336	2.5423
0.6680	0.6346	0.033437	2.5499
0.6700	0.6365	0.033537	2.5575
0.6720	0.6384	0.033637	2.5651
0.6740	0.6403	0.033738	2.5728
0.6760	0.6422	0.033838	2.5804
0.6780	0.6441	0.033938	2.5880
0.6800	0.6460	0.034039	2.5956
0.6820	0.6479	0.034139	2.6033
0.6840	0.6498	0.034239	2.6109
0.6860	0.6517	0.034340	2.6185
0.6880	0.6536	0.034440	2.6261
0.6900	0.6555	0.034541	2.6337
0.6920	0.6574	0.034641	2.6414
0.6940	0.6593	0.034741	2.6490
0.6960	0.6612	0.034842	2.6566
0.6980	0.6631	0.034942	2.6642
0.7000	0.6650	0.035042	2.6719
0.7020	0.6669	0.035143	2.6795
0.7040	0.6688	0.035243	2.6871
0.7060	0.6707	0.035343	2.6947
0.7080	0.6726	0.035444	2.7024
0.7100	0.6745	0.035544	2.7100
0.7120	0.6764	0.035645	2.7176
0.7140	0.6783	0.035745	2.7252
0.7160	0.6802	0.035845	2.7328
0.7180	0.6821	0.035946	2.7405





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.7200	0.6840	0.036046	2.7481
0.7220	0.6859	0.036146	2.7557
0.7240	0.6878	0.036247	2.7633
0.7260	0.6897	0.036347	2.7710
0.7280	0.6916	0.036448	2.7786
0.7300	0.6935	0.036548	2.7862
0.7320	0.6954	0.036648	2.7938
0.7340	0.6973	0.036749	2.8014
0.7360	0.6992	0.036849	2.8091
0.7380	0.7011	0.036950	2.8167
0.7400	0.7029	0.037050	2.8243
0.7420	0.7048	0.037150	2.8319
0.7440	0.7067	0.037251	2.8395
0.7460	0.7086	0.037351	2.8472
0.7480	0.7105	0.037452	2.8548
0.7500	0.7124	0.037552	2.8624
0.7520	0.7143	0.037652	2.8700
0.7540	0.7162	0.037753	2.8776
0.7560	0.7181	0.037853	2.8853
0.7580	0.7200	0.037954	2.8929
0.7600	0.7219	0.038054	2.9005
0.7620	0.7238	0.038155	2.9081
0.7640	0.7257	0.038255	2.9157
0.7660	0.7276	0.038355	2.9234
0.7680	0.7295	0.038456	2.9310
0.7700	0.7314	0.038556	2.9386
0.7720	0.7333	0.038657	2.9462
0.7740	0.7352	0.038757	2.9538
0.7760	0.7371	0.038858	2.9615
0.7780	0.7390	0.038958	2.9691
0.7800	0.7409	0.039059	2.9767
0.7820	0.7428	0.039159	2.9843
0.7840	0.7447	0.039259	2.9919
0.7860	0.7466	0.039360	2.9996
0.7880	0.7485	0.039460	3.0072
0.7900	0.7504	0.039561	3.0148
0.7920	0.7523	0.039661	3.0224
0.7940	0.7542	0.039762	3.0300
0.7960	0.7561	0.039862	3.0376
0.7980	0.7580	0.039963	3.0453
0.8000	0.7599	0.040063	3.0529
0.8020	0.7618	0.040164	3.0605
0.8040	0.7637	0.040264	3.0681
0.8060	0.7656	0.040365	3.0757
0.8080	0.7675	0.040465	3.0834
0.8100	0.7694	0.040566	3.0910
0.8120	0.7713	0.040666	3.0986
0.8140	0.7732	0.040767	3.1062
0.8160	0.7751	0.040867	3.1138
0.8180	0.7770	0.040968	3.1214
0.8200	0.7789	0.041068	3.1291
0.8220	0.7808	0.041169	3.1367
0.8240	0.7827	0.041269	3.1443
0.8260	0.7846	0.041369	3.1519
0.8280	0.7865	0.041470	3.1595
0.8300	0.7884	0.041570	3.1671
0.8320	0.7903	0.041671	3.1748
0.8340	0.7922	0.041772	3.1824
0.8360	0.7941	0.041872	3.1900
0.8380	0.7960	0.041973	3.1976





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.8400	0.7979	0.042073	3.2052
0.8420	0.7998	0.042174	3.2128
0.8440	0.8017	0.042274	3.2204
0.8460	0.8036	0.042375	3.2281
0.8480	0.8055	0.042475	3.2357
0.8500	0.8074	0.042576	3.2433
0.8520	0.8093	0.042676	3.2509
0.8540	0.8112	0.042777	3.2585
0.8560	0.8131	0.042877	3.2661
0.8580	0.8150	0.042978	3.2738
0.8600	0.8169	0.043078	3.2814
0.8620	0.8188	0.043179	3.2890
0.8640	0.8207	0.043280	3.2966
0.8660	0.8226	0.043380	3.3042
0.8680	0.8245	0.043481	3.3118
0.8700	0.8264	0.043581	3.3194
0.8720	0.8283	0.043682	3.3270
0.8740	0.8302	0.043782	3.3347
0.8760	0.8321	0.043883	3.3423
0.8780	0.8340	0.043983	3.3499
0.8800	0.8359	0.044084	3.3575
0.8820	0.8378	0.044185	3.3651
0.8840	0.8397	0.044285	3.3727
0.8860	0.8416	0.044386	3.3803
0.8880	0.8435	0.044486	3.3880
0.8900	0.8454	0.044587	3.3956
0.8920	0.8473	0.044688	3.4032
0.8940	0.8492	0.044788	3.4108
0.8960	0.8511	0.044889	3.4184
0.8980	0.8530	0.044989	3.4260
0.9000	0.8549	0.045090	3.4336
0.9020	0.8568	0.045190	3.4412
0.9040	0.8587	0.045291	3.4489
0.9060	0.8606	0.045392	3.4565
0.9080	0.8625	0.045492	3.4641
0.9100	0.8644	0.045593	3.4717
0.9120	0.8663	0.045693	3.4793
0.9140	0.8682	0.045794	3.4869
0.9160	0.8701	0.045895	3.4945
0.9180	0.8720	0.045995	3.5021
0.9200	0.8739	0.046096	3.5097
0.9220	0.8758	0.046197	3.5174
0.9240	0.8777	0.046297	3.5250
0.9260	0.8796	0.046398	3.5326
0.9280	0.8815	0.046499	3.5402
0.9300	0.8834	0.046599	3.5478
0.9320	0.8853	0.046700	3.5554
0.9340	0.8872	0.046800	3.5630
0.9360	0.8891	0.046901	3.5706
0.9380	0.8910	0.047002	3.5782
0.9400	0.8929	0.047102	3.5858
0.9420	0.8948	0.047203	3.5935
0.9440	0.8967	0.047304	3.6011
0.9460	0.8986	0.047404	3.6087
0.9480	0.9005	0.047505	3.6163
0.9500	0.9024	0.047606	3.6239
0.9520	0.9043	0.047706	3.6315
0.9540	0.9062	0.047807	3.6391
0.9560	0.9081	0.047908	3.6467
0.9580	0.9100	0.048009	3.6543



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
0.9600	0.9119	0.048109	3.6619
0.9620	0.9138	0.048210	3.6695
0.9640	0.9157	0.048310	3.6771
0.9660	0.9176	0.048411	3.6848
0.9680	0.9195	0.048512	3.6924
0.9700	0.9214	0.048613	3.7000
0.9720	0.9233	0.048713	3.7076
0.9740	0.9252	0.048814	3.7152
0.9760	0.9271	0.048915	3.7228
0.9780	0.9290	0.049015	3.7304
0.9800	0.9309	0.049116	3.7380
0.9820	0.9328	0.049217	3.7456
0.9840	0.9347	0.049318	3.7532
0.9860	0.9366	0.049418	3.7608
0.9880	0.9385	0.049519	3.7684
0.9900	0.9404	0.049620	3.7760
0.9920	0.9423	0.049720	3.7836
0.9940	0.9442	0.049821	3.7913
0.9960	0.9461	0.049922	3.7989
0.9980	0.9480	0.050023	3.8065
1.0000	0.9499	0.050123	3.8141
1.0020	0.9518	0.050224	3.8217
1.0040	0.9537	0.050325	3.8293
1.0060	0.9556	0.050426	3.8369
1.0080	0.9575	0.050526	3.8445
1.0100	0.9594	0.050627	3.8521
1.0120	0.9613	0.050728	3.8597
1.0140	0.9632	0.050829	3.8673
1.0160	0.9651	0.050929	3.8749
1.0180	0.9670	0.051030	3.8825
1.0200	0.9689	0.051131	3.8901
1.0220	0.9708	0.051232	3.8977
1.0240	0.9727	0.051332	3.9053
1.0260	0.9746	0.051433	3.9129
1.0280	0.9765	0.051534	3.9205
1.0300	0.9784	0.051635	3.9281
1.0320	0.9803	0.051735	3.9357
1.0340	0.9822	0.051836	3.9433
1.0360	0.9841	0.051937	3.9509
1.0380	0.9860	0.052038	3.9586
1.0400	0.9879	0.052139	3.9662
1.0420	0.9898	0.052240	3.9738
1.0440	0.9917	0.052340	3.9814
1.0460	0.9936	0.052441	3.9890
1.0480	0.9955	0.052542	3.9966
1.0500	0.9974	0.052643	4.0042
1.0520	0.9993	0.052743	4.0118
1.0540	1.0012	0.052844	4.0194
1.0560	1.0031	0.052945	4.0270
1.0580	1.0050	0.053046	4.0346
1.0600	1.0069	0.053147	4.0422
1.0620	1.0088	0.053248	4.0498
1.0640	1.0107	0.053348	4.0574
1.0660	1.0125	0.053449	4.0650
1.0680	1.0144	0.053550	4.0726
1.0700	1.0163	0.053651	4.0802
1.0720	1.0182	0.053752	4.0878
1.0740	1.0201	0.053853	4.0954
1.0760	1.0220	0.053953	4.1030
1.0780	1.0239	0.054054	4.1106



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.0800	1.0258	0.054155	4.1182
1.0820	1.0277	0.054256	4.1258
1.0840	1.0296	0.054357	4.1334
1.0860	1.0315	0.054458	4.1410
1.0880	1.0334	0.054559	4.1486
1.0900	1.0353	0.054660	4.1562
1.0920	1.0372	0.054760	4.1638
1.0940	1.0391	0.054861	4.1714
1.0960	1.0410	0.054962	4.1790
1.0980	1.0429	0.055063	4.1866
1.1000	1.0448	0.055164	4.1942
1.1020	1.0467	0.055265	4.2018
1.1040	1.0486	0.055366	4.2094
1.1060	1.0505	0.055467	4.2170
1.1080	1.0524	0.055568	4.2246
1.1100	1.0543	0.055668	4.2322
1.1120	1.0562	0.055769	4.2398
1.1140	1.0581	0.055870	4.2474
1.1160	1.0600	0.055971	4.2550
1.1180	1.0619	0.056072	4.2626
1.1200	1.0638	0.056173	4.2702
1.1220	1.0657	0.056274	4.2778
1.1240	1.0676	0.056375	4.2854
1.1260	1.0695	0.056476	4.2929
1.1280	1.0714	0.056577	4.3005
1.1300	1.0733	0.056678	4.3081
1.1320	1.0752	0.056779	4.3157
1.1340	1.0771	0.056880	4.3233
1.1360	1.0790	0.056981	4.3309
1.1380	1.0809	0.057082	4.3385
1.1400	1.0828	0.057183	4.3461
1.1420	1.0847	0.057284	4.3537
1.1440	1.0866	0.057385	4.3613
1.1460	1.0885	0.057486	4.3689
1.1480	1.0904	0.057586	4.3765
1.1500	1.0923	0.057687	4.3841
1.1520	1.0942	0.057788	4.3917
1.1540	1.0961	0.057889	4.3993
1.1560	1.0980	0.057990	4.4069
1.1580	1.0999	0.058091	4.4145
1.1600	1.1018	0.058192	4.4221
1.1620	1.1037	0.058293	4.4297
1.1640	1.1056	0.058394	4.4373
1.1660	1.1075	0.058495	4.4448
1.1680	1.1094	0.058596	4.4524
1.1700	1.1113	0.058697	4.4600
1.1720	1.1132	0.058798	4.4676
1.1740	1.1151	0.058899	4.4752
1.1760	1.1170	0.059000	4.4828
1.1780	1.1189	0.059101	4.4904
1.1800	1.1208	0.059202	4.4980
1.1820	1.1227	0.059303	4.5056
1.1840	1.1246	0.059404	4.5132
1.1860	1.1265	0.059506	4.5208
1.1880	1.1284	0.059607	4.5284
1.1900	1.1303	0.059708	4.5360
1.1920	1.1322	0.059809	4.5436
1.1940	1.1341	0.059910	4.5511
1.1960	1.1360	0.060011	4.5587
1.1980	1.1379	0.060112	4.5663





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.2000	1.1398	0.060213	4.5739
1.2020	1.1417	0.060314	4.5815
1.2040	1.1436	0.060415	4.5891
1.2060	1.1455	0.060516	4.5967
1.2080	1.1474	0.060617	4.6043
1.2100	1.1493	0.060718	4.6119
1.2120	1.1512	0.060819	4.6195
1.2140	1.1531	0.060920	4.6270
1.2160	1.1550	0.061021	4.6346
1.2180	1.1569	0.061123	4.6422
1.2200	1.1588	0.061224	4.6498
1.2220	1.1607	0.061325	4.6574
1.2240	1.1626	0.061426	4.6650
1.2260	1.1645	0.061527	4.6726
1.2280	1.1664	0.061628	4.6802
1.2300	1.1683	0.061729	4.6878
1.2320	1.1702	0.061830	4.6954
1.2340	1.1721	0.061931	4.7029
1.2360	1.1740	0.062032	4.7105
1.2380	1.1759	0.062134	4.7181
1.2400	1.1778	0.062235	4.7257
1.2420	1.1797	0.062336	4.7333
1.2440	1.1816	0.062437	4.7409
1.2460	1.1835	0.062538	4.7485
1.2480	1.1854	0.062639	4.7560
1.2500	1.1873	0.062740	4.7636
1.2520	1.1892	0.062842	4.7712
1.2540	1.1911	0.062943	4.7788
1.2560	1.1930	0.063044	4.7864
1.2580	1.1949	0.063145	4.7940
1.2600	1.1968	0.063246	4.8016
1.2620	1.1987	0.063348	4.8092
1.2640	1.2006	0.063448	4.8167
1.2660	1.2024	0.063550	4.8243
1.2680	1.2043	0.063651	4.8319
1.2700	1.2062	0.063752	4.8395
1.2720	1.2081	0.063854	4.8471
1.2740	1.2100	0.063955	4.8547
1.2760	1.2119	0.064056	4.8622
1.2780	1.2138	0.064157	4.8698
1.2800	1.2157	0.064258	4.8774
1.2820	1.2176	0.064360	4.8850
1.2840	1.2195	0.064461	4.8926
1.2860	1.2214	0.064562	4.9002
1.2880	1.2233	0.064663	4.9078
1.2900	1.2252	0.064764	4.9153
1.2920	1.2271	0.064865	4.9229
1.2940	1.2290	0.064967	4.9305
1.2960	1.2309	0.065068	4.9381
1.2980	1.2328	0.065169	4.9457
1.3000	1.2347	0.065270	4.9533
1.3020	1.2366	0.065372	4.9608
1.3040	1.2385	0.065473	4.9684
1.3060	1.2404	0.065574	4.9760
1.3080	1.2423	0.065675	4.9836
1.3100	1.2442	0.065777	4.9912
1.3120	1.2461	0.065878	4.9987
1.3140	1.2480	0.065979	5.0063
1.3160	1.2499	0.066081	5.0139
1.3180	1.2518	0.066182	5.0215





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.3200	1.2537	0.066283	5.0291
1.3220	1.2556	0.066384	5.0366
1.3240	1.2575	0.066486	5.0442
1.3260	1.2594	0.066587	5.0518
1.3280	1.2613	0.066688	5.0594
1.3300	1.2632	0.066789	5.0670
1.3320	1.2651	0.066891	5.0745
1.3340	1.2670	0.066992	5.0821
1.3360	1.2689	0.067093	5.0897
1.3380	1.2708	0.067195	5.0973
1.3400	1.2727	0.067296	5.1049
1.3420	1.2746	0.067398	5.1124
1.3440	1.2765	0.067499	5.1200
1.3460	1.2784	0.067600	5.1276
1.3480	1.2803	0.067701	5.1352
1.3500	1.2822	0.067803	5.1428
1.3520	1.2841	0.067904	5.1503
1.3540	1.2860	0.068005	5.1579
1.3560	1.2879	0.068107	5.1655
1.3580	1.2898	0.068208	5.1731
1.3600	1.2917	0.068309	5.1806
1.3620	1.2936	0.068411	5.1882
1.3640	1.2955	0.068512	5.1958
1.3660	1.2974	0.068614	5.2034
1.3680	1.2993	0.068715	5.2110
1.3700	1.3012	0.068816	5.2185
1.3720	1.3031	0.068918	5.2261
1.3740	1.3050	0.069019	5.2337
1.3760	1.3069	0.069121	5.2413
1.3780	1.3088	0.069222	5.2488
1.3800	1.3107	0.069323	5.2564
1.3820	1.3126	0.069425	5.2640
1.3840	1.3145	0.069526	5.2716
1.3860	1.3164	0.069628	5.2791
1.3880	1.3183	0.069729	5.2867
1.3900	1.3202	0.069830	5.2943
1.3920	1.3221	0.069932	5.3019
1.3940	1.3240	0.070033	5.3094
1.3960	1.3259	0.070135	5.3170
1.3980	1.3278	0.070236	5.3246
1.4000	1.3297	0.070338	5.3322
1.4020	1.3316	0.070439	5.3397
1.4040	1.3335	0.070540	5.3473
1.4060	1.3354	0.070642	5.3549
1.4080	1.3373	0.070743	5.3624
1.4100	1.3392	0.070845	5.3700
1.4120	1.3411	0.070946	5.3776
1.4140	1.3430	0.071048	5.3852
1.4160	1.3448	0.071149	5.3927
1.4180	1.3467	0.071251	5.4003
1.4200	1.3486	0.071352	5.4079
1.4220	1.3505	0.071454	5.4154
1.4240	1.3524	0.071555	5.4230
1.4260	1.3543	0.071657	5.4306
1.4280	1.3562	0.071758	5.4382
1.4300	1.3581	0.071860	5.4457
1.4320	1.3600	0.071961	5.4533
1.4340	1.3619	0.072063	5.4609
1.4360	1.3638	0.072164	5.4684
1.4380	1.3657	0.072266	5.4760



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.4400	1.3676	0.072367	5.4836
1.4420	1.3695	0.072469	5.4912
1.4440	1.3714	0.072571	5.4987
1.4460	1.3733	0.072672	5.5063
1.4480	1.3752	0.072774	5.5139
1.4500	1.3771	0.072875	5.5214
1.4520	1.3790	0.072977	5.5290
1.4540	1.3809	0.073078	5.5366
1.4560	1.3828	0.073180	5.5441
1.4580	1.3847	0.073281	5.5517
1.4600	1.3866	0.073383	5.5593
1.4620	1.3885	0.073484	5.5668
1.4640	1.3904	0.073586	5.5744
1.4660	1.3923	0.073687	5.5820
1.4680	1.3942	0.073789	5.5895
1.4700	1.3961	0.073891	5.5971
1.4720	1.3980	0.073992	5.6047
1.4740	1.3999	0.074094	5.6122
1.4760	1.4018	0.074196	5.6198
1.4780	1.4037	0.074297	5.6274
1.4800	1.4056	0.074399	5.6349
1.4820	1.4075	0.074500	5.6425
1.4840	1.4094	0.074602	5.6501
1.4860	1.4113	0.074704	5.6576
1.4880	1.4132	0.074805	5.6652
1.4900	1.4151	0.074907	5.6728
1.4920	1.4170	0.075009	5.6803
1.4940	1.4189	0.075110	5.6879
1.4960	1.4208	0.075212	5.6955
1.4980	1.4227	0.075314	5.7030
1.5000	1.4246	0.075415	5.7106
1.5020	1.4265	0.075517	5.7182
1.5040	1.4284	0.075618	5.7257
1.5060	1.4303	0.075720	5.7333
1.5080	1.4322	0.075822	5.7408
1.5100	1.4341	0.075923	5.7484
1.5120	1.4360	0.076025	5.7560
1.5140	1.4379	0.076127	5.7635
1.5160	1.4398	0.076229	5.7711
1.5180	1.4417	0.076330	5.7787
1.5200	1.4436	0.076432	5.7862
1.5220	1.4455	0.076534	5.7938
1.5240	1.4474	0.076635	5.8013
1.5260	1.4493	0.076737	5.8089
1.5280	1.4512	0.076839	5.8165
1.5300	1.4531	0.076940	5.8240
1.5320	1.4550	0.077042	5.8316
1.5340	1.4569	0.077144	5.8391
1.5360	1.4588	0.077246	5.8467
1.5380	1.4607	0.077347	5.8543
1.5400	1.4625	0.077449	5.8618
1.5420	1.4644	0.077551	5.8694
1.5440	1.4663	0.077653	5.8769
1.5460	1.4682	0.077754	5.8845
1.5480	1.4701	0.077856	5.8921
1.5500	1.4720	0.077958	5.8996
1.5520	1.4739	0.078060	5.9072
1.5540	1.4758	0.078161	5.9147
1.5560	1.4777	0.078263	5.9223
1.5580	1.4796	0.078365	5.9299



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.5600	1.4815	0.078467	5.9374
1.5620	1.4834	0.078568	5.9450
1.5640	1.4853	0.078670	5.9525
1.5660	1.4872	0.078772	5.9601
1.5680	1.4891	0.078874	5.9676
1.5700	1.4910	0.078976	5.9752
1.5720	1.4929	0.079078	5.9827
1.5740	1.4948	0.079179	5.9903
1.5760	1.4967	0.079281	5.9979
1.5780	1.4986	0.079383	6.0054
1.5800	1.5005	0.079485	6.0130
1.5820	1.5024	0.079587	6.0205
1.5840	1.5043	0.079689	6.0281
1.5860	1.5062	0.079790	6.0356
1.5880	1.5081	0.079892	6.0432
1.5900	1.5100	0.079994	6.0508
1.5920	1.5119	0.080096	6.0583
1.5940	1.5138	0.080198	6.0659
1.5960	1.5157	0.080300	6.0734
1.5980	1.5176	0.080402	6.0810
1.6000	1.5195	0.080503	6.0885
1.6020	1.5214	0.080605	6.0961
1.6040	1.5233	0.080707	6.1036
1.6060	1.5252	0.080809	6.1112
1.6080	1.5271	0.080911	6.1187
1.6100	1.5290	0.081013	6.1263
1.6120	1.5309	0.081115	6.1338
1.6140	1.5328	0.081217	6.1414
1.6160	1.5347	0.081319	6.1489
1.6180	1.5366	0.081421	6.1565
1.6200	1.5385	0.081523	6.1640
1.6220	1.5404	0.081624	6.1716
1.6240	1.5423	0.081726	6.1791
1.6260	1.5442	0.081828	6.1867
1.6280	1.5461	0.081930	6.1942
1.6300	1.5480	0.082032	6.2018
1.6320	1.5499	0.082134	6.2093
1.6340	1.5518	0.082236	6.2169
1.6360	1.5537	0.082338	6.2244
1.6380	1.5556	0.082440	6.2320
1.6400	1.5575	0.082542	6.2395
1.6420	1.5594	0.082644	6.2471
1.6440	1.5613	0.082746	6.2546
1.6460	1.5632	0.082848	6.2622
1.6480	1.5650	0.082950	6.2697
1.6500	1.5669	0.083052	6.2773
1.6520	1.5688	0.083154	6.2848
1.6540	1.5707	0.083256	6.2924
1.6560	1.5726	0.083358	6.2999
1.6580	1.5745	0.083460	6.3075
1.6600	1.5764	0.083562	6.3150
1.6620	1.5783	0.083664	6.3226
1.6640	1.5802	0.083766	6.3301
1.6660	1.5821	0.083868	6.3377
1.6680	1.5840	0.083970	6.3452
1.6700	1.5859	0.084072	6.3528
1.6720	1.5878	0.084174	6.3603
1.6740	1.5897	0.084276	6.3679
1.6760	1.5916	0.084378	6.3754
1.6780	1.5935	0.084480	6.3829





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.6800	1.5954	0.084583	6.3905
1.6820	1.5973	0.084685	6.3980
1.6840	1.5992	0.084787	6.4056
1.6860	1.6011	0.084889	6.4131
1.6880	1.6030	0.084991	6.4207
1.6900	1.6049	0.085093	6.4282
1.6920	1.6068	0.085195	6.4358
1.6940	1.6087	0.085297	6.4433
1.6960	1.6106	0.085399	6.4508
1.6980	1.6125	0.085501	6.4584
1.7000	1.6144	0.085603	6.4659
1.7020	1.6163	0.085706	6.4735
1.7040	1.6182	0.085808	6.4810
1.7060	1.6201	0.085910	6.4886
1.7080	1.6220	0.086012	6.4961
1.7100	1.6239	0.086114	6.5036
1.7120	1.6258	0.086216	6.5112
1.7140	1.6277	0.086318	6.5187
1.7160	1.6296	0.086421	6.5263
1.7180	1.6315	0.086523	6.5338
1.7200	1.6334	0.086625	6.5413
1.7220	1.6353	0.086727	6.5489
1.7240	1.6372	0.086829	6.5564
1.7260	1.6391	0.086932	6.5640
1.7280	1.6410	0.087034	6.5715
1.7300	1.6429	0.087136	6.5790
1.7320	1.6448	0.087238	6.5866
1.7340	1.6467	0.087340	6.5941
1.7360	1.6486	0.087443	6.6017
1.7380	1.6505	0.087545	6.6092
1.7400	1.6524	0.087647	6.6167
1.7420	1.6543	0.087749	6.6243
1.7440	1.6561	0.087851	6.6318
1.7460	1.6580	0.087954	6.6393
1.7480	1.6599	0.088056	6.6469
1.7500	1.6618	0.088158	6.6544
1.7520	1.6637	0.088260	6.6620
1.7540	1.6656	0.088363	6.6695
1.7560	1.6675	0.088465	6.6770
1.7580	1.6694	0.088567	6.6846
1.7600	1.6713	0.088669	6.6921
1.7620	1.6732	0.088772	6.6996
1.7640	1.6751	0.088874	6.7072
1.7660	1.6770	0.088976	6.7147
1.7680	1.6789	0.089079	6.7222
1.7700	1.6808	0.089181	6.7298
1.7720	1.6827	0.089283	6.7373
1.7740	1.6846	0.089385	6.7448
1.7760	1.6865	0.089488	6.7524
1.7780	1.6884	0.089590	6.7599
1.7800	1.6903	0.089692	6.7674
1.7820	1.6922	0.089795	6.7750
1.7840	1.6941	0.089897	6.7825
1.7860	1.6960	0.089999	6.7900
1.7880	1.6979	0.090102	6.7976
1.7900	1.6998	0.090204	6.8051
1.7920	1.7017	0.090306	6.8126
1.7940	1.7036	0.090409	6.8202
1.7960	1.7055	0.090511	6.8277
1.7980	1.7074	0.090613	6.8352





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.8000	1.7093	0.090716	6.8428
1.8020	1.7112	0.090818	6.8503
1.8040	1.7131	0.090921	6.8578
1.8060	1.7150	0.091023	6.8654
1.8080	1.7169	0.091125	6.8729
1.8100	1.7188	0.091228	6.8804
1.8120	1.7207	0.091330	6.8879
1.8140	1.7226	0.091433	6.8955
1.8160	1.7245	0.091535	6.9030
1.8180	1.7264	0.091637	6.9105
1.8200	1.7283	0.091740	6.9181
1.8220	1.7302	0.091842	6.9256
1.8240	1.7321	0.091945	6.9331
1.8260	1.7340	0.092047	6.9406
1.8280	1.7358	0.092150	6.9482
1.8300	1.7377	0.092252	6.9557
1.8320	1.7396	0.092355	6.9632
1.8340	1.7415	0.092457	6.9708
1.8360	1.7434	0.092559	6.9783
1.8380	1.7453	0.092662	6.9858
1.8400	1.7472	0.092764	6.9933
1.8420	1.7491	0.092867	7.0009
1.8440	1.7510	0.092969	7.0084
1.8460	1.7529	0.093072	7.0159
1.8480	1.7548	0.093174	7.0234
1.8500	1.7567	0.093277	7.0310
1.8520	1.7586	0.093379	7.0385
1.8540	1.7605	0.093482	7.0460
1.8560	1.7624	0.093584	7.0535
1.8580	1.7643	0.093687	7.0611
1.8600	1.7662	0.093790	7.0686
1.8620	1.7681	0.093892	7.0761
1.8640	1.7700	0.093995	7.0836
1.8660	1.7719	0.094097	7.0912
1.8680	1.7738	0.094200	7.0987
1.8700	1.7757	0.094302	7.1062
1.8720	1.7776	0.094405	7.1137
1.8740	1.7795	0.094507	7.1213
1.8760	1.7814	0.094610	7.1288
1.8780	1.7833	0.094713	7.1363
1.8800	1.7852	0.094815	7.1438
1.8820	1.7871	0.094918	7.1513
1.8840	1.7890	0.095020	7.1589
1.8860	1.7909	0.095123	7.1664
1.8880	1.7928	0.095226	7.1739
1.8900	1.7947	0.095328	7.1814
1.8920	1.7966	0.095431	7.1889
1.8940	1.7985	0.095533	7.1965
1.8960	1.8004	0.095636	7.2040
1.8980	1.8023	0.095739	7.2115
1.9000	1.8042	0.095841	7.2190
1.9020	1.8061	0.095944	7.2265
1.9040	1.8080	0.096047	7.2340
1.9060	1.8098	0.096149	7.2416
1.9080	1.8117	0.096252	7.2491
1.9100	1.8136	0.096355	7.2566
1.9120	1.8155	0.096457	7.2641
1.9140	1.8174	0.096560	7.2716
1.9160	1.8193	0.096663	7.2792
1.9180	1.8212	0.096765	7.2867



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
1.9200	1.8231	0.096868	7.2942
1.9220	1.8250	0.096971	7.3017
1.9240	1.8269	0.097073	7.3092
1.9260	1.8288	0.097176	7.3167
1.9280	1.8307	0.097279	7.3242
1.9300	1.8326	0.097382	7.3318
1.9320	1.8345	0.097484	7.3393
1.9340	1.8364	0.097587	7.3468
1.9360	1.8383	0.097690	7.3543
1.9380	1.8402	0.097793	7.3618
1.9400	1.8421	0.097895	7.3693
1.9420	1.8440	0.097998	7.3769
1.9440	1.8459	0.098101	7.3844
1.9460	1.8478	0.098204	7.3919
1.9480	1.8497	0.098306	7.3994
1.9500	1.8516	0.098409	7.4069
1.9520	1.8535	0.098512	7.4144
1.9540	1.8554	0.098615	7.4219
1.9560	1.8573	0.098718	7.4294
1.9580	1.8592	0.098820	7.4370
1.9600	1.8611	0.098923	7.4445
1.9620	1.8630	0.099026	7.4520
1.9640	1.8649	0.099129	7.4595
1.9660	1.8668	0.099232	7.4670
1.9680	1.8687	0.099335	7.4745
1.9700	1.8706	0.099437	7.4820
1.9720	1.8725	0.099540	7.4895
1.9740	1.8744	0.099643	7.4970
1.9760	1.8763	0.099746	7.5045
1.9780	1.8782	0.099849	7.5121
1.9800	1.8800	0.099952	7.5196
1.9820	1.8819	0.100054	7.5271
1.9840	1.8838	0.100157	7.5346
1.9860	1.8857	0.100260	7.5421
1.9880	1.8876	0.100363	7.5496
1.9900	1.8895	0.100466	7.5571
1.9920	1.8914	0.100569	7.5646
1.9940	1.8933	0.100672	7.5721
1.9960	1.8952	0.100775	7.5796
1.9980	1.8971	0.100878	7.5871
2.0000	1.8990	0.100981	7.5946
2.0020	1.9009	0.101084	7.6021
2.0040	1.9028	0.101186	7.6096
2.0060	1.9047	0.101289	7.6172
2.0080	1.9066	0.101392	7.6247
2.0100	1.9085	0.101495	7.6322
2.0120	1.9104	0.101598	7.6397
2.0140	1.9123	0.101701	7.6472
2.0160	1.9142	0.101804	7.6547
2.0180	1.9161	0.101907	7.6622
2.0200	1.9180	0.102010	7.6697
2.0220	1.9199	0.102113	7.6772
2.0240	1.9218	0.102216	7.6847
2.0260	1.9237	0.102319	7.6922
2.0280	1.9256	0.102422	7.6997
2.0300	1.9275	0.102525	7.7072
2.0320	1.9294	0.102628	7.7147
2.0340	1.9313	0.102731	7.7222
2.0360	1.9332	0.102834	7.7297
2.0380	1.9351	0.102937	7.7372



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
2.0400	1.9370	0.103040	7.7447
2.0420	1.9389	0.103143	7.7522
2.0440	1.9408	0.103246	7.7597
2.0460	1.9426	0.103349	7.7672
2.0480	1.9445	0.103453	7.7747
2.0500	1.9464	0.103556	7.7822
2.0520	1.9483	0.103659	7.7897
2.0540	1.9502	0.103762	7.7972
2.0560	1.9521	0.103865	7.8047
2.0580	1.9540	0.103968	7.8122
2.0600	1.9559	0.104071	7.8197
2.0620	1.9578	0.104174	7.8272
2.0640	1.9597	0.104277	7.8347
2.0660	1.9616	0.104380	7.8422
2.0680	1.9635	0.104483	7.8497
2.0700	1.9654	0.104587	7.8572
2.0720	1.9673	0.104690	7.8647
2.0740	1.9692	0.104793	7.8722
2.0760	1.9711	0.104896	7.8797
2.0780	1.9730	0.104999	7.8872
2.0800	1.9749	0.105102	7.8947
2.0820	1.9768	0.105205	7.9022
2.0840	1.9787	0.105309	7.9097
2.0860	1.9806	0.105412	7.9172
2.0880	1.9825	0.105515	7.9246
2.0900	1.9844	0.105618	7.9321
2.0920	1.9863	0.105721	7.9396
2.0940	1.9882	0.105825	7.9471
2.0960	1.9901	0.105928	7.9546
2.0980	1.9920	0.106031	7.9621
2.1000	1.9939	0.106134	7.9696
2.1020	1.9958	0.106238	7.9771
2.1040	1.9977	0.106341	7.9846
2.1060	1.9996	0.106444	7.9921
2.1080	2.0015	0.106547	7.9996
2.1100	2.0033	0.106650	8.0071
2.1120	2.0052	0.106754	8.0146
2.1140	2.0071	0.106857	8.0220
2.1160	2.0090	0.106960	8.0295
2.1180	2.0109	0.107064	8.0370
2.1200	2.0128	0.107167	8.0445
2.1220	2.0147	0.107270	8.0520
2.1240	2.0166	0.107374	8.0595
2.1260	2.0185	0.107477	8.0670
2.1280	2.0204	0.107580	8.0745
2.1300	2.0223	0.107683	8.0820
2.1320	2.0242	0.107787	8.0894
2.1340	2.0261	0.107890	8.0969
2.1360	2.0280	0.107993	8.1044
2.1380	2.0299	0.108097	8.1119
2.1400	2.0318	0.108200	8.1194
2.1420	2.0337	0.108303	8.1269
2.1440	2.0356	0.108407	8.1344
2.1460	2.0375	0.108510	8.1419
2.1480	2.0394	0.108613	8.1493
2.1500	2.0413	0.108717	8.1568
2.1520	2.0432	0.108820	8.1643
2.1540	2.0451	0.108924	8.1718
2.1560	2.0470	0.109027	8.1793
2.1580	2.0489	0.109131	8.1868





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
2.1600	2.0508	0.109234	8.1943
2.1620	2.0527	0.109337	8.2017
2.1640	2.0546	0.109441	8.2092
2.1660	2.0565	0.109544	8.2167
2.1680	2.0584	0.109647	8.2242
2.1700	2.0602	0.109751	8.2317
2.1720	2.0621	0.109854	8.2392
2.1740	2.0640	0.109958	8.2466
2.1760	2.0659	0.110061	8.2541
2.1780	2.0678	0.110165	8.2616
2.1800	2.0697	0.110268	8.2691
2.1820	2.0716	0.110372	8.2766
2.1840	2.0735	0.110475	8.2840
2.1860	2.0754	0.110579	8.2915
2.1880	2.0773	0.110682	8.2990
2.1900	2.0792	0.110786	8.3065
2.1920	2.0811	0.110889	8.3140
2.1940	2.0830	0.110993	8.3214
2.1960	2.0849	0.111096	8.3289
2.1980	2.0868	0.111200	8.3364
2.2000	2.0887	0.111303	8.3439
2.2020	2.0906	0.111407	8.3514
2.2040	2.0925	0.111510	8.3588
2.2060	2.0944	0.111614	8.3663
2.2080	2.0963	0.111717	8.3738
2.2100	2.0982	0.111821	8.3813
2.2120	2.1001	0.111924	8.3887
2.2140	2.1020	0.112028	8.3962
2.2160	2.1039	0.112131	8.4037
2.2180	2.1058	0.112235	8.4112
2.2200	2.1077	0.112339	8.4187
2.2220	2.1096	0.112442	8.4261
2.2240	2.1115	0.112546	8.4336
2.2260	2.1133	0.112650	8.4411
2.2280	2.1152	0.112753	8.4486
2.2300	2.1171	0.112857	8.4560
2.2320	2.1190	0.112960	8.4635
2.2340	2.1209	0.113064	8.4710
2.2360	2.1228	0.113168	8.4784
2.2380	2.1247	0.113271	8.4859
2.2400	2.1266	0.113375	8.4934
2.2420	2.1285	0.113479	8.5009
2.2440	2.1304	0.113582	8.5083
2.2460	2.1323	0.113686	8.5158
2.2480	2.1342	0.113790	8.5233
2.2500	2.1361	0.113893	8.5308
2.2520	2.1380	0.113997	8.5382
2.2540	2.1399	0.114101	8.5457
2.2560	2.1418	0.114205	8.5532
2.2580	2.1437	0.114308	8.5606
2.2600	2.1456	0.114412	8.5681
2.2620	2.1475	0.114516	8.5756
2.2640	2.1494	0.114620	8.5830
2.2660	2.1513	0.114723	8.5905
2.2680	2.1532	0.114827	8.5980
2.2700	2.1551	0.114931	8.6055
2.2720	2.1570	0.115034	8.6129
2.2740	2.1589	0.115138	8.6204
2.2760	2.1608	0.115242	8.6279
2.2780	2.1627	0.115346	8.6353





OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
2.2800	2.1645	0.115450	8.6428
2.2820	2.1664	0.115553	8.6503
2.2840	2.1683	0.115657	8.6577
2.2860	2.1702	0.115761	8.6652
2.2880	2.1721	0.115865	8.6727
2.2900	2.1740	0.115969	8.6801
2.2920	2.1759	0.116072	8.6876
2.2940	2.1778	0.116176	8.6951
2.2960	2.1797	0.116280	8.7025
2.2980	2.1816	0.116384	8.7100
2.3000	2.1835	0.116488	8.7174
2.3020	2.1854	0.116592	8.7249
2.3040	2.1873	0.116695	8.7324
2.3060	2.1892	0.116799	8.7398
2.3080	2.1911	0.116903	8.7473
2.3100	2.1930	0.117007	8.7548
2.3120	2.1949	0.117111	8.7622
2.3140	2.1968	0.117215	8.7697
2.3160	2.1987	0.117319	8.7772
2.3180	2.2006	0.117423	8.7846
2.3200	2.2025	0.117527	8.7921
2.3220	2.2044	0.117631	8.7995
2.3240	2.2063	0.117735	8.8070
2.3260	2.2082	0.117839	8.8145
2.3280	2.2101	0.117942	8.8219
2.3300	2.2120	0.118046	8.8294
2.3320	2.2138	0.118150	8.8368
2.3340	2.2157	0.118254	8.8443
2.3360	2.2176	0.118358	8.8518
2.3380	2.2195	0.118462	8.8592
2.3400	2.2214	0.118566	8.8667
2.3420	2.2233	0.118670	8.8741
2.3440	2.2252	0.118774	8.8816
2.3460	2.2271	0.118878	8.8890
2.3480	2.2290	0.118982	8.8965
2.3500	2.2309	0.119086	8.9040
2.3520	2.2328	0.119190	8.9114
2.3540	2.2347	0.119294	8.9189
2.3560	2.2366	0.119398	8.9263
2.3580	2.2385	0.119502	8.9338
2.3600	2.2404	0.119606	8.9412
2.3620	2.2423	0.119711	8.9487
2.3640	2.2442	0.119815	8.9561
2.3660	2.2461	0.119919	8.9636
2.3680	2.2480	0.120023	8.9711
2.3700	2.2499	0.120127	8.9785
2.3720	2.2518	0.120231	8.9860
2.3740	2.2537	0.120335	8.9934
2.3760	2.2556	0.120439	9.0009
2.3780	2.2575	0.120543	9.0083
2.3800	2.2594	0.120647	9.0158
2.3820	2.2612	0.120751	9.0232
2.3840	2.2631	0.120856	9.0307
2.3860	2.2650	0.120960	9.0381
2.3880	2.2669	0.121064	9.0456
2.3900	2.2688	0.121168	9.0530
2.3920	2.2707	0.121272	9.0605
2.3940	2.2726	0.121376	9.0679
2.3960	2.2745	0.121480	9.0754
2.3980	2.2764	0.121585	9.0828



OBSERVED DISTANCE (INCHES)	TRUE DISTANCE (INCHES)	ERROR (INCHES)	RAY ANGLE (DEGREES)
2.4000	2.2783	0.121689	9.0903
2.4020	2.2802	0.121793	9.0977
2.4040	2.2821	0.121897	9.1052
2.4060	2.2840	0.122001	9.1126
2.4080	2.2859	0.122106	9.1201
2.4100	2.2878	0.122210	9.1275
2.4120	2.2897	0.122314	9.1350
2.4140	2.2916	0.122418	9.1424
2.4160	2.2935	0.122523	9.1498
2.4180	2.2954	0.122627	9.1573
2.4200	2.2973	0.122731	9.1647
2.4220	2.2992	0.122835	9.1722
2.4240	2.3011	0.122940	9.1796
2.4260	2.3030	0.123044	9.1871
2.4280	2.3049	0.123148	9.1945
2.4300	2.3067	0.123253	9.2020
2.4320	2.3086	0.123357	9.2094
2.4340	2.3105	0.123461	9.2168
2.4360	2.3124	0.123565	9.2243
2.4380	2.3143	0.123670	9.2317
2.4400	2.3162	0.123774	9.2392
2.4420	2.3181	0.123879	9.2466
2.4440	2.3200	0.123983	9.2541
2.4460	2.3219	0.124087	9.2615
2.4480	2.3238	0.124191	9.2689
2.4500	2.3257	0.124296	9.2764
2.4520	2.3276	0.124400	9.2838
2.4540	2.3295	0.124505	9.2913
2.4560	2.3314	0.124609	9.2987
2.4580	2.3333	0.124713	9.3061
2.4600	2.3352	0.124818	9.3136
2.4620	2.3371	0.124922	9.3210
2.4640	2.3390	0.125027	9.3285
2.4660	2.3409	0.125131	9.3359
2.4680	2.3428	0.125235	9.3433
2.4700	2.3447	0.125340	9.3508
2.4720	2.3466	0.125444	9.3582
2.4740	2.3484	0.125549	9.3657
2.4760	2.3503	0.125653	9.3731
2.4780	2.3522	0.125758	9.3805
2.4800	2.3541	0.125862	9.3880
2.4820	2.3560	0.125967	9.3954
2.4840	2.3579	0.126071	9.4028
2.4860	2.3598	0.126176	9.4103
2.4880	2.3617	0.126280	9.4177
2.4900	2.3636	0.126385	9.4251
2.4920	2.3655	0.126489	9.4326
2.4940	2.3674	0.126594	9.4400
2.4960	2.3693	0.126698	9.4474
2.4980	2.3712	0.126803	9.4549
2.5000	2.3731	0.126907	9.4623



## APPENDIX C

### APPLICATION OF COMPUTER PROGRAM "HOLOFER"

The computer program is an adptation of the inversion first proposed by C. D. Maldonado [ 9, 10, 11] and is designed to invert fringe numbers across a field to the density field. It can be operated in three different modes as described below:

#### (a) Mode 1

Mode 1 is utilized as a self-test of the computer program. It can either generate its own input density field using Subroutine FUNCT or read in a density field through Subroutine FREAD. The program then generates the fringe array and inverts the array back to the original density field. This mode was utilized in the present investigation to determine the value of the scale factor,  $\alpha$ , required to obtain the correct density across the fin.

#### (b) Mode 2

This mode reads in irregularly spaced fringe data and generates the fringe array at regular intervals across the field using Subroutine SHEET. By specifying NCODE = 1, the fringe array can be generated by one of the functions in Subroutine FUNCT. Mode 2 was not utilized.

#### (c) Mode 3

Mode 3 reads in the fringe data at regularly spaced intervals and inverts the array to density data across the field. The Subroutine GARRAY calls Subroutine READ to read in the fringe data. The first two cards preceeding the fringe data provide the program with the fringe field size, location, and symmetry.



The following parameters were used in considering the symmetric field case:

<u>PARAMETER</u>	<u>INPUT</u>
NOF	Run Number
IMAX	201
JMAX	1
ISYM	101
JSYM	1
IMS	2
JMS	100
Z	0.387
XO	0.0
YO	0.0
PHISYM	0.0

References [ 3 ] and [ 12 ] contain further details and applications of the computer program. A print-out of the program is included in the next few pages of this appendix.









```

NPTS=AR(15)
NLINS=AR(16)
SD=AR(18)
PHIZ=AR(19)
DELPHI=AR(20)
YPZERO=AR(21)
YPRNG=AR(22)
XPZERO=AR(23)
XPRNG=AR(24)
NGF=AR(25)
NAF=AR(26)
IPT=AR(27)
KPT=AR(28)
LPT=AR(29)
BND=AR(30)
A=AR(31)
B=AR(32)
C=AR(33)
D=AR(34)
E=AR(35)
P=AR(36)
S=AR(37)
T=AR(38)
U=AR(39)
V=AR(40)
W=AR(41)
Q=AR(42)
WRITE(6,90)
WRITE(6,98)
WRITE(6,91)
WRITE(6,92)
WRITE(6,93)
WRITE(6,94)
WRITE(6,95)
WRITE(6,96)
WRITE(6,97)
WRITE(6,98)
FORMAT(//5X,/,*,IMAX,0)
1, MEXTRA *,/3X,6F10.0)
1, FORMAT(/5X,/,*,ALPHA,3,F10.3,2F10.3,*,SIZE,*,EPS,*,RHO-INF,*,LAMBDA,*,SET00340
1, BETA,*,/3X,*,MODE,*,F10.3)
1, FORMAT(/5X,/,*,/2X,5F10.0,*,DELPHI,*,YPZERO,*,YPRANGE,*,XPZERO,*,SET00360
1, STD.DEV,*,/2X,*,PHIZERO,*,F10.3)
1, FORMAT(/5X,/,*,/4X,6F10.3)
1, XPRANGE *,/5X,/,*,/2X,6F10.0)
1, MAP BND *,/2X,6F10.0)

```

```

CAL00480
CAL00490
CAL00500
CAL00510
CAL00520
CAL00530
CAL00540
CAL00550
CAL00560
CAL00570
CAL00580
CAL00590
CAL00600
CAL00610
CAL00620
CAL00630
CAL00640
CAL00650
CAL00660
CAL00670
CAL00680
CAL00690
CAL00700
CAL00710
CAL00720
CAL00730
CAL00740

```

```

90
91
92
93
94
95

```



```

96  FORMAT (/5X,'* A * B * C * D * E *',SET00420
97  1,  FORMAT (/5X,'* S * T * U * V * W *',SET00440
98  1,  FORMAT (/3X,75A1)
    IF (DGN.GE.4) WRITE (6,89) (AR(I),I=1,42)
    NNN=2
    IF (MODE.LT.0) NNN=1
    IF (MODE.GT.5) NNN=3
    IF (MODE.GT.5) MODE=MODE-10
    NGP=0
    IF (KLIMIT.LT.KEXTRA) KEXTRA=KLIMIT
    IF (MLIMIT.LT.MEXTRA) MEXTRA=MLIMIT
    IF (IPT.LT.0) NGP=IPT
    IF (IPT.LT.0) IPT=-IPT
    ISYM=2.1-(FLOAT(JSYM)/2.-FLOAT(JSYM/2))*2
    IF (JSYM.EQ.0) ISYM=1
    IF (JSYM.GT.JMAX) ISYM=2
    IF (ISYM.EQ.1) JMAX=((JMAX+1)/2)*2
    RJMX=JMAX
    MSYM=JSYM
    IF ((MSYM.EQ.0).OR.(MSYM.GT.JMAX)) MSYM=1
    FCU=ISYM*JSYM*JMAX
    IF ((JSYM.GT.JMAX).OR.(JSYM.EQ.0)) FCU=JMAX
    QSYM=FCU/RJMX
    IMS=(IMAX+ISYM-1)/ISYM
    JMS=JMAX
    IF (ISYM.EQ.1) JMS=(JMAX/2+1)/2
    IF (JSYM.EQ.0) JMS=JMAX/2
    MODE=ABS(AR(13))
    XQ=0.
    YO=0.
    ZD=0.
    PHISYM=0.
    HS=SIZE/2.
    RHOS=1.286
    BOX=RHOFINF*BETA/RHOS/RLAMDA
    RPTS=NPTS
    XPR=0.
    IF (NPTS.GT.1) XPR=XPRNG/(RPTS-1.)/2.
    XPM=-XPR
    PIE=3.141592653589793
    MONE=1
    WRITE (6,58) IMAX,JMAX,IMS,JMS,ISYM,JSYM,MSYM,QSYM,FCU
    1  FORMAT (3X,'I MAX,IMS,JMS,ISYM,JSYM,MSYM,QSYM,FCU',/
      715,2F7.3/)
    NTWO=2
    IX=IMAX+JMAX

```



CAL01180  
CAL01190  
CAL01200  
CAL01210  
CAL01220  
CAL01230  
CAL01240  
CAL01250  
CAL01260  
CAL01270  
CAL01280  
CAL01290  
CAL01300  
CAL01310  
CAL01320  
CAL01330  
CAL01340  
CAL01350  
CAL01360  
CAL01370  
CAL01380  
CAL01390  
CAL01400  
CAL01410  
CAL01420  
CAL01430  
CAL01440  
CAL01450  
CAL01460  
CAL01470  
CAL01480  
CAL01490  
CAL01500  
CAL01510  
CAL01520  
CAL01530  
CAL01540  
CAL01550  
CAL01560  
CAL01570  
CAL01580  
CAL01590  
CAL01600  
CAL01610  
CAL01620  
CAL01630  
CAL01640  
CAL01650

```

NF=IN1
IF ((MODE.EQ.1.) .AND. (NOF.EQ.8) .AND. (DGN.GE.1.)) WRITE (6,69)
IF ((MODE.EQ.1.) .AND. (NOF.EQ.8)) CALL FREAD (NO,RO,NF,ZD)
Z=ZD
IF (DGN.GE.1.) WRITE (6,68)
CALL GARRAY (G,GA,NOF,DGN,MONE,XO,YO,PHISYM)
LM=1
IF ((LPT.EQ.0) .AND. (BND.EQ.0)) LM=0
IF IMX=IMAX+1
JMX=JMAX+1
IJMX=IMAX*JMAX
NBD=1
IF (JSYM.EQ.0) NBD=2
KBD=KLIMIT*NBD
DO 15 IJ=1, IJMX
GA(IJ)=0
IF (NAF.EQ.0) GO TO 16
NF=IN2
IF ((NAF.EQ.8) .AND. (DGN.GE.1.)) WRITE (6,69)
IF (NAF.EQ.8) CALL FREAD (NA,RA,NF,ZD)
MST=MODE
MODE=1
IF (DGN.GE.1.) WRITE (6,68)
IF (NAF.NE.0) CALL GARRAY (GA,G,NAF,DGN,NTWO,XO,YO,PHISYM)
MODE=MST
DO 6 IJ=1, IJMX
G(IJ)=G(IJ)+GA(IJ)
RLINS=NLINS
IF (NAF.EQ.8) WRITE (6,88) NA,(RA(L),L=1,NA)
IF (NOF.EQ.8) WRITE (6,87) NO,(RO(I),I=1,NO)
IF (LM.EQ.0) GO TO 14
RB(1)=-1.
DO 1 I=2,7
RB(I)=RB(I-1)+.5
TPIE=2.*PIE
MPIE=-PIE
DYP=0.
DXP=0.
IF (NLINS.GT.1) DYP=YPRNG/(RLINS-1.)
IF (NPPTS.GT.1) DXP=XPRNG/(RPTS-1.)
IF ((DGN.GE.1.) .AND. (NNN.EQ.2)) WRITE (6,64)
IF (NNN.EQ.2) CALL BDGEN (G,H,SCF,DGN,NBD,BDA,KBD)
DO 5 J=1,NLINS
IF (DGN.GE.1.) WRITE (6,67) J
RJM=J-1
PHI=PHI+DELPHI*RJM
YP(J)=YPZERO+DYP*RJM
PSI=(PHI+90.)*PIE/180.

```

15

6  
16

1





CAL01660  
CAL01670  
CAL01680  
CAL01690  
CAL01700  
CAL01710  
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CAL01800  
CAL01810  
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CAL01940  
CAL01950  
CAL01960  
CAL01970  
CAL01980  
CAL01990  
CAL02000  
CAL02010  
CAL02020  
CAL02030  
CAL02040  
CAL02050  
CAL02060  
CAL02070  
CAL02080  
CAL02090  
CAL02100  
CAL02110  
CAL02120  
CAL02130

```

TAU=PSI-PHISYM
IF (LPT.EQ.0) GO TO 9
IF (LPT.LE.1) WRITE (6,78) (ST,I=1,124)
IF (LPT.GT.1) WRITE (6,74) (ST,I=1,95)
IF ((CMS.EQ.1).AND.(LPT.GT.1)) READ (5,79) ZZ
WRITE (6,86)
WRITE (6,85) Z,PHI,YP(J)
WRITE (6,76)
IF (MODE.EQ.1) WRITE (6,83) (RB(I),I=1,7)
IF (MODE.GT.1) WRITE (6,80) (RB(I),I=1,7)
WRITE (6,81) (DH,I=1,54),(PL,I=1,13)
IC=0
DO 3 I=1,NPTS
  RIM=I-1
  THEO(I,J)=0.
  CA(I,J)=0.
  FA(I,J)=0.
  ERR(I)=0.
  CALC(I,J)=0.
  RHU(I)=0.
  XP(I)=XPZERO+DXP*RIM
  XPI=ABS(XP(I))
  IF (XPI.LT.1.E-10) XP(I)=0.
  RS=SQR(XP(I)**2+YP(J)**2)/HS
  IF (RS.GT.1.) GO TO 13
  THT=ATANM(YP(J),XP(I))
  IF (XPI.EQ.0.) THT=0.
  SIG=TAU-PIE/2.+THT
  IF (SIG.GT.PIE) SIG=SIG-TPIE
  IF (SIG.LT.MPIE) SIG=SIG+TPIE
  SIGI=SIG
  XS=RS*COS(SIG)
  IF (DGN.GE.1.) WRITE (6,44) SIG
  FORMAT (' SIG=',E10.3)
  YS=RS*SIN(SIG)
  IF (DGN.GE.5) WRITE (6,57) PHI,DELPHI,PSI,TAU,THT,SIG,SIGI,XS,YS
  IF FORMAT (' ANGLES=',10E10.3)
  RI=I
  F=0.
  IF (DGN.GE.2.) WRITE (6,66) I
  CALL FUNCT (XS,YS,FA(I,J),NAF,DGN,NTWO)
  IF (MODE.EQ.1) CALL FUNCT (XS,YS,F,NOF,DGN,MONE)
  THEO(I,J)=F
  IF (NNN.GE.2) REWIND 3
  IF (NNN.GE.2) CALL FIELD (RS,SIGI,SOLN,NBD,BDA,DGN,KBD)
  IF (NNN.EQ.1) CALL FIELD2 (RS,SIGI,SOLN,G,H,SCF,DGN)
  CA(I)=SOLN/BOX/HS
  CALC(I,J)=CA(I)-FA(I,J)

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CAL02140  
CAL02150  
CAL02160  
CAL02170  
CAL02180  
CAL02190  
CAL02200  
CAL02210

CAL02260  
CAL02270  
CAL02280  
CAL02290  
CAL02300  
CAL02310  
CAL02320  
CAL02330  
CAL02340  
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CAL02380  
CAL02390  
CAL02400  
CAL02410  
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CAL02470  
CAL02480  
CAL02490  
CAL02500  
CAL02510  
CAL02520  
CAL02530  
CAL02540  
CAL02550  
CAL02560  
CAL02570  
CAL02580  
CAL02590  
CAL02600  
CAL02610

13 RHO(I)=RHOINF\*(CALC(I,J)+1.)  
ERR(I)=CA(I)  
IF (MODE.EQ.1) ERR(I)=(CALC(I,J)-THEO(I,J))  
IF (MODE.GT.1) THEO(I,J)=FA(I,J)  
IF (LPT.EQ.0) GO TO 3  
LC=0  
TL(I)=BL  
TTL=0  
IF ((XP(I).GT.XPM).AND.(XP(I).LT.XPR)) TTL=1.  
IF (IC.EQ.5) IC=0  
IF (IC.EQ.0) TL(I)=PL  
DO 2 L=2,62  
TL(L)=BL  
IF ((I.EQ.1).OR.(TTL.EQ.1).OR.(I.EQ.NPTS)) TL(L)=PL  
IF (LC.EQ.10) LC=0  
IF ((IC.EQ.0).AND.(LC.EQ.0)) TL(L)=PL  
LC=LC+1  
TL(2)=PL  
TL(22)=PL  
TL(62)=PL  
IC=IC+1  
RLW=(CA(I)+1.)\*20.+2.5  
LW=RLW  
IF (LW.GT.62) LW=62  
IF (LW.LT.2) LW=2  
TL(LW)=SC  
RLY=(FA(I,J)+1.)\*20.+2.5  
LY=RLY  
IF (LY.GT.62) LY=62  
IF (LY.LT.2) LY=2  
IF (NAF.NE.0) TL(LY)=ST  
RLX=(THEO(I,J)+1.)\*20.+2.5  
LX=RLX  
IF (LX.GT.62) LX=62  
IF (LX.LT.2) LX=2  
IF (MODE.EQ.1) TL(LX)=OH  
RLZ=(CALC(I,J)+1.)\*20.+2.5  
LZ=RLZ  
IF (LZ.GT.62) LZ=62  
IF (LZ.LT.2) LZ=2  
TL(LZ)=EX  
WRITE (6,82) MOUT,KOUT,INDEX,THEO(I,J),ERR(I),CALC(I,J),RHO(I),  
1 XPR(I),TL(L),L=1,62)  
IF ((NPTS.LE.20).AND.(I.NE.NPTS)) WRITE (6,79)  
CONTINUE  
IF (LPT.NE.0) WRITE (6,81) (DH,I=1,54),(PL,I=1,13)  
IMAX=0.  
TMIN=0.

2

3



CAL02620  
 CAL02630  
 CAL02640  
 CAL02650  
 CAL02660  
 CAL02670  
 CAL02680  
 CAL02690  
 CAL02700  
 CAL02710  
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 CAL02730  
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 CAL02960  
 CAL02970  
 CAL02980  
 CAL02990  
 CAL03000  
 CAL03010  
 CAL03020  
 CAL03030  
 CAL03040  
 CAL03050  
 CAL03060  
 CAL03070  
 CAL03080  
 CAL03090

```

IE=0
BE=0.
DO 4 I=1,NPTS
  TH=THEO(I,J) TMAX=TH
  IF (TH.GT.TMIN) TMIN=TH
  ER=ABS(CALC(I,J)-TH)
  IF (ER.LE.BE) GO TO 4
BE=ER
IE=I
CONTINUE -TMIN
TMM=TMXIN*(CALC(IE,J)-THEO(IE,J))*100./TMM
EB=RHOIN*(CALC(IE,J)-THEO(IE,J)) WRITE (6,75) EB,XP(IE),BE
IF (TMM.NE.0.) BE=(CALC(IE,J)-THEO(IE,J)) WRITE (6,75) EB,XP(IE),BE
IF ((MODE.EQ.1).AND.(LPT.NE.0)) WRITE (6,75) EB,XP(IE),BE
IF (DELPHI.NE.0.) YP(J)=PHI
CONTINUE
IF (BND.EQ.0.) GO TO 14
IF (LPT.EQ.1) WRITE (6,78) (ST,I=1,124)
IF (LPT.GT.1) WRITE (6,74) (ST,I=1,95)
IF ((CMS.EQ.1).AND.(LPT.GT.1)) READ (5,79) ZZ
IF (DGN.GE.1) WRITE (6,63)
CALL MAP(NPTS,NLINS,CALC,NOF,Z,BND)
IF (NAF.EQ.0) GO TO 10
NAO=10*NOF+NAF
IF ((DGN.GE.1).AND.(NGP.EQ.-3)) WRITE (6,62)
IF (NGP.EQ.-3) CALL GPUNCH (Z,XO,YO,PHISYM,NAO,IMAX,JMAX,G)
DO 7 IJ=1,IJMX
  G(IJ)=G(IJ)-GA(IJ)
  IF (IPT.LE.0) GO TO 11
  IF ((IPT.EQ.1).OR.(IPT.EQ.3)) WRITE (6,78) (ST,I=1,124)
  IF ((IPT.EQ.2).OR.(IPT.GE.4)) WRITE (6,74) (ST,I=1,95)
  IF ((CMS.EQ.1).AND.(IPT.EQ.2).OR.(IPT.GE.4)) READ (5,79) ZZ
  CALL GPRINT (G,MONE)
  IF (NGP.EQ.-1) CALL GPUNCH (Z,XO,YO,PHISYM,NOF,IMAX,JMAX,G)
  IF (IPT.EQ.3) WRITE (6,78) (ST,I=1,124)
  IF (IPT.GE.4) WRITE (6,74) (ST,I=1,95)
  IF ((CMS.EQ.1).AND.(IPT.GE.4)) READ (5,79) ZZ
  IF (IPT.GE.3) CALL GPRINT (GA,NTWO)
  IF (KPT.LE.0) GO TO 12
  IF ((KPT.EQ.1).OR.(KPT.EQ.3)) WRITE (6,78) (ST,I=1,124)
  IF ((KPT.EQ.2).OR.(KPT.GE.4)) WRITE (6,74) (ST,I=1,95)
  IF ((CMS.EQ.1).AND.(KPT.EQ.2).OR.(KPT.GE.4)) READ (5,79) ZZ
  IF (DGN.GE.1) WRITE (6,61)
  CALL GPLOTT (G,GA,JMS)
  WRITE (6,78) (EX,I=1,124)
  AGAIN=ST
  12
  11
  10
  7
  14
  5
  4

```



```

89 IF (CMS.NE.1.) READ(5,60) AGAIN
88 IF (AGAIN.EQ.BL) GO TO 20
    WRITE (6,77)
    FORMAT (6F12.7)
87 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
86 1, POINTS) WAS: '7(11F10.3//)'
    THE INVERTED CROSS SECTION FOR: ' )
85 14HY' =,F8.3,' CM.,44X,'O = ORIGINAL FUNCTION'//
    FHY' =,F8.3,' CM.,44X,'O = ORIGINAL FUNCTION'//
27CX,'* = ADD-ON FUNCTION' )
84 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
83 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
82 1, POINTS) WAS: '7(11F10.3//)'
    THE INVERTED CROSS SECTION FOR: ' )
81 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
80 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
79 1, POINTS) WAS: '7(11F10.3//)'
    THE INVERTED CROSS SECTION FOR: ' )
78 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
77 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
76 1, POINTS) WAS: '7(11F10.3//)'
    THE INVERTED CROSS SECTION FOR: ' )
75 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
74 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
69 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
68 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
67 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
66 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
65 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
64 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
63 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
62 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
61 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR ADD-ON FUNCTION NO.8 (' ,I3,
60 1, POINTS) WAS: '7(11F10.3//)'
    THE INPUT DATA FOR THE TEST FUNCTION NO.8, (' ,I3,
    STOP
    END
C0C00001
C

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C          SUBROUTINE BDGEN (G,H,SCF,DGN,NBD,BDA,KBD)
C          BDGEN EVALUATES THE B AND D COEFFICIENTS FOR ALL M AND K, AND WRITES
C          THE ARRAY LINEARLY ON DISK.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /TAB/ INDEX,KEXTRA,MEXTRA,KLIMIT,MLIMIT,KOUT,MOUT
COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
DIMENSION G(IJMX),H(IIMX,5),SCF(JJMX,6),BDA(KBD)
INITIALIZE THE VALUES:
INDEX=0
KL2=NBD*KLIMIT
REWIND 3
JJMX6=JJMX*6
IIMX2=(IIMX+1)/2
PIE=3.141592653589793
RIMAX=IMAX
KLMP=KLIMIT+1
DX=2./RIMAX
RJMAX=JMAX
DXI=2.*PIE/FCU
INITIALIZE THE MODIFIED HERMITE POLYNOMIAL ARRAY:  VECTORS:
(1)=H1, (2)=H2, (3)=ALPHA*X(I), (4)=HM+2 STORED, (5)=HM+1 STORED
DO 1 I=1,IIMX2
RII=IIMX-I+1
IIM=IIMX-I+1
H(II,3)=ALPHA*(RII*DX-DX-1.)
H(II,3)=-H(II,3)
H(II,1)=2.*H(II,3)
H(II,2)=(H(II,3)*H(II,1)-1.)/3.
H(II,1)=-H(II,1)
H(II,2)=H(II,2)
H(II,5)=H(II,2)
H(II,4)=H(II,2)
H(II,5)=H(II,2)
H(II,4)=H(II,1)
H(II,4)=H(II,1)
SIGN=1.
INITIALIZE THE SIN/COS ARRAY:
DO 2 J=1,JJMX
RJM=J-1
SCF(J,1)=0.
SCF(J,2)=1.
SCF(J,3)=SIN(RJM*DXI-PIE/2.)
SCF(J,4)=COS(RJM*DXI-PIE/2.)
SCF(J,5)=0.
SCF(J,6)=0.
MS=0
C          COMMENCE THE M LOOP:

```



```

DO 7 MP=1,MLIMIT
  RM=MP-1
  RM=M
  SIGN=-SIGN
  IF (DGN.LE.-4) WRITE (6,88) SCF(1,1),SCF(2,1),SCF(1,2),SCF(2,2)
  TEST FOR SYMMETRY SKIPS:
  IF (MS.EQ.MSYM) MS=0
  TOTAL=0.
  MS=MS+1
  IF (MS.NE.1) GO TO 6
  COMMENCE THE K LOOP:
  DO 5 KP=1,KLIMIT
    K=KP-1
    PK=KP
    RK=K
    INDEX=INDEX+1
    CALL THE B & D COEFFICIENTS AND WRITE THEM ON DISK:
    CALL BD (M,K,G,H,SCF,B,D,JJMX6)
    IF (DGN.EQ.3.) WRITE (6,89) M,K,B,D
    IF (DGN.LE.-2) WRITE (6,89) M,K,B,D
    IF (DGN.LE.-4) WRITE (6,88) H(1,1),H(1,2),H(1,4),H(1,5)
    KK=K*NBD+1
    K2=KP*NBD
    BDA(K2)=D
    BDA(KK)=B
  GENERATE THE NEXT ORDER OF THE SET OF HERMITE POLYNOMIALS FOR NEW K:
  ORDER=M+2*KP+1
  HA=SQRT(PK*(PK+RM))/ORDER
  HB=2.*SQRT((PK+1.)*(RM+PK+1.))/(ORDER+1.)/(ORDER+2.)
  DO 5 II=1,IIMX2
    IIM=IIMX-II+1
    H(II,1)=2.*(H(II,3)*H(II,2)-HA*H(II,1))
    H(II,1)=SIGN*H(II,1)
    H(II,2)=HB*(H(II,3)*H(II,1)-ORDER*H(II,2))
    ADVANCE THE SIN/COS ARRAY FOR THE NEXT M:
    DO 3 J=1,JJMX
      IF (DGN.LE.-5) WRITE (6,87) (SCF(J,NT),NT=1,6)
      FORM=SCF(J,1)
      STEMP=SCF(J,1)
      SCF(J,1)=SCF(J,1)*SCF(J,4)+SCF(J,2)*SCF(J,3)
      SCF(J,2)=SCF(J,2)*SCF(J,4)-STEMP*SCF(J,3)
      DO 4 J=1,JJMAX
        SCF(J,5)=SCF(J+1,1)-SCF(J,1)
        SCF(J,6)=SCF(J+1,2)-SCF(J,2)
        WRITE (3) (BDA(I),I=1,KBD)
        IF (DGN.LE.-3) WRITE (6,88) (BDA(I),I=1,10)
        IF (JSYM.GT.JMAX) RETURN
      RM=RM+1

```



```

C      REGENERATE THE HERMITE ARRAY FOR NEW M, K=0:
DO 7 II=1,IIMX2
  IIM=IIMX-II+1
  H(II,2)=H(II,5)*(RM+1.)
  H(II,1)=H(II,5)*(RM+2.)
  H(IIM,1)=-SIGN*H(II,1)
  H(II,2)=2.*SQRTRM+1)*H(II,3)*H(II,1)-(RM+1.)*H(II,2))
  H(II,2)=H(II,2))/(RM+2.))/(RM+3.)
  H(II,4)=H(II,1)
  H(II,5)=H(II,2)
  H(II,5)=H(II,2)
  FORMAT ('M=',I4,'', K=',I4,'', B=',E10.4,'', D=',E10.4)
89  FORMAT (2X,10E10.3)
88  RETURN
C000002
C

```

```

SUBROUTINE FIELD (RS,SIG,SOLN,NBD,BDA,DGN,KBD)
C
C      FIELD EVALUATES THE VALUE OF THE FIELD FUNCTION AT A PARTICULAR
C      POINT DESIGNATED IN CYLINDRICAL COORDINATES, BY USING THE INVERSION
C      EQUATION OF MALDONADO, ET.AL. FIELD USES THE ARRAY OF B & D
C      COEFFICIENTS GENERATED IN SUBROUTINE BDGEN.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /TAB/ INDEX,KEXTRA,MEXTRA,KLIMIT,MLIMIT,KOUT,MOUT
COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
DIMENSION BDA(KBD),STK(52),STM(52)
C      INITIALIZE THE VALUES:
INDEX=0
ITIMER=0
KOUT=0
MOUT=0
MMAX=0
KMAX=0
TOTAL=0.
JJMX6=JJMX*6
REWIND 3
AR=ALPHA*RS
ARG=AR**2
EXPON=EXP(-ARG)
PIE=3.141592653589793
APP=ALPHA/PIE/PIE
M=0
RM=M
RIMAX=IMAX
DX=2./RIMAX
C
SUB01140
SUB01150
SUB01160
SUB01170
SUB01180
SUB01190
SUB01200
SUB01210
SUB01220
SUB01230
SUB01240
SUB01250
SUB01260
SUB01270
SUB01280
SUB01290
SUB01300
SUB01310
SUB01320
SUB01330
SUB01340
SUB01350
SUB01360
SUB01370
SUB01380
SUB01390
SUB01400
SUB01410
SUB01420
SUB01430
SUB00980
SUB00990
SUB01000
SUB01010
SUB01020
SUB01030
SUB01040
SUB01050
SUB01060
SUB01070
SUB01080
SUB01090
SUB01100
SUB01110
SUB01120
SUB01130

```



```

16 RJMAX=JMAX
   SIGN=1.
   STK(1)=0.
   STM(1)=0.
   SMS=0.
   CMS=1.
   SMI=1.
   CMI=COS(SIG)
   MEP=MEXTRA+1
   DO 16 MB=1,MEP
     STM(MB)=0.
     FM=1.
     MS=0
2   COMMENCE THE M LOOP:
     SIGN=-SIGN
     K=0
     RK=K
     RM=M
     ARM=1.
     IF (M.NE.0) ARM=AR**M
     KTIMER=0
     KEP=KEXTRA+1
     DO 15 KB=1,KEP
       STK(KB)=0.
       SIGNK=-1.
15  COMPUTE THE K=0 & K=1 ORDERS OF LAGUERRE POLYNOMIAL FOR GIVEN M:
       PM=0.
       P=SQRT(1./FM)
       PP=(RM+1.-ARG)*SQRT(1./FM/(RM+1.))
       TEST FOR SYMMETRY SKIPS:
       IF (MS.EQ.MSYM) MS=0
       MS=MS+1
       IF (MS.NE.1) GO TO 7
       READ A LINE OF B & D COEFFICIENTS FOR GIVEN M:
       READ (3) (BDA(I),I=1,KBD)
       IF (DGN.LE.-6) WRITE (6,88) (BDA(I),I=1,10)
       COMMENCE THE K LOOP:
       INDEX=INDEX+1
       SIGNK=-SIGNK
       COMPUTE THE M,K SUMMATION TERM:
       KK=K*NBD+1
       B=BDA(KK)
       D=0.
       IF (NBD.EQ.2) D=BDA(KK+1)
       BRAKET=B
       IF (RM.EQ.0.) GO TO 4
       BRAKET=B+CMS+D*SMS
       ADD=SIGNK*BRAKET*P*ARM
4

```

SUB01440  
 SUB01450  
 SUB01460  
 SUB01470  
 SUB01480  
 SUB01490  
 SUB01500  
 SUB01510  
 SUB01520  
 SUB01530  
 SUB01540  
 SUB01550  
 SUB01560  
 SUB01570  
 SUB01580  
 SUB01590  
 SUB01600  
 SUB01610  
 SUB01620  
 SUB01630  
 SUB01640  
 SUB01650  
 SUB01660  
 SUB01670  
 SUB01680  
 SUB01690  
 SUB01700  
 SUB01710  
 SUB01720  
 SUB01730  
 SUB01740  
 SUB01750  
 SUB01760  
 SUB01770  
 SUB01780  
 SUB01790  
 SUB01800  
 SUB01810  
 SUB01820  
 SUB01830  
 SUB01840  
 SUB01850  
 SUB01860  
 SUB01870  
 SUB01880  
 SUB01890  
 SUB01900  
 SUB01910





```

TOTAL=TOTAL+ADD
IF (DGN.GT.-5) GO TO 5
STOT=TOTAL*EXPON#APP/BOX/SIZE
WRITE (6,89) M,K,STOT,ADD,BRAKET,P,ARM,B,CMS,D,SMS
ESTABLISH CHECK AS THE RELATIVE SIZE OF THE M,K TERM OF THE SERIES:
5 CHECK=ABS(ADD)
IF (TOTAL.GT.EPS) CHECK=ABS(ADD/TOTAL)
C ADVANCE THE K INDEX:
K=K+1
RK=K
DO 10 KA=1,KEXTRA
KB=KEXTRA-KA+1
STK(KB+1)=STK(KB)
STK(2)=TOTAL
ORDER=M+2*K+1
C GENERATE THE NEXT ORDER OF LAGUERRE POLYNOMIAL FOR NEW K:
PM=PP
P=PP
PP=PP*(ORDER-ARG)-PM*SQRT(RK*(RM+RK))
PP=PP/SQRT((RK+1.)*(RM+RK+1.))
C SET K TIMER TO PROVIDE EXTRA K TERMS AFTER CHECK < EPS:
KTIMER=KTIMER+1
IF (K.GE.KLIMIT) GO TO 6
IF (CHECK.GE.EPS) KTIMER=0
IF (KTIMER.LE.KEXTRA) GO TO 3
GO TO 7
6 KOUT=KOUT+1
IF (KEXTRA.EQ.0) GO TO 7
TOTAL=0.
DO 11 KA=1,KEXTRA
TOTAL=TOTAL+STK(KA+1)
RKX=KEXTRA
TOTAL=TOTAL/RKX
C END OF K LOOP: ADVANCE M:
M=M+1
RM=M
STP=SMS
SMS=SMS*CMI+CMS*SMI
SMS=SMS*CMI-STP*SMI
CMS=CMS*CMI-KMAX
IF (K.GT.KMAX) KMAX=K
FM=FM*RM
DO 12 MA=1,MEXTRA
MB=MEXTRA-MA+1
STM(MB+1)=STM(MB)
STM(2)=TOTAL
C SET M TIMER FOR EXTRA M TERMS:
MTIMER=MTIMER+1
IF (JSYM.GT.JMAX) GO TO 9

```

SUB01920  
SUB01930  
SUB01940  
SUB01950  
SUB01960  
SUB01970  
SUB01980  
SUB01990  
SUB02000  
SUB02010  
SUB02020  
SUB02030  
SUB02040  
SUB02050  
SUB02060  
SUB02070  
SUB02080  
SUB02090  
SUB02100  
SUB02110  
SUB02120  
SUB02130  
SUB02140  
SUB02150  
SUB02160  
SUB02170  
SUB02180  
SUB02190  
SUB02200  
SUB02210  
SUB02220  
SUB02230  
SUB02240  
SUB02250  
SUB02260  
SUB02270  
SUB02280  
SUB02290  
SUB02300  
SUB02310  
SUB02320  
SUB02330  
SUB02340  
SUB02350  
SUB02360  
SUB02370  
SUB02380  
SUB02390



```

13 IF (K.GT.KEXTRA) MTIMER=0
   IF (M.GE.MLIMIT) GO TO 13
   IF (MTIMER.LE.MEXTRA) GO TO 2
   IF (MEXTRA.EQ.0) GO TO 9
   TOTAL=0.
   DO 14 MA=1,MEXTRA
   TOTAL=TOTAL+STM(MA+1)
14 RMX=MEXTRA
   TOTAL=TOTAL/RMX
   END OF M LOOP: COMPUTE OUTPUT SOLN.
9 MOUT=M-1
  IF (KOUT.EQ.0) KOUT=KMAX-1
  SOLN=TOTAL*EXPON*APP/2.
89 FORMAT (' M=',I4,' K=',I4,' SUBTOTAL=',9E10.3)
88 RETURN
   END
C000003
C

```

```

C
C SUBROUTINE BD (M,K,G,H,SCF,B,D,JJMX6)
C
C BD EVALUATES THE FIRST (B) AND SECOND (D) COEFFICIENTS IN THE
C INVERSION EQUATION. FOR A PARTICULAR SET OF INDEXES M & K.
C BD MAKES USE OF THE HERMITE POLYNOMIAL ARRAY GENERATED BY
C SUBROUTINE FIELD AS M & K ADVANCE.
C

```

```

COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
DIMENSION G(IJMX),SCF(JJMX6),H(IIMX)
PIE=3.141592653589793
B=0.
D=0.
RM=M
RK=K
RJMAXX=JMAX
JJMX4=4*JJMX
DXI=2.*PIE/FCU
FORMAT(1X,I10,/)
200 IF (JSYM.LE.0) GO TO 4
   IF (M.NE.0) GO TO 2
   S=DXI
   DO 1 J=1,JMAX
   DO 1 I=1,IMAX
   II=I+1
   IJ=IMAX*(J-1)+I

```

SUB02400  
SUB02410  
SUB02420  
SUB02430  
SUB02440  
SUB02450  
SUB02460  
SUB02470  
SUB02480  
SUB02490  
SUB02500  
SUB02510  
SUB02520  
SUB02530  
SUB02540  
SUB02550  
SUB02560  
SUB02570  
SUB02580

SUB02590  
SUB02600  
SUB02610  
SUB02620  
SUB02630  
SUB02640  
SUB02650  
SUB02660  
SUB02670  
SUB02680  
SUB02690  
SUB02700  
SUB02710  
SUB02720  
SUB02730  
SUB02740  
SUB02750  
SUB02760  
SUB02770

SUB02780  
SUB02790  
SUB02800  
SUB02810  
SUB02820  
SUB02830  
SUB02840



```

1      DH=H(I I)-H(I)
      B=B+G(I J)*S*DH
      B=B*QSYM/2.
      RETURN
2      DO 3 J=1,JMAX
      JS=J+J J M X 4
      S=SCF(JS)/RM
      DO 3 I=1,I MAX
      II=I+1
      IX=I MAX*(J-1)+I
      IJ=I MAX*(J-1)+I
      DH=H(I I)-H(I)
      B=B+G(I J)*S*DH
      B=B*QSYM
      RETURN
3      IF (M.NE.O) GO TO 6
      S=DX I
      DO 5 J=1,JMAX
      DO 5 I=1,I MAX
      II=I+1
      IX=I MAX*(J-1)+I
      IJ=I MAX*(J-1)+I
      DH=H(I I)-H(I)
      B=B+G(I J)*S*DH
      B=B/2.
      RETURN
4      DO 7 J=1,JMAX
      JS=J+J J M X 4
      J2=JS+J J M X
      S=SCF(JS)/RM
      C=SCF(J2)/RM
      DO 7 I=1,I MAX
      II=I+1
      IX=I MAX*(J-1)+I
      IJ=I MAX*(J-1)+I
      DH=H(I I)-H(I)
      B=B+G(I J)*S*DH
      FORMAT (, BD:,4I5,10F6.2)
      D=D-G(I J)*C*DH
      RETURN
      END
      C000004

```

SUBROUTINE FIELD2 (RS,SIG,SOLN,G,H,SCF,DGN)

```

C      FIELD2 COMPUTES THE SAME INVERSION AS SUBROUTINE FIELD, EXCEPT THAT
C      THE COEFFICIENTS B AND D ARE COMPUTED INDIVIDUALLY AS USED BY
C      CALLING BD. DISK STORAGE IS NOT REQUIRED, BUT COMPUTING TIME IS
C      MUCH GREATER. FIELD2 IS UTILIZED BY SPECIFYING A NEGATIVE MODE ON

```

SUB02850  
SUB02860  
SUB02870  
SUB02880  
SUB02890  
SUB02900  
SUB02910  
SUB02920  
SUB02930  
SUB02940  
SUB02950  
SUB02960  
SUB02970  
SUB02980  
SUB02990  
SUB03000  
SUB03010  
SUB03020  
SUB03030  
SUB03040  
SUB03050  
SUB03060  
SUB03070  
SUB03080  
SUB03090  
SUB03100  
SUB03110  
SUB03120  
SUB03130  
SUB03140  
  
SUB03160  
SUB03170  
SUB03180  
SUB03190  
SUB03200  
SUB03210  
SUB03220  
SUB03230  
SUB03240

SUB03250  
SUB03260  
SUB03270  
SUB03280  
SUB03290  
SUB03300



```

C THE INPUT PARAMETER. VALUE OF THE FIELD FUNCTION AT A PARTICULAR
C FIELD EVALUATES THE VALUE OF THE FIELD FUNCTION AT A PARTICULAR
C POINT DESIGNATED IN CYLINDRICAL COORDINATES, BY USING THE INVERSION
C EQUATION OF MALDONADO, ET AL. FIELD CALLS SUBROUTINES BD & GARRAY.
SUB033310
SUB033320
SUB033330
SUB033340
SUB033350
SUB033360
SUB033370
SUB033380
SUB033390
SUB033400
SUB033410
SUB033420
SUB033430
SUB033440
SUB033450
SUB033460
SUB033470
SUB033480
SUB033490
SUB033500
SUB033510
SUB033520
SUB033530
SUB033540
SUB033550
SUB033560
SUB033570
SUB033580
SUB033590
SUB033600
SUB033610
SUB033620
SUB033630
SUB033640
SUB033650
SUB033660
SUB033670
SUB033680
SUB033690
SUB033700
SUB033710
SUB033720
SUB033730
SUB033740
SUB033750
SUB033760
SUB033770
SUB033780

COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /TAB/ INDEX,KEXTRA,MEXTRA,KLIMIT,MLIMIT,KOUT,MOUT
COMMON /SYM/ ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
DIMENSION G(IJMX),H(IIMX,5),SCF(IJMX,6)
C INITIALIZE THE VALUES:
INDEX=0
MTIMER=0
KOUT=0
MOUT=0
MMAX=0
KMAX=0
TOTAL=0.
JJMX6=JJMX*6
IIMX2=(IIMX+1)/2
AR=ALPHA*RS
EXPON=EXP(-ARG)
PIE=3.141592653589793
APP=ALPHA/PIE/PIE
M=0
RM=M
RIMAX=IMAX
DX=2./RIMAX
RJMAX=JMAX
DXI=2.*PIE/FCU
C INITIALIZE THE MODIFIED HERMITE POLYNOMIAL ARRAY: VECTORS:
C (1)=H1, (2)=H2, (3)=ALPHA*X(I), (4)=HM+2 STORED, (5)=HM+1 STORED
DO 1 I=1,IIMX2
RII=II
IIM=IIMX-II+1
H(II,3)=ALPHA*(RII*DX-DX-1.)
H(IIM,3)=-H(II,3)
H(II,1)=2.*H(II,3)
H(II,2)=(H(II,3)*H(II,1)-1.)/3.
H(IIM,1)=-H(II,1)
H(IIM,2)=H(II,2)
H(IIM,5)=H(II,2)
H(IIM,4)=H(II,1)
H(IIM,4)=H(IIM,1)
SIGN=1.
FM=1.
C INITIALIZE THE SIN/COS ARRAY:

```





```

11 DO 11 J=1,JJMX
    RJM=J-1
    SCF(J,1)=0.
    SCF(J,2)=1.
    SCF(J,3)=SIN(RJM*DXI-PIE)
    SCF(J,4)=COS(RJM*DXI-PIE)
    SCF(J,5)=0.
    SCF(J,6)=0.
    MS=0
    COMMENCE THE M LOOP:
2   SIGN=-SIGN
    K=0
    RK=K
    ARM=1.
    IF (M.NE.0) ARM=AR**M
    KTIMER=0
    SIGNK=-1.
    COMPUTE THE K=0 & K=1 ORDERS OF LAGUERRE POLYNOMIAL FOR GIVEN M:
    PM=0.
    P=SQRT(1./FM)
    PP=(RM+1.-ARG)*SQRT(1./FM/(RM+1.))
    ADVANCE THE SIN/COS ARRAY FOR NEW M:
    DO 12 J=1,JJMX
        SCF(J,1)=SCF(J,1)*SCF(J,4)+SCF(J,2)*SCF(J,3)
        SCF(J,2)=SCF(J,2)*SCF(J,4)-SCF(J,1)*SCF(J,3)
    DO 13 J=1,JJMAX
        SCF(J,5)=SCF(J+1,1)-SCF(J,1)
        SCF(J,6)=SCF(J+1,2)-SCF(J,2)
    TEST FOR SYMMETRY SKIPS:
    IF (MS.EQ.MSYM) MS=0
    TOTAL=0.
    MS=MS+1
    IF (MS.NE.1) GO TO 7
    RMS=RM**SIGN
    RMS=COS(RMS)
    SMS=SIN(RMS)
    COMMENCE THE K LOOP:
    INDEX=INDEX+1
    SIGNK=-SIGNK
    CALL THE B & D COEFFICIENTS AND COMPUTE THE M,K SUMMATION TERM:
    CALL BD (M,K,G,H,SCF,B,D,JJMX6)
    IF (DGN.LE.-2.) WRITE (6,89) M,K,B,D
    BRACKET=B
    IF (RM.EQ.0.) GO TO 4
    BRACKET=B*CMS+D*SMS
    ADD=SIGNK*BRACKET*P*ARM
    TOTAL=TOTAL+ADD
4   ESTABLISH CHECK AS THE RELATIVE SIZE OF THE M,K TERM OF THE SERIES:
    C

```

SUB03790  
SUB03800  
SUB03810  
SUB03820  
SUB03830  
SUB03840  
SUB03850  
SUB03860  
SUB03870  
SUB03880  
SUB03890  
SUB03900  
SUB03910  
SUB03920  
SUB03930  
SUB03940  
SUB03950  
SUB03960  
SUB03970  
SUB03980  
SUB03990  
SUB04000  
SUB04010  
SUB04020  
SUB04030  
SUB04040  
SUB04050  
SUB04060  
SUB04070  
SUB04080  
SUB04090  
SUB04100  
SUB04110  
SUB04120  
SUB04130  
SUB04140  
SUB04150  
SUB04160  
SUB04170  
SUB04180  
SUB04190  
SUB04200  
SUB04210  
SUB04220  
SUB04230  
SUB04240  
SUB04250



```

CHECK=ABS(ADD)
IF (TOTAL.GT.EPS) CHECK=ABS(ADD/TOTAL)
C ADVANCE THE K INDEX:
K=K+1
RK=K
ORDER=M+2*K+1
C GENERATE THE NEXT ORDER OF LAGUERRE POLYNOMIAL FOR NEW K:
PM=P
P=PP
PP=PP*(ORDER-ARG)-PM*SQRT(RK*(RM+RK))
PP=PP/SQRT((RK+1.)*(RM+RK+1.))
C GENERATE THE NEXT ORDER OF THE SET OF HERMITE POLYNOMIALS FOR NEW K:
HA=SQRT(RK*(RK+RM))/ORDER
HB=2.*SQRT((RK+1.)*(RM+RK+1.))/(ORDER+1.)/(ORDER+2.)
DO 5 I=1,IIMX2
IIM=IIMX-I+1
H(I,1)=2.*(H(I,3)*H(I,2)-HA*H(I,1))
H(I,1,1)=SIGN*H(I,1)
H(I,1,2)=HB*(H(I,3)*H(I,1)-ORDER*H(I,2))
5 C SET K TIMER TO PROVIDE EXTRA K TERMS AFTER CHECK < EPS:
KTIMER=KTIMER+1
IF (K.GE.KLIMIT) GO TO 6
IF (MTIMER.LE.MEXTRA) GO TO 2
IF (CHECK.GE.EPS) KTIMER=0
IF (KTIMER.LE.KEXTRA) GO TO 3
GO TO 7
C END OF K LOOP: ADVANCE M AND COMPUTE NEW TOTAL:
KOUT=KOUT+1
M=M+1
IF (K.GT.KMAX) KMAX=K
RM=M
FM=FM*RM
C REGENERATE THE HERMITE ARRAY FOR NEW M, K=0:
DO 8 I=1,IIMX2
IIM=IIMX-I+1
H(I,2)=H(I,1,4)*SQRT(RM)/(RM+1.)
H(I,1,1)=H(I,1,5)*(RM+2.)
H(I,1,1)=-SIGN*H(I,1)
H(I,1,2)=2.*SQRT(RM+1.)*(H(I,3)*H(I,1)-(RM+1.)*H(I,2))
H(I,1,2)=H(I,1,2)/(RM+2.)/(RM+3.)
H(I,1,4)=H(I,1,1)
H(I,1,5)=H(I,1,2)
C SET M TIMER FOR EXTRA M TERMS:
IF (MS.EQ.1) MTIMER=MTIMER+1
IF (JSYM.GT.JMAX) GO TO 9
IF (K.GT.KEXTRA) MTIMER=0
IF (M.GE.MLIMIT) GO TO 9
C END OF M LOOP: COMPUTE OUTPUT SOLN.

```



SUB04750  
SUB04760  
SUB04770  
SUB04780  
SUB04790  
SUB04800  
SUB04810  
SUB04820  
SUB00010  
SUB00020

```

9      MOUT=M-1      EQ.0) KOUT=KMAX-1
      IF (KOUT.EQ.0) KOUT=KMAX-1
      SOLN=TOTAL*EXPON*APP/2.
      FORMAT (1,M=' ,I4,', K=' ,I4,', B=' ,E10.4,', D=' ,E10.4)
      RETURN
      END
89

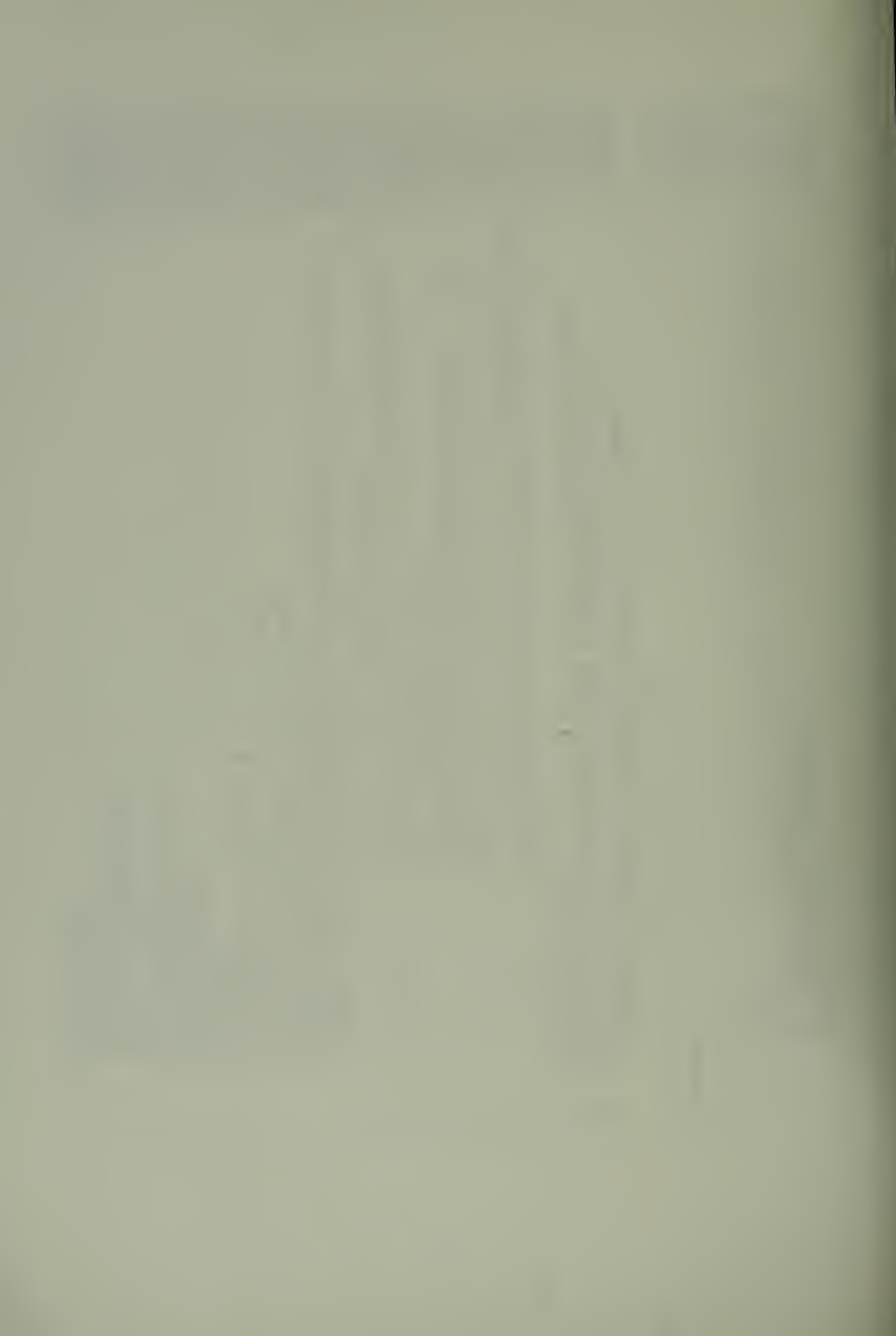
```

C000005  
C

```

C      SUBROUTINE GARRAY (G,GA,NOF,DGN,NUMB,XO,YO,PHISYM)
C
C      GARRAY FILLS THE DATA ARRAY OVER AN ORTHOGONAL AREA WITH
C      THE REGULAR DATA OBTAINED BY THE METHOD CORRESPONDING TO THE
C      PARTICULAR MODE:
C
C      MODE 1 - DATA OBTAINED BY SAMPLING A KNOWN FUNCTION SUPPLIED
C               IN SUBROUTINE FUNCT AND SAMPLED IN SUBROUTINE GOLF.
C
C      MODE 2 - DATA OBTAINED BY GENERATING A REGULAR ARRAY FROM
C               IRREGULAR EXPERIMENTAL INPUT DATA READ IN. CALLS
C               SUBROUTINE SHEET. (EXPERIMENTAL DATA MAY
C               BE SIMULATED, SEE 'SHEET')
C
C      MODE 3 - UTILIZES RAW DATA TAKEN AT THE PROPER INTERVAL,
C               OR PREVIOUSLY GENERATED, AND READ DIRECTLY INTO THE
C               GARRAY. CALLS SUBROUTINE READ.
C
C      COMMON I MAX,JMAX,I IMX,JJMX,I JMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
C      COMMON /SYM/ I SYM,J SYM,M SYM,FCU,IMS,JMS,QSYM
C      COMMON /IO/ CMS,IN1,IN2,IN4
C      DIMENSION G(I MAX,J MAX),GA(I MAX,J MAX)
C      PIE=3.141592653589793
C      HS=SIZE/2.
C      IF (MODE.GT.3) MODE=1
C      RIMX=I MAX
C      RJMX=J MAX
C      DELR=SIZE/RIMX
C      DELXI=2.*PIE/FCU
C      IF (MODE.GT.1) GO TO 2
C      DO 1 J=1,JMS
C      RJ=J
C      XI=(RJ-.5)*DELXI-PIE
C      J2=J+2*(JMS-J)
C      J3=J+JMAX/2
C      J4=J2+JMAX/2
C
C      SUB00030
C      SUB00040
C      SUB00050
C      SUB00060
C      SUB00070
C      SUB00080
C      SUB00090
C      SUB00100
C      SUB00110
C      SUB00120
C      SUB00130
C      SUB00140
C      SUB00150
C      SUB00160
C      SUB00170
C      SUB00180
C      SUB00190
C      SUB00200
C      SUB00210
C      SUB00220
C      SUB00230
C      SUB00240
C      SUB00250
C      SUB00260
C      SUB00270
C      SUB00280
C      SUB00290
C      SUB00300
C      SUB00310
C      SUB00320
C      SUB00330
C      SUB00340
C      SUB00350
C      SUB00360
C      SUB00370
C      SUB00380

```



SUB00390  
SUB00400  
SUB00410  
SUB00420  
SUB00430  
SUB00440  
SUB00450  
SUB00460  
SUB00470  
SUB00480  
SUB00490  
SUB00500  
SUB00510  
SUB00520  
SUB00530  
SUB00540  
SUB00550  
SUB00560  
SUB00570  
SUB00580  
SUB00590  
SUB00600  
SUB00610  
SUB00620

```

DO 1 I=1,IMS
RI=I
II=I*MAX+1-I
R=(RI-.5)*DELR-HS
CALL GOLF (R,XI,GIJ,NOF,DGN,NUMB)
G(I,J)=GIJ
IF (ISYM.EQ.2) G(II,J)=GIJ
IF (ISYM.EQ.2) GO TO 1
G(II,J3)=GIJ
IF (JSYM.EQ.0) GO TO 1
G(I,J2)=GIJ
G(II,J4)=GIJ
CONTINUE
GO TO 4
1 IF (MODE.GT.2) GO TO 3
2 CALL SHEET (G,GA,XO,YO,PHISYM,NOF)
GO TO 4
3 CALL READ (Z,XO,YO,PHISYM,NOF,IMAX,JMAX,G)
4 IF (DGN.GE.2) WRITE (6,39)
39 RETURN
FORMAT (' GARRAY RETURNS ')
END
C000006
C

```

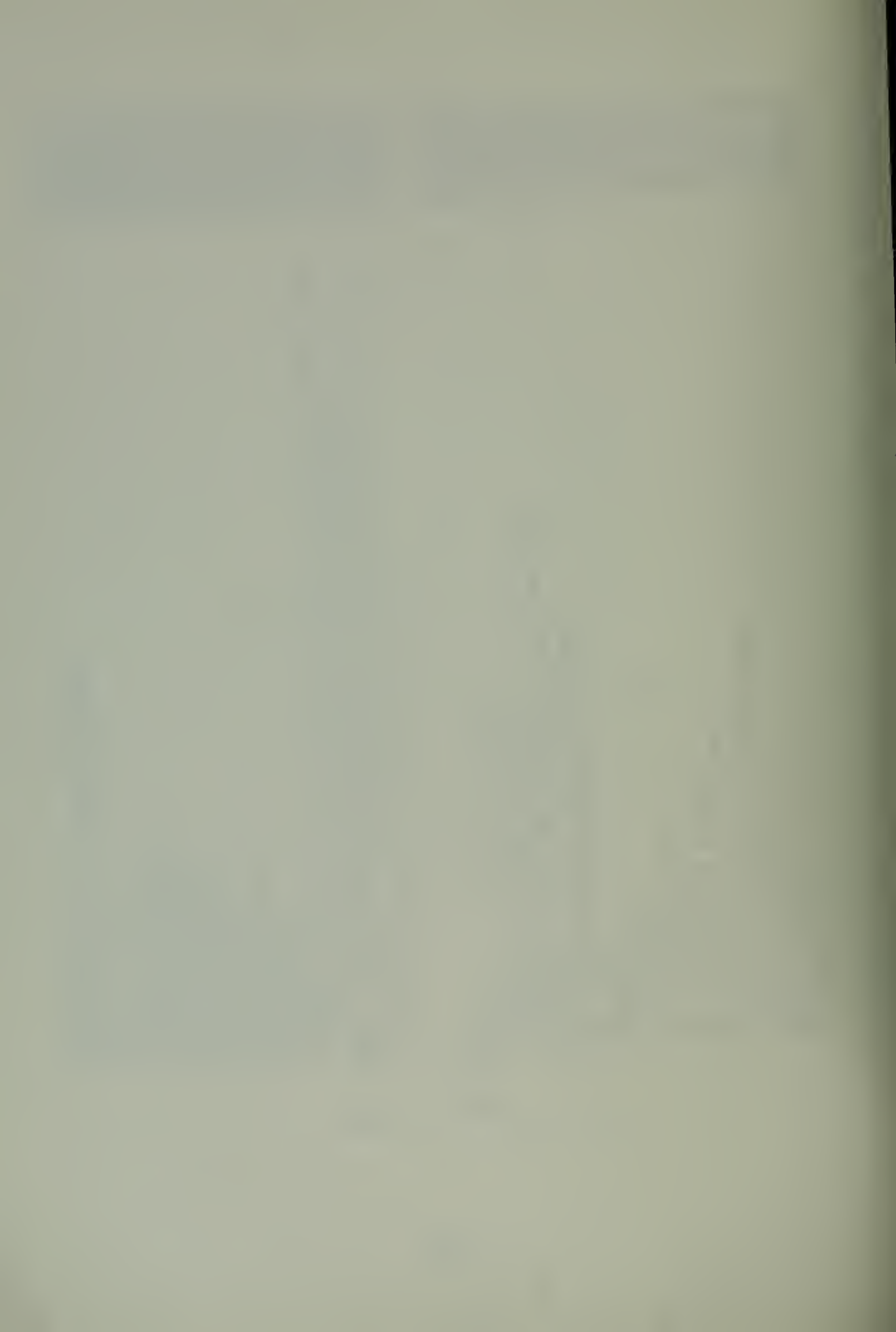
SUB00630  
SUB00640  
SUB00650  
SUB00660  
SUB00670  
SUB00680  
SUB00690  
SUB00700  
SUB00710  
SUB00720  
SUB00730  
SUB00740  
SUB00750  
SUB00760  
SUB00770  
SUB00780  
SUB00790  
SUB00800  
SUB00810  
SUB00820  
SUB00830  
SUB00840

```

SUBROUTINE GOLF (R,XI,GIJ,NOF,DGN,NUMB)
C
C GOLF COMPUTES THE FUNCTION G(R,XI) FOR A PARTICULAR LINE OF SIGHT
C FROM A KNOWN FUNCTION CONTAINED IN SUBROUTINE FUNCT.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
ZERO=0.
LMAX=IMAX*3
RLMAX=LMAX
DELXP=SIZE/RLMAX
SXI=SIN(XI)
CXI=COS(XI)
DELYS=DELXP*CXI
DELYS=DELXP*SXI
XP=DELXP*.5-SIZE/2.
XS=XP*CXI-R*SXI
YS=XP*SXI+R*CXI
GIJ=0.
DO 1 L=1,LMAX
RL=L
CALL FUNCT (XS,YS,F,NOF,DGN,NUMB)
GIJ=GIJ+F

```





```

1  XS=XS+DELYS
   YS=YS+DELYS
   IF (GIJ.NE.0.) GIJ=GIJ*DELXP*BOX
   IF ((SD.EQ.0.) .OR. (NUMB.EQ.1)) GO TO 2
   IF (DGN.GE.3) WRITE (6,28) IX
   CALL GAUSS (IX,SD,ZERO,RV)
   GIJ=GIJ+RV
2  IF (DGN.GE.3) WRITE (6,29) R,XI,GIJ
   RETURN
29  FORMAT (1, R='F8.3',1, XI='F8.3',1, GIJ='F8.3')
28  FORMAT (1, GAUSS, IX='I8')
   END
C000007
C

C  SUBROUTINE FUNCT (XS,YS,F,NOF,DGN,NUMB)
C
CP67USERID 1095BOXJ
C
C  FUNCT EVALUATES AS INPUT FUNCTION AT POSITION (X,Y) IN THE TEST
C  SECTION COORDINATE SYSTEM. NOF IDENTIFIES THE EQUATION USED.
C
COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /EQPARA/ A,B,C,D,E,P,Q,S,T,U,V,W,RO,RA,NO,NA,N1,N2
DIMENSION RO(101),RA(101)
AA=A
BB=B
CC=C
DD=D
EE=E
PP=P
IF (NUMB.LE.1) GO TO 50
AA=S
BB=T
CC=U
DD=V
EE=W
PP=Q
PIE=3.141592653589793
HS=SIZE/2.
R=SQRT(XS**2+YS**2)/HS
F=0.
IF (R.GT.1.) GO TO 11
IF (NOF.LE.0) GO TO 11
C 1. AXISYMMETRIC GAUSSIAN:
SUB000850
SUB000860
SUB000870
SUB000880
SUB000890
SUB000900
SUB000910
SUB000920
SUB000930
SUB000940
SUB000950
SUB000960
SUB000970
SUB000980

SUB000990
SUB01000

SUB01020
SUB01030
SUB01040
SUB01050
SUB01060
SUB01070
SUB01080
SUB01090
SUB01100
SUB01110
SUB01120
SUB01130
SUB01140
SUB01150
SUB01160
SUB01170
SUB01180
SUB01190
SUB01200
SUB01210
SUB01220
SUB01230
SUB01240
SUB01250
SUB01260
SUB01270
SUB01280

```



```

1      IF (NOF.GT.1) GO TO 2
      F=AA*EXP(-1.*(R*HS/BB)**2)
      GO TO 11
C
C
2      ADJUSTABLE RECTANGULAR STEP FUNCTION:
      IF (NOF.GT.2) GO TO 3
      F=PP
      IF ((ABS(XS-DD).LE.BB).AND.(ABS(YS-EE).LE.CC)) F=AA
      GO TO 11
C
C
3      DISPLACABLE ELLIPTICAL GAUSSIAN:
      IF (NOF.GT.3) GO TO 4
      F=AA*EXP(-1.*(XS-DD)/BB)**2+((YS-EE)/CC)**2)
      GO TO 11
C
C
4      CONSTANT:
      IF (NOF.GT.4) GO TO 5
      F=AA
      GO TO 11
C
C
5      ADJUSTABLE AND DISPLACABLE ELLIPTIC RAMP FUNCTION:
      IF (NOF.GT.5) GO TO 6
      RBC=SQRT(((XS-DD)/BB)**2+((YS-EE)/CC)**2)
      F=0.
      IF (RBC.LT.1.) F=AA*((1.-RBC)**PP)
      GO TO 11
C
C
6      DISPLACABLE ELLIPTIC STEP FUNCTION:
      IF (NOF.GT.6) GO TO 7
      RBC=SQRT(((XS-DD)/BB)**2+((YS-EE)/CC)**2)
      F=0.
      IF (RBC.LT.1.) F=AA
      GO TO 11
C
C
7      CIRCULAR COSINE-SQUARED FUNCTION OF BB MAXIMA:
      IF (NOF.GT.7) GO TO 8
      F=AA*COS((2.*BB-1.)*PIE/R/2.))**2
      GO TO 11
C
C
8      NUMERICAL FUNCTION: REQUIRES AN INPUT ARRAY READ IN BY
      SUBROUTINE FREAD: N FOLLOWED BY N POINT VALUES. (101 MAX)
      A CONSTANT VALUE AA IS ADDED TO THE FUNCTION.
      IF (NOF.GT.8) GO TO 9
      IF (NUMB.LE.1) N=NO
      IF (NUMB.GT.1) N=NA
      NM=N-1
      NMM=N-2
      RN=N

```

SUB01290  
SUB01300  
SUB01310  
SUB01320  
SUB01330  
SUB01340  
SUB01350  
SUB01360  
SUB01370  
SUB01380  
SUB01390  
SUB01400  
SUB01410  
SUB01420  
SUB01430  
SUB01440  
SUB01450  
SUB01460  
SUB01470  
SUB01480  
SUB01490  
SUB01500  
SUB01510  
SUB01520  
SUB01530  
SUB01540  
SUB01550  
SUB01560  
SUB01570  
SUB01580  
SUB01590  
SUB01600  
SUB01610  
SUB01620  
SUB01630  
SUB01640  
SUB01650  
SUB01660  
SUB01670  
SUB01680  
SUB01690  
SUB01700  
SUB01710  
SUB01720  
SUB01730  
SUB01740  
SUB01750  
SUB01760



```

SUB01770
SUB01780
SUB01790
SUB01800
SUB01810
SUB01820
SUB01830
SUB01840
SUB01850
SUB01860
SUB01870
SUB01880
SUB01890
SUB01900
SUB01910
SUB01920
SUB01930
SUB01940
SUB01950
SUB01960
SUB01970
SUB01980
SUB01990
SUB02000
SUB02010
SUB02020
SUB02030
SUB02040
SUB02050
SUB02060
SUB02070
SUB02080
SUB02090
SUB02100
SUB02110
SUB02120
SUB02130
SUB02140
SUB02150
SUB02160

RI=R*(RN-1.)+1.
IR=INT(RI)
IR=FLOAT(IR)
DI=RI-IR
IF (NUMB.LE.1) F=RO(IR)
IF (NUMB.GT.1) F=RA(IR)
IF ((IR.NE.N).AND.(NUMB.LE.1)) F=F+DI*(RO(IR+1)-RO(IR))
IF ((IR.NE.N).AND.(NUMB.GT.1)) F=F+DI*(RA(IR+1)-RA(IR))
F=F*AA+BB
GO TO 11

9. SPECIAL FUNCTION: MAY BE WRITTEN FOR THE OCCASION AND
   INSERTED IN SUBROUTINE SPFUN
   IF (NOF.GT.9) GO TO 10
   CALL SPFUN (XS,YS,F)
   GO TO 11

EQUATIONS NO. 10 AND BEYOND ARE SET TO ZERO.
F=0.

IF (DGN.GE.4) WRITE (6,99) XS,YS,F
FORMAT ('F8.3','F8.3','F8.3')
RETURN
END
C0000008
C

SUBROUTINE SPFUN (XS,YS,F)

SPFUN IS A SPECIAL ROUTINE FOR EQ'N NO. 9. ANY FUNCTION MAY BE
ENTERED.

COMMON /EQPARA/ A,B,C,D,E,P,Q,S,T,U,V,W,RO,RA,NO,NA,N1,N2
DIMENSION RO(101),RA(101)
F=0.
IF ((ABS(XS).LE.B).AND.(ABS(YS).LE.C)) F=A
RETURN
END
C0000009
C

```



```

SUBROUTINE SHEET (G,D,XO,YO,PHISYM,NOF)
SHEET READS IRREGULARLY SPACED VALUES OF THE LINE INTEGRAL, AS
OBTAINED FROM HOLOGRAPHIC INTERFEROGRAMS. THE INTEGRAL LINES MAY BE
DEFINED EITHER BY GRID INTERCEPT POSITIONS, OR BY ANGLE AND RADIUS
ABOUT THE CENTER OF THE LABORATORY COORDINATE SYSTEM CENTER. LINES
MUST BE ENTERED IN CONSECUTIVE ORDER FROM LOWEST (NEG.) TO HIGHEST
(POS.) RADIUS. DATA MAY BE SIMULATED BY SPECIFYING NCODE=1,
FOLLOWED BY APERTURE POSITIONS FOR A FUNCTION NUMBER IN 'SUBFUNCT'.
SUB02170
SUB02180
SUB02190
SUB02200
SUB02210
SUB02220
SUB02230
SUB02240
SUB02250
SUB02260
SUB02270
SUB02280
SUB02290
SUB02300
SUB02310
SUB02320
SUB02330
SUB02340
SUB02350
SUB02360
SUB02370
SUB02380
SUB02390
SUB02400
SUB02410
SUB02420
SUB02430
SUB02440
SUB02450
SUB02460
SUB02470
SUB02480
SUB02490
SUB02500
SUB02510
SUB02520
SUB02530
SUB02540
SUB02550
SUB02560
SUB02570
SUB02580
SUB02590
SUB02600
SUB02610
SUB02620
SUB02630
SUB02640

COMMON IMAX,JMAX,IMX,JMX,IJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
COMMON /SYM/ISYM,JSYM,MSYM,FCU,IMS,JMS,QSYM
COMMON /IO/ CMS,IN1,IN2,IN4
DIMENSION G(IMAX,JMAX),D(IMAX,JMAX),XI(303),RR(303)
DIMENSION XG(303),XD(303),YG(303),YD(303),XY(303)
NAR=303
PIE=3.141592653589793
MPIE=-PIE
MPIE=2.*PIE/2.
MPIE=PIE/2.
PIE=PIE/2.
ZERO THE ARRAYS:
DO 1 J=1,JMAX
DO 1 I=1,IMAX
G(I,J)=0.
D(I,J)=0.
DO 2 I=1,NAR
XG(I)=0.
XD(I)=0.
YG(I)=0.
YD(I)=0.
XY(I)=0.
XI(I)=0.
RR(I)=0.
READ THE BASIC DATA:
IF (CMS.EQ.1.) REWIND 1
READ (IN1,59) NOF,NCODE
READ (IN1,58) Z,XO,YO,PHISYM,XX,XMN,YMX,YMN
READ (IN1,59) JM
RIMX=IMAX
DR=SIZE/RIMX
R=(-DR-SIZE)/2.
RZO=SQRT(XO**2+YO**2)
GAM=ATANM(YO,XO)
TP=3-ISYM
BT=JSYM
DAN=PIE*TP/BT
HS=SIZE/2.

```









SUB03130  
 SUB03140  
 SUB03150  
 SUB03160  
 SUB03170  
 SUB03180  
 SUB03190  
 SUB03200  
 SUB03210  
 SUB03220  
 SUB03230  
 SUB03240  
 SUB03250  
 SUB03260  
 SUB03270  
 SUB03280  
 SUB03290  
 SUB03300  
 SUB03310  
 SUB03320  
 SUB03330  
 SUB03340  
 SUB03350  
 SUB03360  
 SUB03370  
 SUB03380  
 SUB03390  
 SUB03400  
 SUB03410  
 SUB03420  
 SUB03430  
 SUB03440  
 SUB03450  
 SUB03460  
 SUB03470  
 SUB03480  
 SUB03490  
 SUB03500  
 SUB03510  
 SUB03520  
 SUB03530  
 SUB03540  
 SUB03550  
 SUB03560  
 SUB03570  
 SUB03580  
 SUB03590  
 SUB03600

```

RRH=100GO.
XH=RRH*COS(XIH)
YH=RRH*SIN(XIH)
IF (LPR.EQ.1) GO TO 7
YTX=-RR(IMT)*SIN(XI(IMT))-YO
YTN=-RR(IMT)*SIN(XI(IMT))-YO
XTX=RR(IMT)*COS(XI(IMT))-XO
XTN=RR(IMT)*COS(XI(IMT))-XO
UA=TAN(XI(IMT))
UC=TAN(XI(IMT))
UB=YTX-UA*XTX
UD=YTN-UC*XTN
UH=(UD-UB)/(UA-UC)
YH=XH*UA+UB
RRH=SQRT((XH-XO)**2+(YH-YO)**2)
XIH=ATANM((YH-YO),(XH-XO))
CONTINUE
7 C  FILL THE ANGLE AND RADIUS FOR ANY CODE 3 OR 4 LINES:
DO 9 I=1,IM
IF(XY(I).NE.3.) GO TO 8
BAS=SQRT(RRH**2-RR(I)**2)
XI(I)=XIH-ATANM(RR(I),BAS)
GO TO 9
8 XI(I)= ATANM((YH-YD(I)),(XH-XD(I)))
RR(I)=RRH*SIN(XI(I)-XIH)
9 CONTINUE
C  ANGLES AND RADII ARE NOW FILLED FOR ALL POINTS IN THIS LINE.
C  VACATE THE SET OF VECTORS TO BE USED AS TEMPORARY STORAGE:
DO 10 I=1,IM
XD(I)=0.
YD(I)=0.
XG(I)=0.
YG(I)=RR(I)
XY(I)=D(I,J)
RR(I)=0.
D(I,J)=0.
XI(I)=C.
10 C  CONVERT THE LINE TO REGULAR RADII USING INTERPOLATION:
RR(1)=R+DR
CALL SPLINE(YG,XY,IM,RR(1),D(1,J))
DO 11 I=2,IMAX
RI=I
RR(I)=R+DR*RI
CALL SPLINR(YG,XY,IM,RR(I),D(I,J))
11 C  GENERATE THE VECTOR OF ANGLES FOR THIS COLUMN AND STORE IN G ARRAY:
DO 12 I=1,IMAX
BAS=SQRT(RRH**2-RR(I)**2)
G(I,J)=XIH-ATANM(RR(I),BAS)

```



```

12      YG(I)=0.
C      D(I,J)=XY(I)
C      XY(I)=0.
C      COLUMNS ARE NOW ALL REGULARLY FILLED.
C      NEXT, INTERPOLATE EACH ROW REGULARLY OVER THE ANGLES.
C      DO 23 I=1,IMAX
C      EXPAND THE DATA TO 2 SETS TO ESTABLISH SMOOTH INTERPOLATION.
      JM3=3*JM
      II=IMAX+1-I
      IF (JSYM.NE.0) GO TO 14
      DO 13 J=1,JMS
      J2=J+JMS
      J3=J2+JMS
      XD(J)=D(I,J)
      XD(J2)=D(I,J)
      XD(J3)=D(I,J)
      XG(J)=G(I,J)-TPIE-PHISYM
      XG(J2)=G(I,J)-PIE-PHISYM
      XG(J3)=G(I,J)-PHISYM
      GO TO 16
13      DO 15 J=1,JMS
14      J1=JMS+1-J
      J2=JMS+J
      J3=JM3+1-J
      XD(J1)=D(I,J)
      XD(J2)=D(I,J)
      XD(J3)=D(I,J)
      XG(J1)=G(I,J)-2.*(G(I,J)-PHISYM)-PIE-PHISYM
      XG(J2)=G(I,J)-PIE-PHISYM
      XG(J3)=G(I,J)+2.*(DAN+PHISYM-G(I,J))-PIE-PHISYM
      CONTINUE
      JM2=2*JMS
      JP=JMS/2
      DO 17 J=1,JM2
      XD(J)=XD(J+JP)
      XG(J)=XG(J+JP)
      JJS=JM2+1
      DO 18 J=JJS,JM3
      XD(J)=0.
      XG(J)=0.
18      FIND THE SMALLEST ANGLE
C      XY(1)=1.
      SA=XG(1)
      DO 19 J=1,JM2
      IF (XG(J).GE.SA) GO TO 19
      SA=XG(J)
      XY(1)=J
      CONTINUE
19

```



```

C      FIND THA MAX ANGLE IN THE ROW:
      XY(JM2)=JM2
      SB=XG(JM2)
      DO 20 J=1,JM2
      IF (XG(J).LE.SB) GO TO 20
      SB=XG(J)
      XY(JM2)=J
      CONTINUE
20     C      DETERMINE THE ORDER OF INCREASING ANGLE IN THE ROW
      SB=XG(JM2)
      J=2
      JSA=XY(JJ-1)
      SA=XG(JJ-1)
      JTS=0
      DO 22 J=1,JM2
      IF (XG(J).LE.SA) GO TO 22
      IF (XG(J).GT.SB) GO TO 22
      SB=XG(J)
      XY(JJ)=J
      JTS=1
      CONTINUE
22     IF (JTS.EQ.0) JM2=JJ
      JJ=JJ+1
      IF (JJ.LE.JM2) GO TO 21
      DO 23 J=1,JM2
      JX=XY(J)
      YD(J)=XD(JX)
      INTERPOLATE:
      DXI=2.*PIE/FCU
      XI(J)=DXI/2.-PIE-PHISYM
      CALL SPLINE (XG,YD,JM2,XI(J),G(I,J))
      DO 24 J=2,JMS
      XI(J)=XI(J-1)+DXI
      CALL SPLIN (XG,YD,JM2,XI(J),G(I,J))
      DO 25 J=1,JMS
      XIJ=XI(J)
      XU=XMX
      IF ((XIJ.GE.0.).AND.(XIJ.LT.PIE)) XU=XMN
      YU=YMN
      IF ((XIJ.GE.MPIT).AND.(XIJ.LT.PIT)) YU=YMX
      XL=XMN
      IF ((XIJ.GE.0.).AND.(XIJ.LT.PIE)) XL=XMX
      YL=YMX
      IF ((XIJ.GE.MPIT).AND.(XIJ.LT.PIT)) YL=YMN
      XIJ=SI(XIJ)
      CXIJ=COS(XIJ)
      RMN=(XO-XL)*SXIJ-(YO-YL)*CXIJ
      RMX=(XO-XU)*SXIJ-(YO-YU)*CXIJ

```

SUB04090  
SUB04100  
SUB04110  
SUB04120  
SUB04130  
SUB04140  
SUB04150  
SUB04160  
SUB04170  
SUB04180  
SUB04190  
SUB04200  
SUB04210  
SUB04220  
SUB04230  
SUB04240  
SUB04250  
SUB04260  
SUB04270  
SUB04280  
SUB04290  
SUB04300  
SUB04310  
SUB04320  
SUB04330  
SUB04340  
SUB04350  
SUB04360  
SUB04370  
SUB04380  
SUB04390  
SUB04400  
SUB04410  
SUB04420  
SUB04430  
SUB04440  
SUB04450  
SUB04460  
SUB04470  
SUB04480  
SUB04490  
SUB04500  
SUB04510  
SUB04520  
SUB04530  
SUB04540  
SUB04550  
SUB04560





```

25 DO 25 I=1,IMAX
   IF (RR(I).LT.RMN) G(I,J)=0.
   IF (RR(I).GT.RMX) G(I,J)=0.
C CONTINUE
C EXPAND SYMMETRY SECTOR INTO AN ORTHOGONAL INTERVAL.
25 IF (ISYM.EQ.2) GO TO 27
DO 26 J=1,JMS
J2=JMAX/2+1-J
J3=JMAX/2+J
J4=JMAX+1-J
DO 26 I=1,IMAX
II=IMAX+1-I
G(I,J2)=G(I,J)
G(II,J3)=G(I,J)
G(II,J4)=G(I,J)
26 RETURN
C FOR EVEN SYMMETRY, AVERAGE THE GARRAY COLUMNS.
27 IMS=(2*IMAX+1)/2
DO 28 J=1,JMAX
DO 28 I=1,IMS
II=IMAX+1-I
GST=(G(I,J)+G(II,J))/2.
G(I,J)=GST
G(II,J)=GST
28 RETURN (5I5)
59 FORMAT (10F7.3)
58
C000010
C
C FUNCTION ATANM(Y,X)
C COMPUTES THE ARCTAN OF Y/X BETWEEN -PI AND +PI.
C
PIE=3.141592653589793
PI2=PIE/2.
ATANM=SIGN(PI2,Y)
IF(X.NE.0.) ATANM=ATAN(Y/X)
IF(X.EQ.0.) RETURN
IF(Y.GE.0.) ATANM=PIE+ATANM
IF(Y.LT.0.) ATANM=-PIE+ATANM
RETURN
END
C000011
C
SUB04570
SUB04580
SUB04590
SUB04600
SUB04610
SUB04620
SUB04630
SUB04640
SUB04650
SUB04660
SUB04670
SUB04680
SUB04690
SUB04700
SUB04710
SUB04720
SUB04730
SUB04740
SUB04750
SUB04760
SUB04780
SUB04790
SUB04800
SUB04810
SUB04820
SUB04830
SUB04840
SUB04850
SUB04860
SUB04870
SUB04880
SUB04890
SUB04900
SUB04910
SUB04920
SUB04930
SUB04940
SUB04950
SUB04960
SUB04970
SUB04980
SUB04990
SUB05000
SUB05010

```



```

SUBROUTINE SIM (XD,YD,XG,YG,D,R,XI,XY,XO,YO,PS,XM,XN,YN,I,IM,
1DG,NF)
SIM SIMULATES THE FRINGE NUMBER DATA ONE WOULD OBTAIN FROM THE
HOLOGRAPHIC INTERFEROGRAM PROCESS FOR A KNOWN FUNCTION AS
CONTAINED IN SUBROUTINE FUNCT. THE GRID BOX DIMENSIONS MUST
EXCEED THE INVERSION CIRCLE SIZE, AND APERTURE POINTS SPECIFIED
MUST FALL BETWEEN XI=-40 DEGREES, AND XI=+130 DEGREES.
SUB050200
SUB050230
SUB05030
SUB05030
SUB05040
SUB05050
SUB05050
SUB05060
SUB05070
SUB05080
SUB05090
SUB05100
SUB05110
SUB05120
SUB05130
SUB05140
SUB05150
SUB05160
SUB05170
SUB05180
SUB05190
SUB05200
SUB05210
SUB05220
SUB05230
SUB05240
SUB05250
SUB05260
SUB05270
SUB05280
SUB05290
SUB05300
SUB05310
SUB05320
SUB05330
SUB05340
SUB05350
SUB05360
SUB05370
SUB05380
SUB050010
SUB050020

COMMON IMAX,JMAX,IIMX,IJMX,ALF,SIZ,EPS,MOD,BOX,SD,IX,Z
COMMON /IO/ CMS,IN1,IN2,IN4
READ (IN1,29) XH,YH
ZER=0.
RIM=IM
RI=I
DX=(XM-XN)/(RIM-1.)
DY=(YM-YN)/(RIM-1.)
X1=XM-(RIM-1.)*DX
Y1=YM-(RIM-1.)*DY
XI1=ATANM (YH-Y1,XH-X1)-PS
RRH=SQRT((XH-XO)**2+(YH-YO)**2)
XIO=ATANM(YH-YO,XH-XO)-PS
R1=RRH*SIN(XI1-XIO)
IF ((I.GT.1).AND.(I.LT.IM)) GO TO 1
IF (ABS(R1).LT.SIZ/2.) R1=SIGN(SIZ/2.,R1)
XY=3.
GO TO 2
1 CALL GOLF (R,XI,D,NR,ZER,ZER)
XD=XN
YD=YN
IF (XH.NE.X1) YD=Y1-(X1-XN)*(YH-Y1)/(XH-X1)
IF (YS.GE.YN) GO TO 2
YD=YN
IF (YH.NE.Y1) XD=X1-(Y1-YN)*(XH-X1)/(YH-Y1)
RETURN
FORMAT (10F7.3)
END

C000012
C

SUBROUTINE FREAD (NO,RO,NF,ZZ)
FREAD READS THE NUMERIC ARRAY WHICH IS USED FOR EQUATION 8 OF
SUBROUTINE FUNCT. FIRST CARD IS NUMBER OF POINTS (N.GE.1),
FOLLOWED BY ONE POINT PER CARD.
C
C
C DIMENSION RO(101)
SUB000030
SUB000040
SUB000050
SUB000060
SUB000070
SUB000080
SUB000090

```



```

SUB000100
      READ (NF,89) NO,ZZ
      WRITE(6,90) NO,ZZ
      DO 10 I=1,NO
      READ(NF,88) RO(I)
      WRITE(6,88) RO(I)
      CONTINUE
      10 FORMAT(15,F9.3)
      88 FORMAT(F8.5)
      90 FORMAT(1X,15,F9.3)
      RETURN
      END
      C0000013
      C

SUB000120
      SUBROUTINE GPRINT (G,NUMB)
      C
      C      GPRINT PRINTS THE DATA ARRAY 'G' WHICH WAS INPUT TO
      C      THE PROGRAM IN SUBROUTINE GARRY.
      C
      COMMON IMAX,JMAX,IIMX,JJMX,IJMX,ALPHA,SIZE,EPS,MODE,BOX,SD,IX,Z
      DIMENSION G(IJMX)
      DIMENSION X(15)
      DATA HYP,VERT/1H-,1H1/
      IF (NUMB.EQ.1) WRITE (6,99) MODE,Z
      IF (NUMB.EQ.2) WRITE (6,92) Z
      JM2=JMAX/2
      RIJMX=IJMX/2
      RIMAX=IMAX
      RJMAX=JMAX
      DX=SIZE/RIJMX
      DXI=360./RJMAX
      INTRVL SETS THE NUMBER OF TERMS PRINTED PER LINE. IF IT IS ALTERED,
      ONE MUST ALSO REDIMENSION X AND ALTER FORMATS 98, 97, AND 95.
      INTRVL=15
      IB=1
      IT=IB+INTRVL-1
      IF (IT.GT.IMAX) IT=IMAX
      IBT=IT-IB+1
      WRITE (6,98) (II,II=IB,IT)
      DO 2 I=1,IBT
      RI=IB-1+I
      X(I)=-SIZE/2.+(RI-.5)*DX
      LM=7*I
      WRITE (6,97) (X(I),I=1,IBT)
      WRITE (6,96) (HYP,L=1,LM),VERT
      JMH=JM2+1
      DO 3 J=JMH,JMAX
      RJ=J

```









```

COMMON /IO/ CMS, IN1, IN2, IN4
DIMENSION G(IMAX, JMAX)
READ (IN1, 39) NOF, IMAX, JMAX, ISYM, JSYM, IMS, JMS
READ (IN1, 40) Z, XO, YO, PHISYM
DO 10 J=1, JMS
  READ (IN1, 38) (G(I, J), I=1, IMS)
  WRITE (6, 37) NOF, Z, XO, YO, PHISYM, IMAX, JMAX, JSYM
  RJMX=JMAX
  MSYM=JSYM
  IF ((MSYM.EQ.0).OR.(MSYM.GT.JMAX)) MSYM=1
  FCU=ISYM*JSYM*JMAX
  IF (JSYM.GT.JMAX) FCU=JMAX
  QSYM=FCU/RJMX
  DO 4 J=1, JMS
    IF (ISYM.EQ.1) GO TO 2
    DO 1 I=1, IMS
      II=IMAX+1-II
      G(II, J)=G(I, J)
    GO TO 4
  J2=JMAX/2+1-J
  J3=JMAX/2+J
  J4=JMAX+1-J
  DO 3 I=1, IMAX
    II=IMAX+1-II
    G(I, J2)=G(II, J)
    G(II, J3)=G(I, J)
    G(II, J4)=G(I, J)
  CONTINUE
  FORMAT(10I5)
  FORMAT(4F7.3)
  FORMAT(10F7.3)
  MODE 3 READS GARRAY DIRECTLY: NOF='I4',
1, ' , JMAX='I4', PHISYM='F7.3',
2, ' , XO='F7.3', JSYM='I4',
  RETURN
END
C000016
C

SUBROUTINE MAP (IM, JM, A, N, Z, BAND)
C
C MAP CALLS SUBROUTINE MIMPII AND PLOTS A CONTOUR MAP OF THE ARRAY
C
  DIMENSION A(IM, JM), T(24)
  DATA BL/1H/
  DO 1 I=1, 24
    T(I)=BL
  1

```

SUB000950  
SUB000960  
SUB000970

SUB010000  
SUB01010  
SUB01020  
SUB01030  
SUB01040  
SUB01050  
SUB01060  
SUB01070  
SUB01080  
SUB01090  
SUB01100  
SUB01110  
SUB01120  
SUB01130  
SUB01140  
SUB01150  
SUB01160  
SUB01170  
SUB01180  
SUB01190  
SUB01200  
SUB01210  
SUB01220

SUB01240  
SUB01250  
SUB01260  
SUB01270  
SUB01280  
SUB01290  
SUB01300

SUB01310  
SUB01320  
SUB01330  
SUB01340  
SUB01350  
SUB01360  
SUB01370  
SUB01380







```

RINK=RANGE/80.
TOP=BIG+RINK
CEN=BIG
BOT=BIG-RINK
KC=0
DO 7 K=1,41
  IC=0
  DO 6 I=1,101
    ROW(I)=BL
    IF((I.EQ.1).OR.(I.EQ.51).OR.(I.EQ.101))ROW(I)=PL
    IF((I.EQ.1).OR.(K.EQ.41))ROW(I)=PL
    IF((K.EQ.1).OR.(K.EQ.41))ROW(I)=PL
    IF((TOP.GE.C).AND.(BOT.LE.0.))ROW(I)=DH
    IF((IC.EQ.5)GO TO 3
    GO TO 4
  IC=0
  IF(KC.EQ.10)ROW(I)=PL
  IF(KPT.LE.2)GO TO 5
  IF((D(I).LE.TOP).AND.(D(I).GE.BOT))ROW(I)=ST
  IF((B(I).LE.TOP).AND.(B(I).GE.BOT))ROW(I)=EX
  IC=IC+1
  IF(KC.EQ.5)KC=0
  IF(KC.NE.0)WRITE (6,65) (ROW(I),I=1,101)
  IF(KC.EQ.0)WRITE (6,68) CEN,(ROW(I),I=1,101)
  TOP=TOP-2.*RINK
  CEN=CEN-2.*RINK
  BOT=BOT-2.*RINK
  KC=KC+1
  WRITE (6,66) (ST,I=1,120)
  FORMAT (//,J=,I3//)
  FORMAT (1X,F8.3,1X,101A1)
  FORMAT (1H1,//121A1//)
  FORMAT (//121A1//)
  FORMAT (10X,101A1)
  RETURN
END
C000018
C

```

```

SUBROUTINE INTERP (A,IM,AS,B,JM,BS)
C
C INTERP CONVERTS A REGULAR VECTOR A OF IM POINTS TO A REGULAR VECTOR
C B OF JM POINTS. OS=.5 FOR A VECTOR WITH POINTS DEFINED IN THE
C CENTER OF THE INTERVAL, AS AND BS ARE THE % OF AN INTERVAL FROM THE
C EDGE OF THE FIELD TO THE FIRST POINT (.0 OR .5 FOR EDGE OR CENTER
C DEFINED POINTS)
C
C DIMENSION A(IM),B(JM)
C
SUB02220
SUB02230
SUB02240
SUB02250
SUB02260
SUB02270
SUB02280
SUB02290
SUB02300

```



```

RIM=IM
RJM=JM
RAT=(RIM-1.+2.*AS)/(RJM-1.+2.*BS)
DO 2 I=1,JM
  BI=I
  AI=RAT*(BI+BS)-AS
  IA=AI
  F=AI-FLOAT(IA)
  IF((IA.EQ.0).OR.(IA.EQ.JM)) GO TO 1
  B(I)=A(IA)+F*(A(IA+1)-A(IA))
  GO TO 2
1 IF(IA.EQ.0) B(I)=A(1)*(F-AS)/(1.-AS)
2 IF(IA.EQ.JM) B(I)=A(JM)*F/(1.-AS)
CONTINUE
RETURN
END

```

C000019

```

C-----
C SUBROUTINE SPLINE
C
C PURPOSE
C PROVIDES INTERPOLATED VALUE USING "CUBIC SPLINE FITTING"
C
C USAGE
C FIRST CALL TO SUBROUTINE:
C CALL SPLINE(X,Y,M,XINT,YINT)
C
C SUBSEQUENT CALLS:
C CALL SPLINN(X,Y,M,XINT,YINT)
C
C DESCRIPTION OF PARAMETERS
C X: MONOTONICALLY INCREASING ABSCISSA ARRAY
C Y: ONE-FOR-ONE CORRESPONDING ORDINATE ARRAY
C M: NUMBER OF X AND Y VALUES SUPPLIED < OR = 300
C XINT: VALUE OF ABSCISSA FOR WHICH CORRESPONDING ORDINATE
C IS TO BE INTERPOLATED (OR EXTRAPOLATED)
C YINT: INTERPOLATED (OR EXTRAPOLATED) ORDINATE VALUE
C
C REMARKS
C IF SPECIFIED X FALLS OUTSIDE OF RANGE, AN EXTRAPOLATED
C VALUE WILL BE SUPPLIED
C
C SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED
C SUBROUTINE SPLICO IS INCLUDED IN SUBROUTINE SPLIN PACKAGE
C
C-----
C SPL00010
C SPL00020
C SPL00030
C SPL00040
C SPL00050
C SPL00060
C SPL00070
C SPL00080
C SPL00090
C SPL00100
C SPL00110
C SPL00120
C SPL00130
C SPL00140
C SPL00150
C SPL00160
C SPL00170
C SPL00180
C SPL00190
C SPL00200
C SPL00210
C SPL00220
C SPL00230
C SPL00240
C SPL00250
C SPL00260
C SPL00270
C SPL00280
C SPL00290
C SPL00300
C SPL00310

```





C C C C C C C C C C C C C C C C  
 MATHEMATICAL METHOD  
 UPON FIRST ENTRY TO SPLIN, A CALL TO SPLICO IS MADE TO  
 DETERMINE THE COEFFICIENTS TO BE USED IN PERFORMING THE  
 INTERPOLATIONS. SEARCH FOR BRACKETING ABSCISSA VALUES IS  
 ALWAYS MADE FROM THE REFERENCE LAST USED IN INTERPOLATING.  
 REFERENCE  
 PENNINGTON, RALPH H., "INTRODUCTORY COMPUTER METHODS AND  
 NUMERICAL ANALYSIS", THE MACMILLAN COMPANY, NEW YORK, 1965  
 -----  
 C C C C C C C C C C C C C C C C

SUBROUTINE SPLINE(X,Y,M,XINT,YINT)  
 DIMENSION X(M),Y(M),C(4,300)  
 CALL SPLICO(X,Y,M,C)  
 K=1  
 ENTRY SPLIN(X,Y,M,XINT,YINT)  
 IF(XINT-X(1)) 70,1,2  
 3 70 K=1  
 GO TO 7  
 1 YINT=Y(1)  
 RETURN  
 2 IF(XINT-X(K+1)) 6,4,5  
 4 YINT=Y(K+1)  
 RETURN  
 5 K=K+1  
 IF(M-K) 71,71,3  
 71 K=M-1  
 GO TO 7  
 6 IF(XINT-X(K)) 13,12,11  
 12 YINT=Y(K)  
 RETURN  
 13 K=K-1  
 GO TO 6  
 7 PRINT 101,XINT  
 101 FORMAT(8H,XINT = E18.9,32H, OUT OF RANGE FOR INTERPOLATION)  
 11 YINT=(X(K+1)-XINT)\*(C(1,K))\*X(K+1)-XINT)\*\*2+C(3,K))  
 YINT=YINT+(XINT-X(K))\*(C(2,K))\*X(K))\*\*2+C(4,K))  
 RETURN  
 END



```

SUBROUTINE SPLICO(X,Y,M,C)
DIMENSION X(M),Y(M),C(4,300),D(300),P(300),E(300),A(300,3),B(300),
1 Z(300)
MM=M-1
DO 2 K=1,MM
D(K)=X(K+1)-X(K)
P(K)=D(K)/6.
2 E(K)=(Y(K+1)-Y(K))/D(K)
DO 3 K=2,MM
3 B(K)=E(K)-E(K-1)
A(1,2)=-1.-D(1)/D(2)
A(1,3)=D(1)/D(2)
A(2,3)=P(2)-P(1)*A(1,3)
A(2,2)=2.*(P(1)+P(2))-P(1)*A(1,2)
A(2,3)=A(2,3)/A(2,2)
B(2)=B(2)/A(2,2)
DO 4 K=3,MM
4 A(K,2)=2.*(P(K-1)+P(K))-P(K-1)*A(K-1,3)
B(K)=B(K)-P(K-1)*B(K-1)
A(K,3)=P(K)/A(K,2)
B(K)=B(K)/A(K,2)
Q=D(M-2)/D(M-1)
A(M,1)=1.+Q+A(M-2,3)
A(M,2)=-Q-A(M,1)*A(M-1,3)
B(M)=B(M-2)-A(M,1)*B(M-1)
Z(M)=B(M)/A(M,2)
MN=M-2
DO 6 I=1,MN
6 K=M-I
Z(K)=B(K)-A(K,3)*Z(K+1)
Z(1)=-A(1,2)*Z(2)-A(1,3)*Z(3)
DO 7 K=1,MM
7 Q=1./(6.*D(K))
C(1,K)=Z(K)*Q
C(2,K)=Z(K+1)*Q
C(3,K)=Y(K)/D(K)-Z(K)*P(K)
C(4,K)=Y(K+1)/D(K)-Z(K+1)*P(K)
RETURN
END

```

SPL00730  
SPL00750  
SPL00760  
SPL00770  
SPL00780  
SPL00790  
SPL00800  
SPL00810  
SPL00820  
SPL00830  
SPL00840  
SPL00850  
SPL00860  
SPL00870  
SPL00880  
SPL00890  
SPL00900  
SPL00910  
SPL00920  
SPL00930  
SPL00940  
SPL00950  
SPL00960  
SPL00970  
SPL00980  
SPL00990  
SPL01000  
SPL01010  
SPL01020  
SPL01030  
SPL01040  
SPL01050  
SPL01060  
SPL01070  
SPL01080  
SPL01090  
SPL01100  
SPL01110  
SPL01120

C000020

MET00010  
MET00020  
MET00030  
MET00040  
MET00050  
MET00060

\*\*\*\*\*

SUBROUTINE MTMPII

PURPOSE



MET00070  
 MET00080  
 MET00090  
 MET00100  
 MET00110  
 MET00120  
 MET00130  
 MET00140  
 MET00150  
 MET00160  
 MET00170  
 MET00180  
 MET00190  
 MET00200  
 MET00210  
 MET00220  
 MET00230  
 MET00240  
 MET00250  
 MET00260  
 MET00270  
 MET00280  
 MET00290  
 MET00300  
 MET00310  
 MET00320  
 MET00330  
 MET00340  
 MET00350  
 MET00360  
 MET00370  
 MET00380  
 MET00390  
 MET00400  
 MET00410  
 MET00420  
 MET00430  
 MET00440  
 MET00450  
 MET00460  
 MET00470  
 MET00480  
 MET00490  
 MET00500  
 MET00510  
 MET00520  
 MET00530  
 MET00540

MTMPII WILL PRODUCE, ON THE PRINTER, A CONTOUR MAP OF ANY SINGLE PRECISION TWO DIMENSIONAL ARRAY.	
CALL MTMPII(Y,N,M,T,BND,AZ,BZ,AMIN,IJT,ICON)	
DESCRIPTION OF PARAMETERS	
Y - THE ARRAY TO BE CONTOURED. DIMENSIONED Y(N,M)	
N - NUMBER OF ROWS IN Y.	
M - NUMBER OF COLUMNS IN Y.	
T - ARRAY FOR PLOT TITLE. REAL*4. T(24).	
BND - BANDWIDTH FOR THE CONTOURING. IF BND IS ZERO A BANDWIDTH WILL BE CALCULATED AS FOLLOWS BND=(MAX(Y)-MIN(Y))/15.	
AZ - A LINEAR TRANSFORMATION MAYBE PERFORMED ON THE ARRAY Y OF THE FOLLOWING FORM AZ=Y+BZ. IF AZ=0 THEN AZ WILL BE COMPUTED SUCH THAT MAX(IMAX(Y)),IMIN(Y)) WILL BE LESS THAN 1, AND BZ WILL BE LEFT AS INPUT.	
BZ - SEE UNDER AZ	
AMIN - THE LEVEL AT WHICH COUTOURING WILL BEGIN. IF AMIN > MIN(Y) THEN AMIN WILL BE CALCULATED TO BE, WITH REFERENCE AT ZERO, THE NEXT LOWER CONTOUR LEVEL FROM MIN(Y) AS DETERMINED BY BND.	
IJT - IF IJT=0 AMIN WILL BE CALCULATED AS DESCRIBED IF IJT>0 AMIN WILL BE CALCULATED AS DESCRIBED IF IJT<0 AMIN WILL BE CALCULATED AS DESCRIBED	
ICON - IF ICON=0 NO CONTOURING WILL BE DONE BUT THE ARRAY Y WILL BE PRINTED IN THE PLOT FORMAT.	
REMARKS	
MTMPII REQUIRES A PRINTER WITH 132 PRINT POSITIONS. IF NECESSARY THE MAP WILL BE SEGMENTED COLUMNWISE. THE ROWS AND COLUMNS ARE NUMBERED ALONG THE EDGES. SO THAT A SEGMENTED MAP MAYBE EASILY JOINED TOGETHER. ONLY THREE SIGNIFICANT FIGURES WILL BE PRINTED AT EACH POINT. THE POSITION OF THE FIRST SIGNIFICANT DIGIT WILL BE DETERMINED BY MAX(IMAX(Y)),IMIN(Y)). THE PLOT WILL BE PRODUCED ON A INCH INCH GRID. IT WILL BE ASSUMED THAT THE SPACING BETWEEN POINTS IN BOTH DIRECTIONS IS THE SAME AND EQUAL FOR ALL POINTS	
SUBROUTINES REQUIRED	
NONE	
METHOD	
THE CONTOUR LEVELS ARE DETERMINED BY SIMPLE LINEAR	



```

C C C C C C C
      INTERPOLATION FROM THE FOUR SURROUNDING PIONTS.
      *****
MET00550
MET00560
MET00570
MET00580
MET00590
MET00600
MET00610
MET00620

MTMPII SUBROUTINE FOR ONE-INCH GRID SPACING
OAKES CODE 5105 15 JAN 69

      *****
SUBROUTINE MTMPII(Y,N,M,T,BND,AZ,BZ,AMIN,IJT,ICON)
REAL*4 IH,KG,ITJZ
DIMENSION A(140),B(140),C(140),D(140),IH(20),Y(N,M),TP(10),TPX(10)
Z,TPM(10),XMT(10),BTM(10),BTX(10),BT(10),KG(10),T(24)
DIMENSION E(140),F(140),G(140),H(140)

DATA DUE/4H /,EPL/4H+ /,EMI/4H- /,IH/1H0,1H /,1H1,1H /,1H2,
11H,1H3,1H /,1H4,1H,1H5,1H,1H6,1H,1H7,1H,1H8,1H /,1H9,1H /,KG/
21H0,1H1,1H2,1H3,1H4,1H5,1H6,1H7,1H8,1H9/,BLK/4H /

YMIN=Y(1,1)
YMAX=Y(1,1)
DO 20 I=1,M
DO 10 J=1,N
YMIN=AMIN1(YMIN,Y(J,I))
YMAX=AMAX1(YMAX,Y(J,I))
10 CONTINUE
20 DELY=YMAX-YMIN
IF(BND) 25,25,30
BND=DELY/15.0
30 IF (AMIN-YMIN) 31,31,32
31 IF (IJT) 33,32,33
32 PD=YMIN/BND
PF=ABS(PD-INT(PD))
IF (YMIN) 2,1,1
1 AMIN=YMIN-PF*BND
GO TO 33
2 AMIN=YMIN-(1.0-PF)*BND
33 AHLD=AZ
IF(AZ) 55,35,55
35 SM=AMAX1(ABS(YMIN),ABS(YMAX))
NS=0
40 NS=NS+1
SM=10.0*SM
IF(SM-1.0)40,50,45
45 NS=NS-1

```





```

SM=SM/10.0
IF(SM-1.0)50,50,50,45
50 AHLD=10.0*NS
55 HBND=BND/2.0
PRINT 70
PRINT 6,T
6 FORMAT(5X,24A4,/)
PRINT 57,AHLD,BZ
57 FORMAT(1H0,65H THE FOLLOWING TRANSFORMATION WAS PERFORMED ON THE IN
1 PUT MATRIX /5X,1H(,E12.5,8H*Y(I,J)+,E12.5,1H) //2X,73HAND THREE
2 DIGITS TO THE RIGHT OF THE DECIMAL POINT ARE PRINTED IN THE MAP )
MET0101010
MET0101020
MET0101030
MET0101040
MET0101050
MET0101060
MET0101070
MET0101080
MET0101090
MET0101100
MET0101110
MET0101120
MET0101130
MET0101140
MET0101150
MET0101160
MET0101170
MET0101180
MET0101190
MET0101200
MET0101210
MET0101220
MET0101230
MET0101240
MET0101250
MET0101260
MET0101270
MET0101280
MET0101290
MET0101300
MET0101310
MET0101320
MET0101330
MET0101340
MET0101350
MET0101360
MET0101370
MET0101380
MET0101390
MET0101400
MET0101410
MET0101420
MET0101430
MET0101440
MET0101450
MET0101460
MET0101470
MET0101480

C
54 PRINT 54,YMAX,YMIN
54 FORMAT(/4X,5HYMAX=,E15.7,5X,5HYMIN=,E15.7)
5 IF (ICON)5,58,5
5 PRINT 11,BND
11 FORMAT(2X,17H THE BAND WIDTH IS,E12.5,6H UNITS //4X,14H CONTOUR LEVE
1LS
I=0
YTOP=AMIN
IF(ABS(YMIN-YMAX)-100.0*BND)53,53,58
53 YB=YTOP
YTOP=YTOP+BND
I=I+1
J=MOD(I,20)
ITJZ=IH(J)
IF(YB-YMAX)59,58,58
59 PRINT 61,YB,YTOP,ITJZ
61 FORMAT(/4X,E10.3,4H TO ,E10.3,2H =,1X,A1)
GO TO 53
58 NCCP=0
NCP=0
PRINT 70
70 FORMAT(1H1)
PRINT 6,T
NLINE=0
NCCP=NCP+1
NCP=NCP+13
73 IF(NCP-M)80,80,75
75 NCP=M
80 CONTINUE
J=-2
NLINE=NLINE+1
LLINE=N-NLINE+1
C SET UP HEADING
85 IF(NCCP-1) 85,85,90
90 J=-1
DO 100 I = 1,135

```



```

      A(I)=BLK
      B(I)=BLK
      H(I)=BLK
100  CONTINUE
110  DO 160 L=NCCP,NCP
      J=J+8
      IF(KI-100) 130,120,120
120  LL=KI/100
      A(J)=KG(LL+1)
      KI=KI-100*LL
      GO TO 135
130  A(J)=KG(1)
135  J=J+1
      IF(KI-10) 150,140,140
140  LL=KI/10
      A(J)=KG(LL+1)
      KI=KI-10*LL
      GO TO 155
150  A(J)=KG(1)
155  J=J+1
      A(J)=KG(KI+1)
160  CONTINUE FIRST ROW OF ARRAY
      C SETUP
      GO TO 260
170  NLINE=NLINE+1
      LLINE=N-NLINE+1
      IF(NLINE-N) 180,180,380
180  DO 190 I=1,135
      A(I)=BLK
      B(I)=BLK
      C(I)=BLK
      D(I)=BLK
      E(I)=BLK
      F(I)=BLK
      G(I)=BLK
      H(I)=BLK
190  CONTINUE
195  IF (ICON)195,260,195
      NCV=NCCP-1
      J=4
      IF(NCY)200,200,210
200  J=5
210  NCV=NCV+1
      IF(NCV-NCP) 220,220,260
220  IF(NCV-M) 230,260,260
230  NLINE=NLINE-1
      YD1=Y(NLINE,NCY)-Y(NLINE+1,NCY)

```

```

MET01490
MET01500
MET01510
MET01520
MET01530
MET01540
MET01550
MET01560
MET01570
MET01580
MET01590
MET01600
MET01610
MET01620
MET01630
MET01640
MET01650
MET01660
MET01670
MET01680
MET01690
MET01700
MET01710
MET01720
MET01730
MET01740
MET01750
MET01760
MET01770
MET01780
MET01790
MET01800
MET01810
MET01820
MET01830
MET01840
MET01850
MET01860
MET01870
MET01880
MET01890
MET01900
MET01910
MET01920
MET01930
MET01940
MET01950
MET01960

```



```

YD2=Y(NLINE,NCY+1)-Y(NLINE+1,NCY+1)
TP(1)=Y(NLINE,NCY)-0.125*YD1
TPX(1)=Y(NLINE,NCY)-0.250*YD1
TPM(1)=Y(NLINE,NCY)-0.375*YD1
XMT(1)=Y(NLINE,NCY)-0.500*YD1
BTM(1)=Y(NLINE,NCY)-0.625*YD1
BTX(1)=Y(NLINE,NCY)-0.750*YD1
BT(1)=Y(NLINE,NCY)-0.875*YD1
TPX(10)=Y(NLINE,NCY+1)-0.125*YD2
TPM(10)=Y(NLINE,NCY+1)-0.250*YD2
XMT(10)=Y(NLINE,NCY+1)-0.375*YD2
BTM(10)=Y(NLINE,NCY+1)-0.500*YD2
BTX(10)=Y(NLINE,NCY+1)-0.625*YD2
BT(10)=Y(NLINE,NCY+1)-0.750*YD2
NLINE=NLINE+1
D1=0.1***(TP(10)-TP(1))
D2=0.1***(TPX(10)-TPX(1))
D3=0.1***(TPM(10)-TPM(1))
D4=0.1***(XMT(10)-XMT(1))
D5=0.1***(BTM(10)-BTM(1))
D6=0.1***(BTX(10)-BTX(1))
D7=0.1***(BT(10)-BT(1))
DO 240 I=2,9
  TP(I)=TP(I-1)+D1
  TPM(I)=TPM(I-1)+D2
  XMT(I)=XMT(I-1)+D3
  BTM(I)=BTM(I-1)+D4
  BTX(I)=BTX(I-1)+D5
  BT(I)=BT(I-1)+D6
CONTINUE
DO 250 I=1,10
  J=J+1
  I1=MOD(IFIX((TP(I)-AMIN)/BND),20)+1
  I2=MOD(IFIX((TPX(I)-AMIN)/BND),20)+1
  I3=MOD(IFIX((TPM(I)-AMIN)/BND),20)+1
  I4=MOD(IFIX((XMT(I)-AMIN)/BND),20)+1
  I5=MOD(IFIX((BTM(I)-AMIN)/BND),20)+1
  I6=MOD(IFIX((BTX(I)-AMIN)/BND),20)+1
  I7=MOD(IFIX((BT(I)-AMIN)/BND),20)+1
  A(J)=IH(I1)
  B(J)=IH(I2)
  C(J)=IH(I3)
  D(J)=IH(I4)
  E(J)=IH(I5)
  F(J)=IH(I6)
  G(J)=IH(I7)

```

```

MET01970
MET01980
MET01990
MET02000
MET02010
MET02020
MET02030
MET02040
MET02050
MET02060
MET02070
MET02080
MET02090
MET02100
MET02110
MET02120
MET02130
MET02140
MET02150
MET02160
MET02170
MET02180
MET02190
MET02200
MET02210
MET02220
MET02230
MET02240
MET02250
MET02260
MET02270
MET02280
MET02290
MET02300
MET02310
MET02320
MET02330
MET02340
MET02350
MET02360
MET02370
MET02380
MET02390
MET02400
MET02410
MET02420
MET02430
MET02440

```



```

250 CONTINUE
260 GO TO 210
265 NCY=NCCP-1
270 J=-2
275 IF(NCY) 265,265,270
280 J=-1
285 GO TO 330
290 NCY=NCY+1
295 IF(NCY-NCP) 280,280,310
300 J=J+7
305 THLD=AHLD*Y(NLINE,NCY)+BZ
310 IF(THLD) 285,290,290
315 H(J)=EMI
320 GO TO 295
325 H(J)=EPL
330 NUM=INT(ABS(THLD-INT(THLD)))*1000.0+0.5)
335 NDS=100
340 DO 300 KK=1,3
345 J=J+1
350 KI=NUM/NDS
355 H(J)=KG(KI+1)
360 NUM=NUM-KI*NDS
365 NDS=NDS/10
370 CONTINUE
375 GO TO 270
380 IF(NCP-M) 360,320,320
385 IF(J-127)330,330,360
390 J=J+3
395 KI=NLINE
400 IF(KI-100) 340,335,335
405 LL=KI/100
410 H(J)=KG(LL+1)
415 KI=KI-100*LL
420 GO TO 343
425 H(J)=KG(1)
430 J=J+1
435 IF(KI-10) 350,345,345
440 LL=KI/10
445 H(J)=KG(LL+1)
450 KI=KI-10*LL
455 GO TO 355
460 H(J)=KG(1)
465 J=J+1
470 H(J)=KG(KI+1)
475 J=J-5
480 IF(NCY-1) 270,270,360
485 IF(NLINE-1)362,362,368
490 PRINT 370,(A(I),I=1,132),(B(IP1),IP1=1,132),(H(IP2),IP2=1,132)

```

```

MET02450
MET02460
MET02470
MET02480
MET02490
MET02500
MET02510
MET02520
MET02530
MET02540
MET02550
MET02560
MET02570
MET02580
MET02590
MET02600
MET02610
MET02620
MET02630
MET02640
MET02650
MET02660
MET02670
MET02680
MET02690
MET02700
MET02710
MET02720
MET02730
MET02740
MET02750
MET02760
MET02770
MET02780
MET02790
MET02800
MET02810
MET02820
MET02830
MET02840
MET02850
MET02860
MET02870
MET02880
MET02890
MET02900
MET02910
MET02920

```





```

368 GO TO 170
    PRINT 370, (A(I), I=1,132), (B(IP1), IP1=1,132), (C(IP2), IP2=1,132),
    1(D(IP3), IP3=1,132), (E(IP4), IP4=1,132), (F(IP5), IP5=1,132),
    2(G(IP6), IP6=1,132), (H(IP7), IP7=1,132)
370 FORMAT(132A1)
380 GO TO 170
    DO 390 I=1,135
    A(I)=BLK
    B(I)=BLK
    C(I)=BLK
    D(I)=BLK
390 CONTINUE
    J=-2
    IF(NCCP-1) 395,395,400
395 J=-1
400 DO 430 L=NCCP,NCP
    J=J+8
    KI=L
    IF(KI-100) 410,405,405
405 LL=KI/100
    C(J)=KG(LL+1)
    KI=KI-100*LL
    GO TO 412
410 C(J)=KG(1)
412 J=J+1
415 IF(KI-10) 420,415,415
    LL=KI/10
    C(J)=KG(LL+1)
    KI=KI-10*LL
    GO TO 422
420 C(J)=KG(1)
422 J=J+1
430 C(J)=KG(KI+1)
    CONTINUE
    PRINT 370, (B(IP1), IP1=1,132), (C(IP2), IP2=1,132)
    IF(NCP-M) 60,500,500
500 RETURN
    END

```

MET02930  
 MET02940  
 MET02950  
 MET02960  
 MET02970  
 MET02980  
 MET02990  
 MET03000  
 MET03010  
 MET03020  
 MET03030  
 MET03040  
 MET03050  
 MET03060  
 MET03070  
 MET03080  
 MET03090  
 MET03100  
 MET03110  
 MET03120  
 MET03130  
 MET03140  
 MET03150  
 MET03160  
 MET03170  
 MET03180  
 MET03200  
 MET03210  
 MET03220  
 MET03230  
 MET03240  
 MET03250  
 MET03260  
 MET03270  
 MET03280  
 MET03290  
 MET03300

```

C000021
C .....
C SUBROUTINE GAUSS
C .....
C PURPOSE
C COMPUTES A NORMALLY DISTRIBUTED RANDOM NUMBER WITH A GIVEN
C .....
    GAUS 10
    GAUS 20
    GAUS 30
    GAUS 40
    GAUS 50
    GAUS 60
    GAUS 70

```







CC

```
SUBROUTINE RANDU
PURPOSE
  COMPUTES UNIFORMLY DISTRIBUTED RANDOM REAL NUMBERS BETWEEN
  0 AND 1.0 AND RANDOM INTEGERS BETWEEN ZERO AND 2**31. EACH
  ENTRY USES AS INPUT AN INTEGER RANDOM NUMBER AND PRODUCES
  A NEW INTEGER AND REAL RANDOM NUMBER.
USAGE
  CALL RANDU(IX,IY,YFL)
DESCRIPTION OF PARAMETERS
  IX - FOR THE FIRST ENTRY THIS MUST CONTAIN ANY ODD INTEGER
  NUMBER WITH NINE OR LESS DIGITS. AFTER THE FIRST ENTRY,
  IX SHOULD BE THE PREVIOUS VALUE OF IY COMPUTED BY THIS
  SUBROUTINE.
  IY - A RESULTANT INTEGER RANDOM NUMBER REQUIRED FOR THE NEXT
  ENTRY TO THIS SUBROUTINE. THE RANGE OF THIS NUMBER IS
  BETWEEN ZERO AND 2**31
  YFL-- THE RESULTANT UNIFORMLY DISTRIBUTED, FLOATING POINT,
  RANDOM NUMBER IN THE RANGE 0 TO 1.0
REMARKS
  THIS SUBROUTINE IS SPECIFIC TO SYSTEM/360 AND WILL PRODUCE
  2**29 TERMS BEFORE REPEATING. THE REFERENCE BELOW DISCUSSES
  SEEDS (65539 HERE), RUN PROBLEMS, AND PROBLEMS CONCERNING
  RANDOM DIGITS USING THIS GENERATION SCHEME. MACLAREN AND
  MARSAGLIA, JACM 12, P. 83-89, DISCUSS CONGRUENTIAL
  GENERATION TYPE, ONE FILLING A TABLE AND ONE PICKING FROM THE
  TABLE, IS OF BENEFIT IN SOME CASES. 65549 HAS BEEN
  SUGGESTED AS A SEED WHICH HAS BETTER STATISTICAL PROPERTIES
  FOR HIGH ORDER BITS OF THE GENERATED DEVIATE.
  SEEDS SHOULD BE CHOSEN IN ACCORDANCE WITH THE DISCUSSION THAT
  GIVEN IN THE POINT RANDOM NUMBERS ARE DESIRED, AS ARE
  IF FLOATING POINT RANDU, THE RANDOM CHARACTERISTICS OF THE
  AVAILABLE FROM RANDU, THE RANDOM CHARACTERISTICS OF THE
  FLOATING POINT DEVIATES ARE MODIFIED AND IN FACT THESE
  DEVIATES HAVE HIGH PROBABILITY OF HAVING A TRAILING LOW
  ORDER ZERO BIT IN THEIR FRACTIONAL PART.
SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED
  NONE
METHOD
  POWER RESIDUE METHOD DISCUSSED IN IBM MANUAL C20-8011,
  RANDOM NUMBER GENERATION AND TESTING
```



```

C ..... RAND 520
C ..... RAND 530

SUBROUTINE RANDU(IX,IY,YFL)
IY=IX*65539
IF(IY) 5,6,6
5 IY=IY+2147483647+1
6 YFL=IY
YFL=YFL*.4656613E-9
RETURN
END

RAND 540
RAND 550
RAND 560
RAND 570
RAND 580
RAND 590
RAND 600

```





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14

## KEY WORDS

## LINK A

## LINK B

## LINK C

ROLE

WT

ROLE

WT

ROLE

WT

fin-flat plate flow field

wing-root flow field

holography

three-dimensional density field

three-dimensional flow field

fin

flat plate

interferogram

flow field

fin flat plate junction

wing root junction

holographic interferometry

supersonic flow



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